

**Energy and vapor transfer in evaporation processes over saturated bare soil and water  
surface-lysimeter studies**

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## Abstract

Potential evaporation (PE) is a significant input in many hydrological numerical models for the estimation of actual evaporation. Evaporation from water ( $PE_w$ ) is generally considered equivalent to evaporation from saturated bare soils ( $PE_s$ ). The influences of the underlying surface on PE as well as the energy and vapor transfer in potential evaporation processes over different surfaces are rarely discussed. In this research, lysimeter experiments were set up to measure the diurnal cycles of evaporation from two saturated sandy soils and water at a high temporal resolution in the Guanzhong Basin, China. Evaporation from Class A Pan, meteorological variables and temperatures were also measured during the experiment. Observation results show that  $PE_s$  is ~12% higher than  $PE_w$  on a yearly scale. There were also some clear differences in diurnal and seasonal PE dynamics between saturated bare soils and water. In summer,  $PE_s$  is higher than  $PE_w$  at day but smaller at night, with the peak value of  $PE_w$  lagging ~4 hours behind  $PE_s$ .

These observed PE dynamics and energy transfer processes can be quantitative explained on the basis of a full analysis of the energy balance equation. A comprehensive description of the flux transfer processes showed that these differences in PE are governed by differences in available energy (including albedo and thermal properties) between soils and water. Moreover, the observed differences in PE and vapor transfer processes were reproduced and described by improving the vapor diffusion equation, with considering the influence of different surfaces and boundary layer thicknesses. It is found that dynamics in PE were mainly characterized by surface temperature, which further determined the vapor gradients between the evaporation surface and air flow. The results suggested that differences between  $PE_s$  and  $PE_w$  cannot be neglected in hydrological applications. This study can act as both an experimental and theoretical reference for estimating potential evaporation rates.

**Keywords:** saturated bare soil evaporation, water evaporation, lysimeter, energy balance, vapor diffusion

46    **Abbreviations**

47    **PE**   potential evaporation

48    **PE<sub>s</sub>**   evaporation from saturated bare soils

49    **PE<sub>fine</sub>**   evaporation from saturated fine sand

50    **PE<sub>coarse</sub>**   evaporation from saturated coarse sand

51    **PE<sub>w</sub>**   evaporation from water surfaces

52    **PE<sub>col</sub>**   evaporation from water filled lysimeter column

53    **PE<sub>pan</sub>**   evaporation from water filled Class A Pan

54

## 1. Introduction

Evaporation from a surface with sufficient water supply is known as potential evaporation (PE). PE is an important variable considered in many disciplines and numerical models (hydrology, hydrogeology, agriculture, and land surface modelling) to estimate actual evaporation ( $E_a$ ) from soil (Ben Neriah et al., 2014; Bormann, 2011; Han et al., 2014; Kirono et al., 2010) or from phreatic water (Zhidong et al., 1988). Although several definitions of PE have been proposed (J.-P. Lhomme, 1997), they generally fall within two categories according to the surface type. They are evaporation from saturated bare soils ( $PE_s$ ) and evaporation from water surface ( $PE_w$ ).

$PE_s$  is close to the Stage I Evaporation (early stages of evaporation from saturated soils with an initially high and relatively constant rate).  $PE_s$  is considered to be driven by atmosphere forcings and soil properties (Allen et al., 2005; Eberbach et al., 2011). However, in those experiments which investigate how soil type affects  $PE_s$ , soil columns were only saturated initially. There remain limited quantitative studies on the influence of soil texture on evaporation rate under a constantly saturated soil. In fact,  $PE_s$  and  $PE_w$ , as well as the evaporation rate from pan ( $PE_{pan}$ ) have all been used as PE to estimate actual evaporation rate in applications (Allen et al., 1989; Granger, 1989; Kirono et al., 2010; Lehmann & Or, 2013; J. P. Lhomme & Guilioni, 2010; L. Wang et al., 2021; Westra, 2009). While this simplification relies on several hypotheses, differences in PE over saturated textures and water have not been quantitatively explored, nor have these differences been compared with  $PE_{pan}$  based on measured results at a high temporal resolution.

PE is a highly dynamic process reflecting the interplay between atmospheric conditions (energy input related to meteorological elements) and flow process (vapor diffusion) (Lim et al., 2012, 2013, 2016; Shahraeeni et al., 2012). The energy transfer process over different evaporation surfaces can be expressed based on the energy balance equation (Hamdani et al., 2018; Lim et al., 2013). However, applications sometimes adopt a simplified flux term in the full-form energy balance equation. For example, since there are rarely measurements of total ground heat flux for soil  $G$  (or the heat storage variation for water  $N$ ) at meteorological stations,

they are often ignored within daily calculations of PE or assumed to be a fraction of net radiation over a sub-daily scale (Allen et al., 2005). Since  $G$  and  $N$  must be considered over shorter periods, the dynamics of evaporation should be described by a comprehensive analysis of the energy balance, considering the influences of different surfaces on flux terms.

Vapor from saturated bare soil (water filled porous) and water is mainly transported by diffusion across the boundary layer above the evaporation surface (Bird et al., 1961; Haghighi et al., 2013; Lehmann & Or, 2013; Shahraeeni et al., 2012). Solutions based on diffusion transport (such as Fick's First Law of Diffusion) determine that evaporation rate is characterized by vapor concentration gradients between evaporation surface and the flowing air as well as the thickness of the boundary layer  $\delta$  (Ben Neriah et al., 2014; Dalton, 1802; Haghighi & Or, 2013; Lehmann et al., 2018; Lehmann & Or, 2013; Lim et al., 2012; Schlünder, 1988). However, previous research (Haghighi et al., 2013) has determined that  $\delta$  can be defined by velocity convections. There have been no previous comparisons of the differences in  $\delta$  among saturated soil and water.

Energy transfer and vapor transfer are inter-dependent processes essential for characterizing evaporation dynamics. These two drivers are thought to be influenced by the properties of evaporation surfaces, and are regulated by various factors including available energy, vapor pressure, and climate elements. To the best of our knowledge, exploration of this two-component interpretation in evaporation process between saturated bare soils and water over sub-daily scale for semi-arid region is a research gap, thus should be addressed in detail.

The present study implemented a unique lysimeter experiment in which evaporation over saturated fine sand ( $PE_{\text{fine}}$ ), coarse sand ( $PE_{\text{coarse}}$ ), water in the lysimeter columns ( $PE_{\text{col}}$ ) as well as water in the Class A Pan ( $PE_{\text{pan}}$ ) were measured at a high temporal resolution in the Guanzhong Basin to investigate the following objectives:

1. Investigating the differences of total evaporation over one year as well as the differences of sub-daily transient evaporation in different seasons between saturated fine sand, coarse sand and water surfaces.
2. Characterizing the differences in PE from a fully analysis of the energy balance and the vapor transfer in evaporation processes.

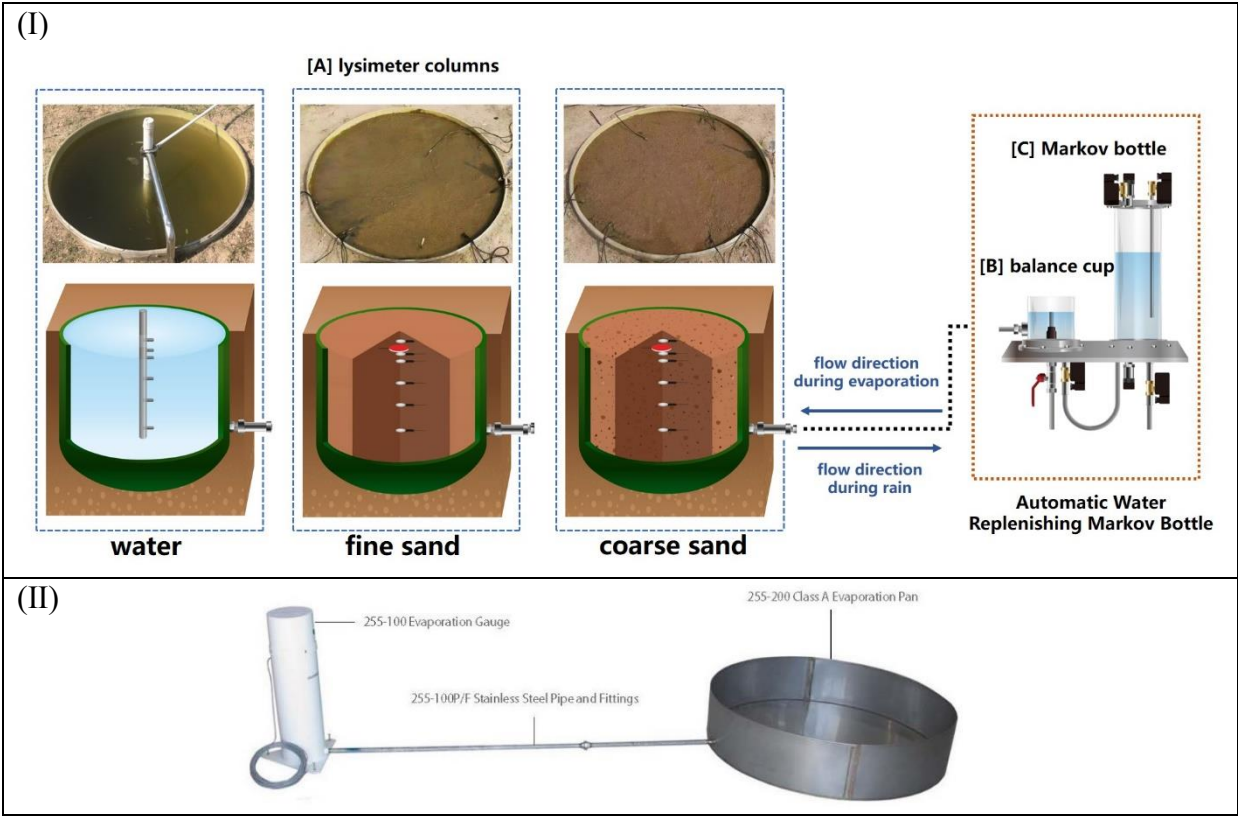
111 3. Evaluating the PE-estimation methods for characterizing the variables regulating energy  
 112 transfer and vapor transfer processes over saturated bare soils and water.  
 113

114 **2 Materials and methods**

115 **2.1 Materials**

116 **2.1.1 Lysimeter columns**

117 The lysimeter experiment was conducted at Chang’an University, Xi’an, Guanzhong Basin,  
 118 China (Latitude: 34°28’ N; Longitude: 108°93’ E). The setup of the experiments is show in  
 119 Figure 1.



120 *Figure 1: (I) Setup of the lysimeter experiment. (II) Setup of the Class A Pan.*

121 Lysimeters allow to measure evaporation rates directly (Allen, Howell, et al., 1991; Allen,  
 122 Pruitt, et al., 1991; Goss & Ehlers, 2010; Pütz et al., 2018). As Figure 1 (I) illustrates, lysimeter  
 123 columns (area 1.0 m<sup>2</sup>, depth 70cm) are filled with water and two different homogeneous soil  
 124 textures. Evaporation rate from saturated soil textures (PE<sub>s</sub>) and water (PE<sub>w</sub>) were observed by  
 125 the “Automatic Water Replenishing Markov Bottle” (installed in the basement). This Markov

126 Bottle based lysimeter device got a Utility Model Patent in 2014 (ZL2014200056418.x), and can  
127 provide reliable potential evaporation rate directly (Jensen et al., 2016; Truong, 2012; K. Wang  
128 & Dickinson, 2012).

129 The bottom of each column (A) was connected to a Markov Bottle (C) via the balance cup  
130 (B). The water table in the columns can be maintained at the soil and water surface so that  
131 potential evaporation rate equals to mass losses of water from the Markov Bottle to the lysimeter.  
132 In this experiment, a drop of the water level in the Markov Bottle of 94 cm coincides with an  
133 evaporation rate in the lysimeter column of 1 cm. The water level variation in the Markov Bottle  
134 were automatically measured using pressure sensors (MPM489, Microsensor).

135 Temperatures (for soil: MPS 6, for water: 1/3DIN) in the columns were observed at 3, 5, 10,  
136 20, 30 and 50 cm depth. Heat flux  $G_0$  (HFP01, Hukseflux Inc.) were observed at 5 cm depth in  
137 the soil column. Data are automatically recorded by the data logger (CR1000, Campbell  
138 Scientific Inc. USA) every 5 minutes.

139

#### 140 **2.1.2 Class A Pan evaporation and meteorological stations**

141 As Figure 1 (II) illustrates, the Class A Pan (diameter 120.6 cm, depth 25 cm) was equipped  
142 with an evaporation gauge (255-100 Novalynx Analog Output Evaporation Gauge and 255-200  
143 Novalynx Class A Evaporation Pan, Campbell Scientific Inc. USA). The pan evaporation rate  
144 can be measured directly by recording the water level variations. Temperature sensors were  
145 placed at the center of the pan (10 cm away from the bottom) to measure the bulk water  
146 temperature. Temperatures on water profile were measured at 3, 10 and 15cm depth below the  
147 surface.

148 The standard meteorological station is located in the southwestern part of the experimental  
149 site. Humidity-temperature sensors (083E-1-6, Met One Instrument, USA), air pressure sensor  
150 (Vaisala PTB110 Barometer-CS106, Campbell Scientific Inc. USA) and net radiometer (CNR4,  
151 Kipp&Zonen, Netherlands) were installed at 1.5 m above the ground. Wind speed sensors were  
152 located at 2 m and 10 m respectively. The measured data were automatically recorded every 5

minutes by CR-3000 (Campbell Scientific Inc. USA). The present study determined PE and other parameters at a 1-hour temporal resolution.

## 2.2 Methods and theory

### 2.2.1 Energy balance equation

The energy balance between an evaporation surface and the atmosphere can be expressed as:

$$\begin{aligned} R_{n,s} - G &= LE_s + H_s \\ R_{n,w} - N &= LE_w + H_w \end{aligned} \quad (1)$$

where  $R_n$  is the net radiative flux ( $\text{W}\cdot\text{m}^{-2}$ ),  $G$  is the total ground heat flux ( $\text{W}\cdot\text{m}^{-2}$ ),  $N$  is the variation in water heat storage ( $\text{W}\cdot\text{m}^{-2}$ ),  $H$  is the convective energy exchanged between air and the evaporation surface ( $\text{W}\cdot\text{m}^{-2}$ ),  $LE$  is the latent heat flux due to evaporation ( $\text{W}\cdot\text{m}^{-2}$ ), and subscripts  $s$  and  $w$  refer to soil and water, respectively.

$R_n$  includes net shortwave radiation ( $R_{ns}$ ) and net longwave radiation ( $R_{nl}$ ):

$$R_n = R_{ns} + R_{nl} = (1 - \chi) \cdot R_{sd} + \varepsilon \cdot R_{ld} - R_{lu} \quad (2)$$

Net radiation is a critical input for estimating the evaporation rate (Han et al., 2014). In Equation (2),  $R_{sd}$  ( $\text{W}\cdot\text{m}^{-2}$ ) and  $R_{ld}$  ( $\text{W}\cdot\text{m}^{-2}$ ) represent measured downwards shortwave and downwards longwave radiation at the meteorological station, respectively and  $R_{lu}$  ( $\text{W}\cdot\text{m}^{-2}$ ) is the outgoing longwave radiation emitted from the surface (Lawrence et al., 2019; Saito et al., 2006):

$$R_{lu} = \varepsilon \cdot \sigma \cdot (T + 273.14)^4 \quad (3)$$

where  $\varepsilon$  is the surface emissivity,  $\sigma$  is the Stefan-Boltzman constant ( $=5.67 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$ ) and  $T$  ( $^{\circ}\text{C}$ ) is the surface temperature. The present study used temperatures measured at depths of 3 cm and 5 cm to estimate  $T$  using linear extrapolation.  $R_{ns}$  and  $R_{lu}$  represent the differences in albedo  $\chi$  and surface temperature between soil and water, respectively.

For soil,  $G$  can be estimated as:

$$G = G_0 + \Delta S \quad (4)$$

Soil heat flux  $G_0$  is measured using heat flux plates. The rate of soil heat storage  $\Delta S$  above the heat flux plate is calculated as (Ochsner et al., 2007; Roxy et al., 2014):



$$\Delta S = \Delta z \cdot c_s \cdot \frac{(T_i - T_{i-1})}{\Delta t} \quad (5)$$

where soil temperature measured at 3 cm is used to estimate  $\Delta S$ ,  $T_i - T_{i-1}$  (K) is the temporal variation in temperature (where  $i$  represents time),  $\Delta z$  is the thickness of the layer (5 cm for soil),  $\Delta t$  is the time step,  $c_s$  (MJ/m<sup>3</sup>K) is the volumetric heat capacity for saturated soil texture [Appendix, Equation (A2)], the rate of water heat storage  $N$  is calculated by Equation (5), and  $\Delta z$  is the depth of water in the column ( $\Delta z$  is 25 cm for water in Class A Pan evaporation, the present study divided the depth of the water column into 10-cm layers).

The sensible heat flux  $H$  can be estimated according to a resistance-type formulation (Cellier et al., 1996; Griend & Owe, 1994; Stewart et al., 1994):

$$H = \frac{\rho \cdot c_p \cdot (T - T_a)}{r_a} \quad (6)$$

where  $T_a$  is the air temperature,  $c_p$  is the specific heat of air,  $\rho$  (kg·m<sup>-3</sup>) is the air density, and  $r_a$  (s·m<sup>-1</sup>) is the aerodynamic resistance [given in the Appendix Equation (A3)].

The latent heat flux  $LE$  for evaporation can be estimated from the evaporation rate, where  $L$  is the latent heat of vaporization.

193

### 2.2.2 Vapor transfer equation

Vapor transfer from the evaporation surface to air  $E_{vapor}$  can be expressed based on Fick's First Law of Diffusion and the boundary layer theory (Lehmann et al., 2018; Lim et al., 2012):

$$E_{vapor} = \frac{M_w \cdot D_v \cdot \alpha \cdot [e_s(T) / T - e_a(T_a) / T_a]}{R \cdot \rho_w \cdot \delta} \quad (7)$$

where  $M_w$  (= 0.018 kg·mol<sup>-1</sup>) is the molecular mass of water,  $R$  (= 8.3 J·mol<sup>-1</sup>·K<sup>-1</sup>) is the ideal gas constant,  $\rho_w$  (= 1,000 kg·m<sup>-3</sup>) is the water density,  $\alpha$  is continuous water availability at the surface, for saturated texture  $\alpha$  equals to the saturated content  $\theta_s$ , for water  $\alpha$  equals to 1.  $D_v$  (m<sup>2</sup>·s<sup>-1</sup>) is the diffusion coefficient for water vapor in air (Ben Neriah et al., 2014),  $e_s$  is the vapor pressure at the surface (Pa), which is assumed to be equivalent to the saturated vapor pressure at the surface temperature  $T$ ,  $e_a(T_a)$  is the actual vapor pressure of air at the reference

height, and  $e_s(T)/T - e_a(T_a)/T_a$  represents the vapor pressure gradient between the surface and air ( $\Delta e$ ).

According to the boundary layer theory,  $\delta$  (mm) is the thickness of the boundary layer above the surface. Quantifying  $\delta$  is difficult due to complex “mixed convection” between porous media, water, and the atmosphere, an alternative empirical calculation of  $\delta$  is given as:

$$\delta = a \cdot U^{-0.5} \quad (8)$$

where  $\delta$  is inversely proportional to the square root of the free-stream velocity  $U$  (Ben Neriah et al., 2014; Ghezzehei et al., 2004; Haghighi et al., 2013; Shahraeeni et al., 2012),  $U$  ( $\text{m} \cdot \text{s}^{-1}$ ) can be represented as the wind velocity at 2 m. Previous studies have determined  $\delta$  based on velocity convections in which  $a = 3.9$  and  $a = 2.26$  were assumed under laminar flow and turbulent flow respectively (Ben Neriah et al., 2014; Haghighi & Or, 2013; Shahraeeni et al., 2012). Parameter  $a$  is generally determined through calibration against measured data. The present study selected  $a=2.06$  and  $a=5.8$  for saturated textures and the water surface respectively, to close the measured PE under consideration of monthly sums of PE and diurnal courses of PE. The selection of parameter  $a$  in the present study will be introduced in Results section.

The calculations for determining the parameters for Equation (1) and Equation (7) are provided in the Appendix [Appendix Equations (A1–A7)].

### 3. Results

#### 3.1 Meteorological condition of the research area

Figure 2 shows the average diurnal cycles of meteorological elements across different seasons during the study period.

Meteorological variables typically showed diurnal cycles. The peak yearly  $T_a$  occurred in summer due to higher incoming radiation, whereas the coldest days occurred between December and January. The peak  $T_a$  lagged the  $R_n$  due to thermal inertia of the air. Relative humidity ( $RH$ ) showed an inverse relationship to  $T_a$ . Wind speed showed a clear diurnal cycle with higher value during the day and lower value at night. The study area generally showed a relatively low wind speed. Consistent with Equation (A7), the pattern within the vapor pressure ( $e_a$ ) curve differed

from those of other meteorological variables, with a peak in the mid-morning, besides for in winter during which the peak occurred in the late morning. Net radiation ( $\text{W}\cdot\text{m}^{-2}$ ) curves are in bell-shape,  $R_n$  peaks around 12:40 (Beijing time). Since there was virtually no incoming shortwave radiation at night, the negative  $R_n$  could be mainly attributed to upward longwave emissivity.

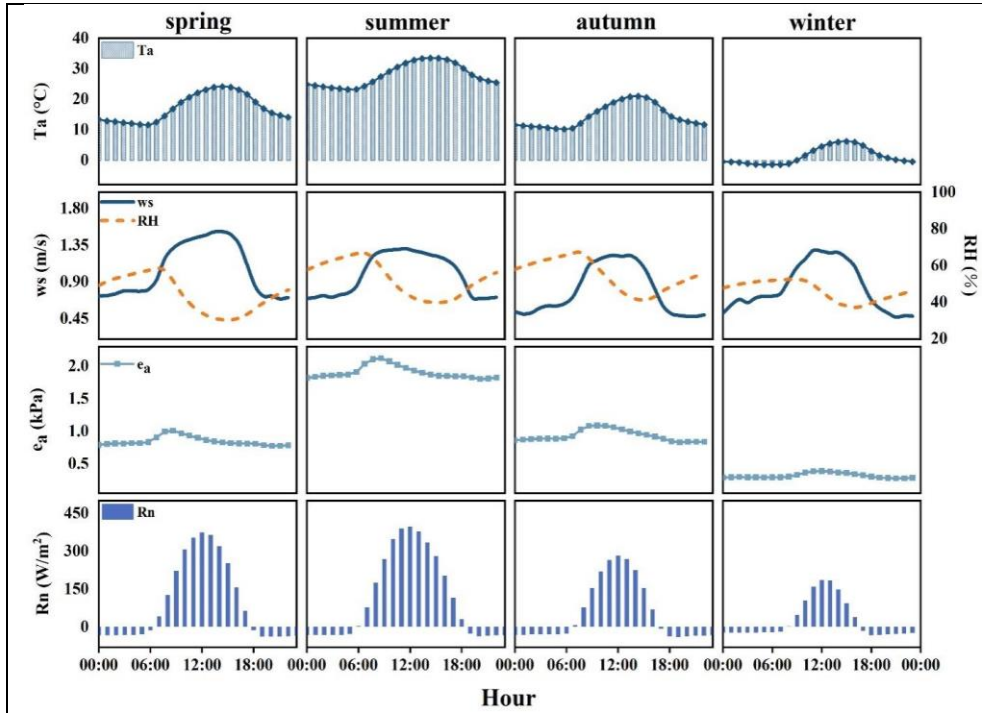
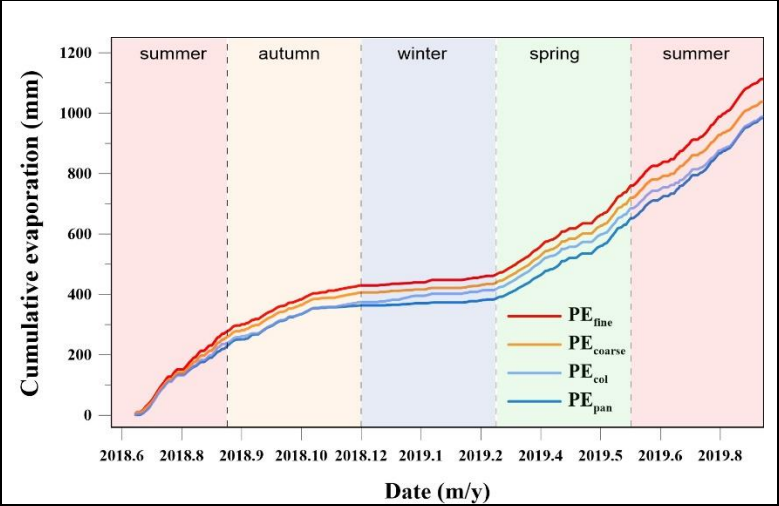


Figure 2: Average diurnal cycles of air temperature, wind speed, relative humidity, actual vapor pressure, and net radiation in different seasons.

### 3.2 Measured PE over saturated bare soils and water

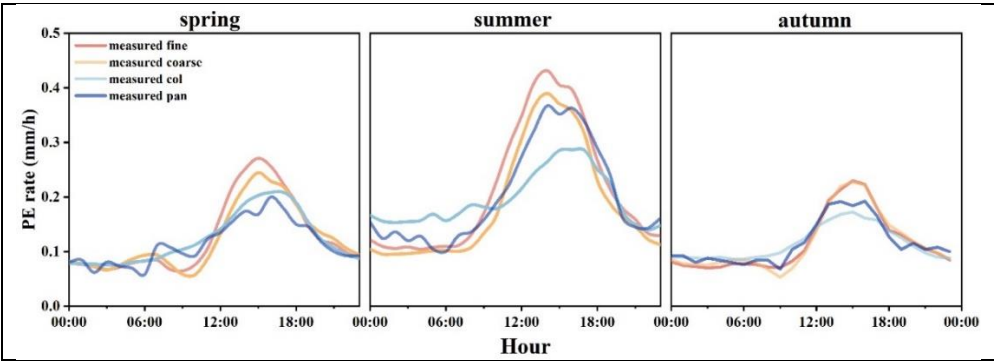
Cumulative evaporation from two saturated sandy soils, water, and Class A Pan over the study period are plotted in Figure 3. There were clear differences in PE among the different surfaces with increasing cumulative evaporation. The maximum and minimum PE were observed for fine sand and water, respectively.  $PE_{\text{fine}}$  exceeded  $PE_{\text{col}}$ ,  $PE_{\text{pan}}$ , and  $PE_{\text{coarse}}$  by 12.6%, 13.1%, and 7.3%, respectively during the measurement period. These differences were more pronounced in summer. For example, from June to August,  $PE_{\text{fine}}$  exceeded  $PE_{\text{col}}$ ,  $PE_{\text{pan}}$ , and  $PE_{\text{coarse}}$  by 15.1%, 11.9%, and 9.4%, respectively. The measured cumulative evaporation curves were parallel to each other from September until the end of November, indicating only

249 minor differences in PE among different surfaces in autumn). The duration of the frost period  
 250 (daily minimum temperature below 0 °C) from December to January was 68 days. The frozen  
 251 conditions in winter effectively prevented any increase in cumulative evaporation.



252 *Figure 3 Cumulative evaporation over saturated fine sand, coarse sand, water, and Class A Pan water over the*  
 253 *study period*

254 The cumulative PE for spring, summer, and autumn was basically equal to that for the year  
 255 in the study period. The present study selected March, August, and October as months  
 256 representative of spring, summer, and autumn, respectively, and conducted detailed analyses for  
 257 these three months. Figure 4 shows the measured average diurnal PE for saturated fine sand,  
 258 coarse sand, water, and Class A Pan water over the three seasons.



259 *Figure 4 Measured diurnal potential evaporation (PE) curves for saturated fine sand, coarse sand, water in column,*  
 260 *and water in Class A Pan over different seasons.*

262 There was a daily peak in the diurnal PE cycles of saturated soils in the afternoon. In spring  
 263 and autumn, there is another “small PE peak” exists after sunrise. The effect of soil texture on

264 PE was more obvious under a high evaporation demand, such as during the daytime in spring  
265 and summer. Moreover, the observations showed considerable nighttime  $PE_s$  and an average  
266 nighttime PE over the two saturated sandy soils exceeding 1.0 mm/day in spring and summer.  
267 Specifically, PE showed diurnal variation, decreasing after sunset and stabilizing until the  
268 second part of the night (Figure 4).

269 There were absolute and temporal differences in the diurnal dynamics of  $PE_s$  and  $PE_w$   
270 curves (Li et al., 2020). The cumulative daily  $PE_s$  exceeded that of  $PE_w$ , with  $PE_s > PE_w$  before  
271 noon till the sunset, and  $PE_s < PE_w$  before dawn until morning. As shown in Figure 4, ~66% of  
272 the cumulative daily  $PE_s$  occurred between sunrise and sunset in August, with the reminding 34%  
273 occurring at night. Nighttime evaporation from soil was ~32% smaller than that during the day.  
274 For water, nighttime evaporation is only ~10% smaller than daytime evaporation. In addition,  
275 the peak values of  $PE_{fine}$  and  $PE_{coarse}$  exceeded that of  $PE_{col}$ , and occurred at ~14:00 for saturated  
276 sandy soils, whereas the peak  $PE_{col}$  (~17:00) lagged ~4 h behind that in  $PE_s$ .

277 The results also showed that Class A Pan evaporation approximated the evaporation in the  
278 deeper lysimeter column on a yearly scale. At a sub-daily scale,  $PE_{pan}$  exceeded  $PE_{col}$  after noon  
279 but was less than  $PE_{col}$  at night. The temporal dynamics of the  $PE_{pan}$  curve were similar to those  
280 of the PE curves for saturated bare soils. The measured  $PE_{pan}$  curve showed fluctuations,  
281 particularly under a smaller evaporation demand (such as between sunset and sunrise). This  
282 result could be attributed to accuracy of the Class A Pan evaporation gauge is 0.1mm, water  
283 level variation smaller than 0.1mm cannot be captured.

284

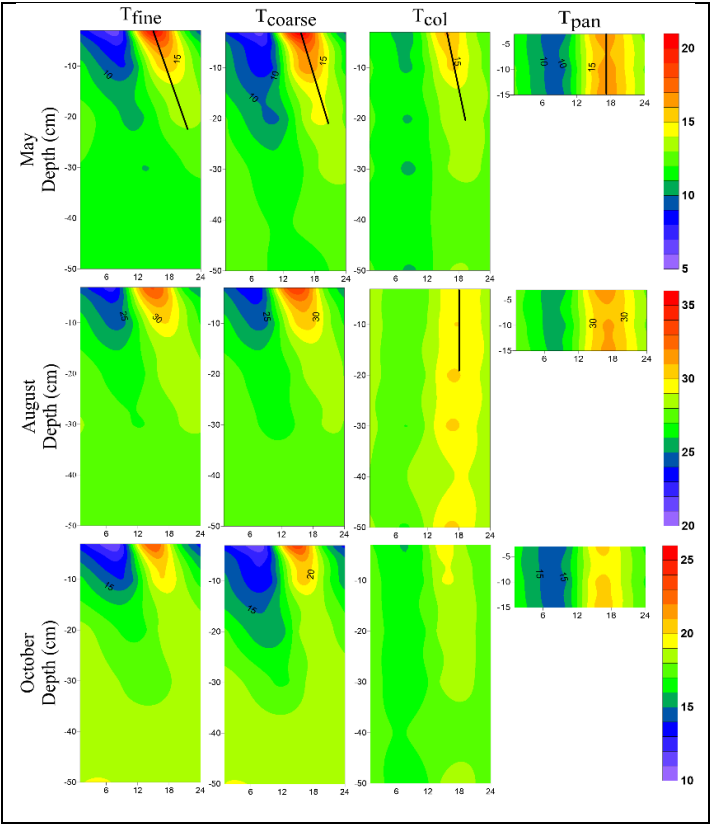
### 285 3.3 Temperature variations over different evaporation surfaces

286 Figure 5 shows the average diurnal temperature variations for saturated fine sand ( $T_{fine}$ ),  
287 coarse sand ( $T_{coarse}$ ), water in the column ( $T_{col}$ ) and water in the Class A Pan ( $T_{pan}$ ) in different  
288 seasons.

289 There were slight differences in surface temperature among the saturated sandy soils. The  
290 variation in surface temperature of coarse sand exceeded that for fine sand due to its smaller heat  
291 capacity. Different to that for saturated textures and Pan water, the surface temperature of water

292 in the column is the lowest at noon, but highest around the dusk.

293 There were considerable differences in temperature distributions with depth among soil and  
294 water. Surfaces with saturated textures showed a delayed downward heat transport compared to  
295 those of other surfaces. As shown in Figure 5, the delay was indicated by an angle, with slower  
296 transmission associated with larger angles. Temperature discretization in water was not as  
297 obvious as that in soil, resulting in a smaller angle in the water column than in saturated textures.  
298 Since  $T_{col}$  is relatively uniform in summer, the gradient angle along the vertical profile in  
299 summer approximated 0 (Figure 5). In particular, shallow Pan water was well mixed with a  
300 vertically homogenous temperature and angle of 0 over the different seasons.



301 *Figure 5 Variations in average diurnal temperature in the soil columns, water column, and water in Class A Pan*  
302 *over different season during the study period and in the study area.*  
303

304 The temperature dynamics observed in the results could be attributed to: (1) a smaller heat  
305 capacity resulting in more obvious variations in soil temperature; (2) higher heat capacity  
306 enhancing water convection, eventually resulting in a more homogeneous vertical temperature in  
307 water; (3) penetration of the water by shortwave radiation, resulting in warming of a larger

vertical column compared to that in soil (incoming radiation does not penetrate soil profile, resulting in heating of only the soil surface).

### 3.4 Energy balance components

Figure 6 shows the average diurnal cycles of energy flux terms [in Equation (1)] over saturated fine sand, coarse sand, water in the column, and Pan water over different seasons.

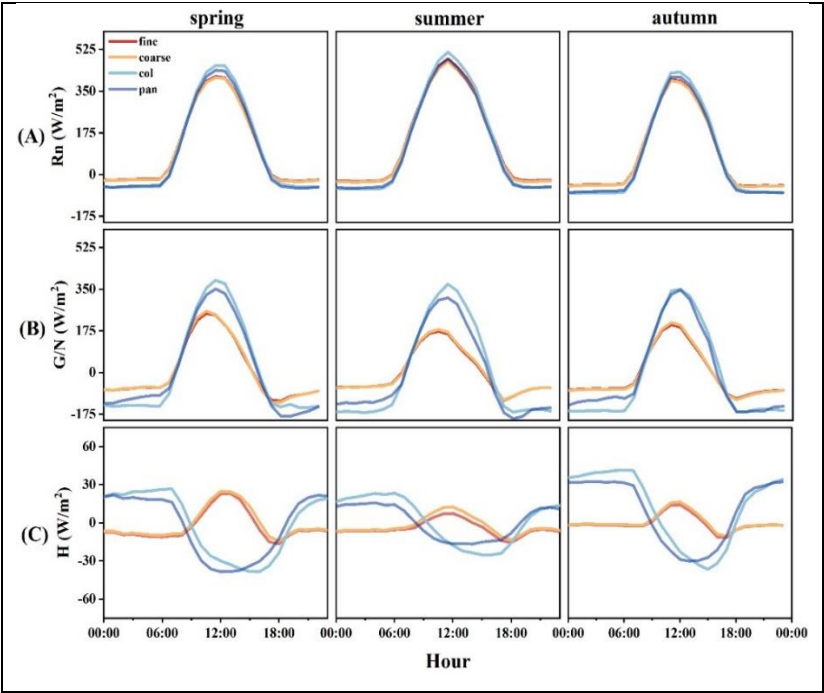


Figure 6 Average diurnal cycles of energy fluxes of the energy balance equation for saturated fine sand, coarse sand, water in column and water in Class A Pan over different seasons: net radiation (A), heat flux/heat storage variation (B), sensible heat flux (C).

#### 3.4.1 Net radiation

As shown in Figure 6A,  $R_n$  was slightly negative at night, following which it increased to a positive value after sunrise (summer: between 5:00 and 6:00, spring and autumn: between 6:00 and 7:00). The value of  $R_n$  exceeded those of other flux terms and was influenced by albedo (which affects net shortwave radiation  $R_{ns}$ ) and surface temperature (which affects upward longwave radiation  $R_{lu}$ ). Based on our measurements, coarse sand is slightly higher in albedo (Appendix Table , which results in smaller  $R_{ns}$ ) and surface temperature (which enhances the outgoing longwave radiation  $R_{lu}$ ) than fine sand. Consequently,  $R_n$  for coarse is slightly smaller

326 than fine sand both during the day and at night. Furthermore, a weaker variation in surface  
327 temperature for column water (Figure 5) resulted a smaller  $R_{lu}$  and a more obvious change in  $R_n$   
328 compared to that of saturated textures and Pan water.

329

### 330 **3.4.2 Heat flux and heat storage variation**

331 Total ground heat flux  $G$  and the variation in water heat storage  $N$  represent the below-  
332 surface redistribution of energy. The values of  $G$  and  $N$  were highly related to variations in  
333 temperature along depths and surface thermal properties. Figure 6B shows the average diurnal  
334 cycles of  $G$  and  $N$  for different surfaces over different seasons.

335 According to Equation (5),  $N$  is determined by temporal temperature variations  $T_i - T_{i-1}$  and  
336 water depth  $\Delta z$ . Under a vertically homogenous temperature (for example in summer), the  
337 amplitude of  $N_{col}$  for the deeper column exceeds that of  $N_{pan}$  due to its larger depth (Figure 6B).  
338 While,  $N_{col}$  and  $N_{pan}$  did not show significant differences during daytime in autumn. This is  
339 because more pronounced temperature variations ( $T_i - T_{i-1}$ ) in Pan will compensate the influence  
340 of smaller  $\Delta z$  for the Class A Pan. As for soils,  $G_{coarse}$  has bigger amplitude than  $G_{fine}$ . This is  
341 related to a more pronounced temperature variation of the coarse sand.  $N$  for water have a  
342 particularly evident amplitude than  $G$ . This is related to the increased heat capacity of water  
343 compared to the soil.

344  $N$  and  $G$  also showed differences in temporal dynamics. Both  $N$  and  $G$  turned from negative  
345 to positive simultaneously in the morning with increasing solar radiation,  $G$  peaked in the late  
346 morning, whereas  $N$  showed an obvious lag response to  $G$  (Figure 6B). Both  $G$  and  $N$  were  
347 positive in the day and negative at night, indicating storage of energy under the surface during  
348 the day and its release at night.

349 On seasonal basis, the amplitudes of  $G$  and  $N$  in spring and autumn exceeded those in  
350 summer, which could be attributed to a stronger heating-cooling process (which results in higher  
351 temperature gradients shown in Figure 5) below the evaporation surface throughout these two  
352 seasons. Over longer time scales, the positive values of  $G$  and  $N$  in spring and summer allowed



the storage of heat, whereas their negative values in autumn and winter allowed this stored heat to dissipate.

### 3.4.3 Sensible heat flux

Figure 6C shows the average diurnal cycles of  $H$  for different surfaces among the different seasons. Daytime (9:00 to 15:00) soil  $H$  was positive since heat was transferred from soil to air, whereas  $H$  was negative during the night as heat was transferred from the atmosphere to the soil surface. In contrast,  $H$  of water was negative during the day (10:00 to 19:00) and positive at night. These opposite trends in nighttime  $H$  between soil and water were regulated by surface temperature. Generally, the absolute value of  $H$  over well-watered surface is the smallest among major fluxes in Equation (1) (Finch, 2001; Gallego-Elvira et al., 2011).

### 3.5 Vapor transfer components

Following Equation (7), PE through vapor diffusion is influenced by wind speed (Figure 7A) and air temperature (Figure 7B). In particular, PE was mostly driven by vapor gradients  $e_s(T)/T - e_a(T_a)/T_a$  (abbreviated to “ $\Delta e$ ”) between the evaporation surface and the thin layer of air above the surface (Figure 7C).

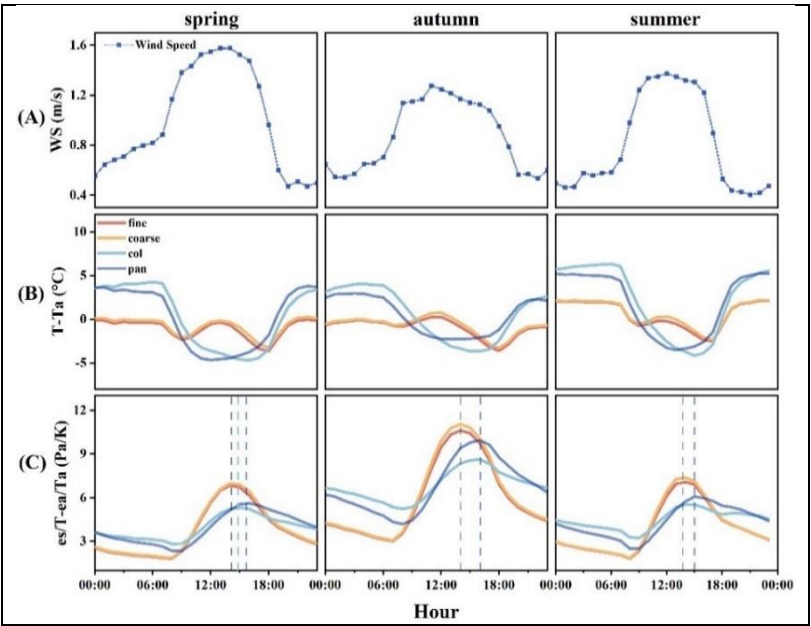


Figure 7 Average diurnal cycles of: (A) wind speed, (B) deficits between surface temperature and air temperature, and (C) vertical vapor gradients ( $\Delta e$ ) between different evaporation surfaces and air over different seasons

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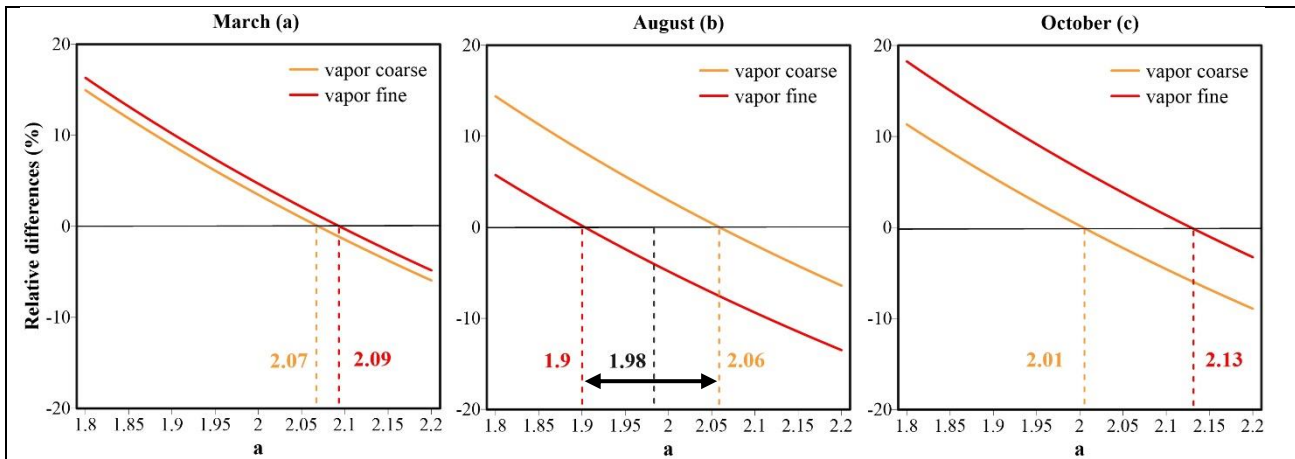
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The average diurnal cycles of  $\Delta e$  above the evaporation surface showed strong correlations with surface temperature, relative humidity, and air temperature. As shown in Figure 7C, the variations of  $\Delta e$  in spring and autumn were moderate compared to those in summer. Over sub-daily scale  $\Delta e$  for saturated sands showed a more pronounced amplitude compared to that for water, and it also peaked before that of water.  $\Delta e$  of Pan water showed a more pronounced amplitude than that of deeper column water. The water column  $\Delta e$  peaked before that of Pan over all seasons, except for summer.

PE through vapor diffusion is also influenced by the boundary layer thickness  $\delta$ . Following Equation (8), the optimization on parameter  $a$  is important for the performance of the vapor transfer equation. Compared with water,  $a$  was supposed to have more uncertainty for soil texture. This is because vapor transport from wet porous across the viscous boundary layer will exhibit more complex dynamics compared with smooth free water surface. We changed a series of  $a$  to obtain better match between measured and calculated PE for saturated textures. Results are plotted in Figure 10.



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Figure 8 Influence of parameter  $a$  on relative differences between calculated and measured PE for saturated fine sand (red solid curve) and coarse sand (orange solid curve) in different seasons.

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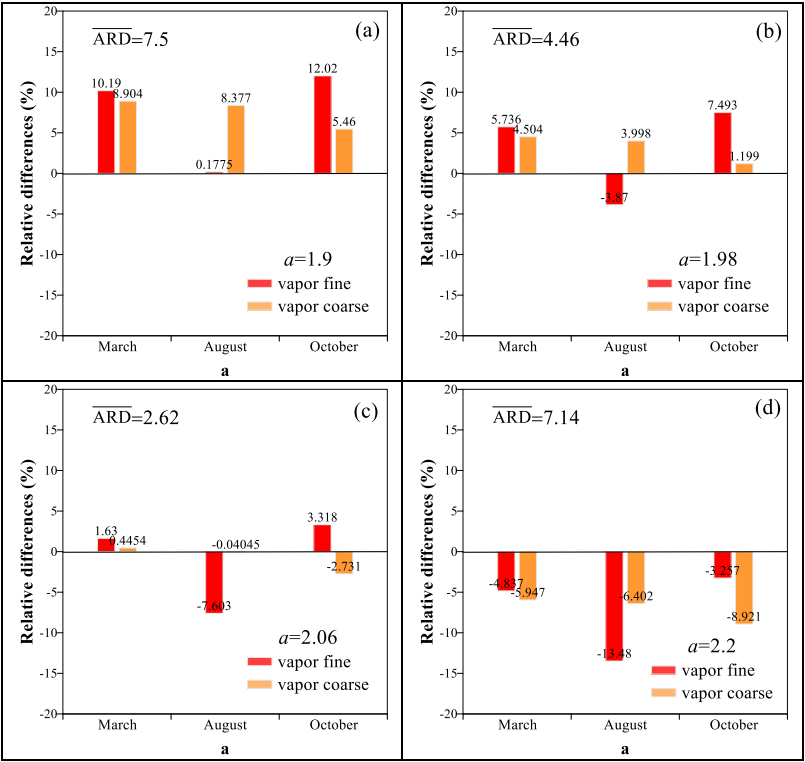
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The relatively larger contribution of PE during summer to the yearly PE demonstrate the importance of PE in hot period. As a result, we constitute a calibration procedure (by modifying  $a$ ) to close the measured evaporation rate in summer seasons. Relative differences on monthly PE approximate zero when  $a=1.9$  (red solid curve) and  $a=2.06$  (orange solid curve) were

394 employed for saturated fine sand and coarse sand respectively in August (Figure 8B). 1.9, 2.06,  
 395 their average 1.98 as well as 2.2 were selected to calculate PE according to the vapor transfer  
 396 equation. Relative differences between calculated and measured PE for each of the months and  
 397 for each of the soil textures were plotted in Figure 9. Average value of the absolute relative  
 398 differences  $\overline{ARD}$  (%) between calculated and measured monthly PE for soils were conducted to  
 399 evaluate the performance of  $a$ .



400 *Figure 9 Relative differences between calculated (according to the vapor transfer equation) and measured PE for*  
 401 *saturated fine sand and coarse sand in March, August and October.*  
 402 *(a)  $a = 1.9$ , (b)  $a = 1.98$ , (c)  $a = 2.06$ , (d)  $a = 2.2$*   
 403

404 Besides the total PE, the optimization on  $a$  also needs to consider its applicability in  
 405 estimating diurnal PE dynamics. The influence of  $a$  on PE dynamics were evaluated as well.  
 406 Measured and calculated average diurnal PE curves for saturated textures in different seasons  
 407 were plotted in Figure B1 and Figure B2 (Appendix B). Evaluation statistics results between  
 408 calculated and measured PE diurnal PE cycles are shown in Table B (Appendix B).

409 After a exhaustively optimization, we selected  $a=2.06$  for soils since it corresponds the best  
 410 to the measured total PE (with smallest  $\overline{ARD}$ , Figure 9b), as well as smaller deviations on

411 diurnal PE cycles (Table B,  $a=2.06$  provides the smallest MBE, and relatively smaller RMSE)  
412 from the measurements (Figure B1 and Figure B2).

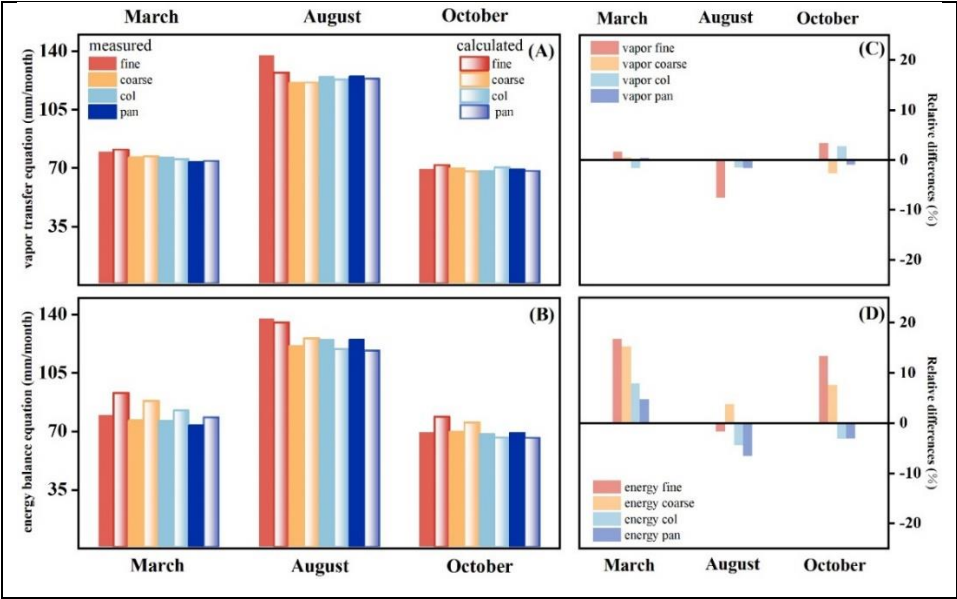
413 As for water in the column and water in Pan, the value of  $a$  were modified through trial and  
414 error to close the measured PE. We determined that  $a=5.8$  can be recommended for estimating  
415 water evaporation.

416

417 **3.6 Calculated PE**

418 **3.6.1 Total PE**

419 Figure 10 shows the measured and calculated monthly sums of PE for saturated fine sand,  
420 coarse sand, column water, and Class A Pan water according to the vapor transfer equation  
421 (Figure 10A) and the energy balance equation (Figure 10B). Differences between measured and  
422 calculated PE were plotted for each of the month and for each of the surface in Figure 10C and  
423 Figure 10D.



424 *Figure 10 Monthly sums of measured and calculated PE according to the vapor transfer equation (A), the energy*  
425 *balance equation (B). Relative differences between calculated and measured PE sums according to: (C) the vapor*  
426 *transfer equation, (D) the energy balance equation. Results are given for saturated fine sand, coarse sand, water,*  
427 *and water in Class A Pan.*  
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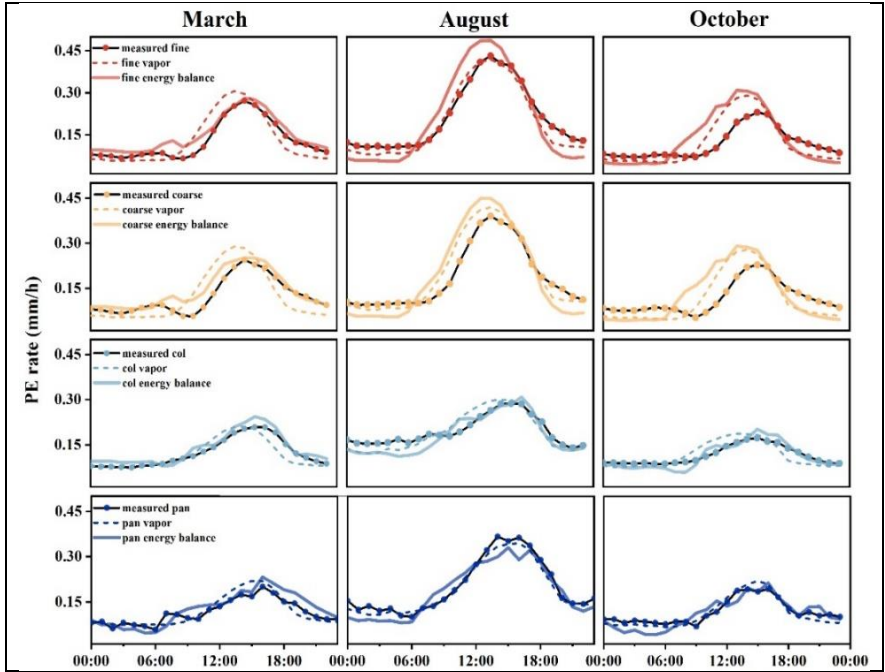
429 Figure 10C illustrates that, results by the vapor transfer equation showed minor relative  
430 differences between measurement and calculation. This method showed the closest match to

431 observed total  $PE_{col}$  and  $PE_{pan}$  in different seasons, but slightly underestimated  $PE_{fine}$  when the  
 432 evaporative demand was high (in August). Figure 10D indicates that the energy balance method  
 433 overestimated soil evaporations in most seasons but underestimated water evaporations in  
 434 August and October. Figure 10C and Figure 10D indicate that, the absolute values of relative  
 435 differences by the vapor transfer method were smaller than by the energy balance method.

436

### 437 3.6.2 Diurnal variations of PE

438 The comparison of measured and calculated seasonal-specific diurnal PE cycles over  
 439 different surfaces are plotted in Figure 11 to evaluate the applicability of the vapor transfer  
 440 equation and energy balance equation on sub-daily scale,



441 Figure 11 Measured and calculated average diurnal PE curves from saturated fine sand, coarse sand, water in  
 442 column and water in Class A Pan in different seasons.  
 443

444 A statistical evaluation of the fit between calculated and measured results (Figure 11;

445 Table ) indicates that. Both methods can describe the diurnal PE cycles, but they all resulted  
 446 in relatively larger errors in estimating  $PE_s$ . The vapor transfer method appears to provide a  
 447 reliable estimation of PE for both soil and water, particularly at night. With calculated results

deviate less from measured data (with smaller RMSE, and MBE in most conditions). However, this method predicts  $PE_{col}$ -peak too early in spring and autumn compared with measured results.

The energy balance method resulted in relatively larger deviations from measurements. This energy-unclosure seems to be caused by  $G$  (for soil) or  $N$  (for water). The measurements and calculations of  $G$  or  $N$  are generally subject to relatively larger uncertainties (Hamdani et al., 2018; Lensky et al., 2018). It should be noted that the energy balance method has advantages in estimating temporal PE dynamics by reproducing the time lag of the evaporation peak over soil and water in spring, as well as reproducing the evaporation peak over water surfaces.

**Table 1 Root mean square error (RMSE) and mean bias error (MBE) between calculated and measured evaporation rate (mm/h) for the saturated fine sand, coarse sand, water in column and water in Class A Pan, by the vapor transfer equation and energy balance method for season-specific average diurnal cycles**

Evaluation statistics	Method	March				August				October			
		fine	coarse	water	pan	fine	coarse	water	pan	fine	coarse	water	pan
RMSE	vapor	0.04	0.05	0.022	<b>0.021</b>	<b>0.027</b>	<b>0.033</b>	<b>0.026</b>	<b>0.033</b>	<b>0.048</b>	<b>0.005</b>	0.023	0.026
	energy	<b>0.026</b>	<b>0.027</b>	<b>0.017</b>	0.025	0.063	0.062	0.029	0.035	0.065	0.065	<b>0.017</b>	<b>0.024</b>
MBE	vapor	<b>0.002</b>	<b>0.001</b>	<b>-0.005</b>	<b>0.003</b>	-0.015	<b>0.001</b>	<b>-0.003</b>	<b>-0.01</b>	<b>0.005</b>	<b>-0.002</b>	<b>0.001</b>	<b>-0.004</b>
	energy	0.022	0.019	0.008	0.011	<b>-0.003</b>	0.007	-0.014	-0.017	0.016	0.01	-0.003	-0.005

## 4. Discussion

### 4.1 Energy transfer in evaporation process

#### 4.1.1 Energy transfer during daytime

The absolute value of  $H$  over a well-watered surface was the smallest among the flux terms. Energy considerations indicated that the differences in PE between over-saturated soils and water were mostly driven by available energy ( $R_n-G$  or  $R_n-N$ ).

Albedo determines incoming shortwave radiation. The albedo of soil is mainly influenced by water content and the color of the surface. In this research, the albedo for the coarse sand is slightly higher than fine sand due to the lower porosity and a slightly different color. Albedo of water in the lysimeter column was slightly smaller than those of saturated soil textures since radiation can penetrate deeper below the water column. In general, the higher albedo of water in shallow Class A Pan water could be attributed to some of the received radiation being reflected by the bottom of Pan. Most importantly, surface temperature ( $T$ ) will influence the outgoing

474 longwave radiation. Overall, as shown in Figure 6, the results showed the ranking of different  
475 surfaces according to daytime net radiative flux to be:  $R_{n, col} > R_{n, pan} > R_{n, fine} > R_{n, coarse}$ .

476 A smaller amount of absorbed  $R_n$  does not imply that PE for saturated soil is smaller. The  
477 energy re-distribution of the underlying material, which is determined by temperature variations  
478 and thermal properties, identified significant differences between saturated soils ( $G$ ) and water  
479 ( $N$ ).  $R_n$  cannot penetrate the soil during the day, and the delay in heat transfer through porous  
480 media with depth results in a smaller ground heat flux  $G$  (compared with  $N$ , Figure 6). Therefore,  
481 a significant portion of  $R_n$  was converted to evaporation over saturated soil surfaces. In particular,  
482 daytime energy available for evaporation over fine sand exceeded that for coarse sand as  
483 measurements showed that  $R_{n, fine} - G_{0, fine} > R_{n, coarse} - G_{0, coarse}$  during the day. In addition, the small  
484 peak in  $PE_s$  at around sunrise could be attributed to the temporal dynamics of available energy.  
485 The point at which  $G$  changed from negative to positive in spring and autumn lagged behind that  
486 of  $R_n$ . This delay increased the latent heat flux due to the large deficits between  $R_n - G$  (Li et al.,  
487 2022).

488 As for water, a significant portion of  $R_n$  was used to heat the deeper column water due to  
489 the water having the highest heat capacity. This resulted in less energy available for the latent  
490 heat flux ( $R_{n, w} - N < R_{n, s} - G$ ) and a delay in the release of stored energy to the latent heat flux and  
491 sensible heat flux in water. As a result, the  $PE_w$  curves lagged behind  $PE_s$ . This time lag was  
492 relatively shorter for  $PE_{pan}$  since shallow water bodies showed a reduced stored energy  $N$ .

493

#### 494 **4.1.2 Energy transfer at nighttime**

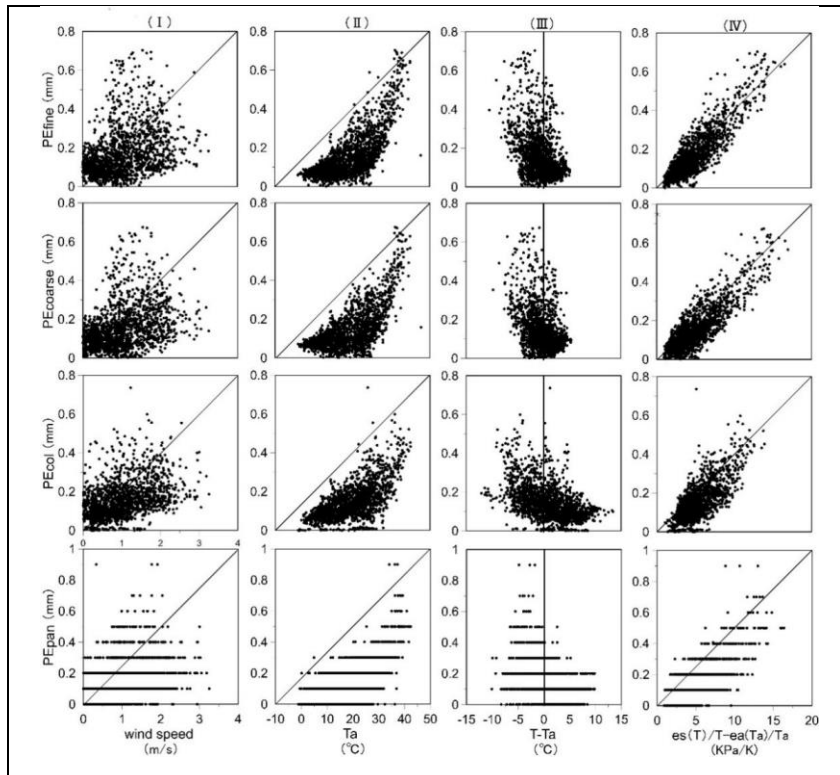
495 Except for water availability, nocturnal water loss achieved the highest positive correlation  
496 with wind speed, vapor pressure deficit, downward longwave radiation and a strong negative  
497 with relative humidity (Groh et al., 2019; K. E. Skaggs & S. Irmak, 2011; Maheu et al., 2014;  
498 Novick et al., 2009; Padrón et al., 2020). Light wind speed and high air relative humidity will  
499 inhibit evaporation as water vapor concentration gradient between soil (or water) and air will be  
500 too low to reach the latent heat status (Monteith, 1956; Jia et al., 2019). Under this climate  
501 condition, nighttime evaporation barely occurred and can be neglected.

502 Measured results imply that PE from wet surface is considerable and not constant at night.,  
 503 Given the lack of incoming shortwave radiation at night, nighttime  $G$  and  $N$  were negative due to  
 504 the release of heat by underlying material stored during the day. Released energy was mainly  
 505 transferred to the latent heat flux and upward longwave radiation emitted by the evaporation  
 506 surface.  $N$  has more negative value than  $G$  at night, as a result,  $PE_w$  is higher than  $PE_s$  in this  
 507 period (with  $R_{n,w}-N > R_{n,s}-G$ ). On the other hand, there were gradual increases in both  $G$  and  $N$   
 508 (from more negative to less negative) after sunset, leading to a decrease in PE from sunset to  
 509 dawn. Compared with measured results, nighttime PE is underestimated by the energy balance  
 510 method (Figure 11). This is because vapor diffusion process also an important contribution to  
 511 evaporation at night).

512

#### 513 4.2 Vapor transfer in evaporation process

514 Figure 12 illustrates the measured PE rates versus wind speed, air temperature, differences  
 515 between surface temperature and air temperature ( $T-T_a$ ), and vertical vapor gradients ( $\Delta e$ ).



516 Figure 12 Measured evaporation rates of saturated fine sand, coarse sand, and water vs. wind speed (I), air  
 517 temperature (II), differences between surface temperature and air temperature (III), and vertical vapor gradients  
 518 (IV).



519

520 The observations suggest that high wind speed and air temperature  $T_a$  enhance PE (Figure  
521 12I, Figure 12II). While wind speed has a more erratic pattern than  $T_a$ .

522 The deficits between surface temperature and air temperature ( $T-T_a$ ) indicating thermal  
523 stability of the air above the evaporation surface ( $T-T_a < 0$  represents a stable condition)  
524 (Hamdani et al., 2018).  $T-T_a > 0$  will result in increased mixing convection, and this unstable  
525 thermal condition generally enhances the evaporation rate under most conditions due to vapor  
526 transfer. However, the results shown in Figure 10III imply that PE was not directly proportional  
527 to thermal instability. The PE peaks for both water and saturated textures occurred mostly under  
528 stable conditions (when  $T-T_a < 0$ ).

529 Fick's first law of diffusion stipulates that vapor transfer during the evaporation process  
530 follows the "*vapor gradients* ( $\Delta e$ )" pattern (Hamdani et al., 2018). The measured PE cycles  
531 mostly followed those of  $\Delta e$ . While peaks in PE calculated by Equation (7) were not always  
532 aligned with measured peaks (Figure 11, Figure 12) since the vapor diffusion equation does not  
533 consider the complex energy transfer process below the evaporation surface.

534 The good performance of the vapor transfer equation achieved in the present study was  
535 highly dependent on the selection of suitable values of parameters  $a$  for saturated soil and water  
536 [in Equation (8)]. The present study ensured the same environmental conditions across the  
537 saturated sands and water. Lysimeter columns and Class A Pan are cylindrical with same area.  
538 Calibration results showed that the parameter  $a$  differed between saturated soil and water. This  
539 difference could be attributed to the relationships of  $\delta$  with airflow as well as with the  
540 characteristic length of the evaporating surface according to boundary layer theory (Haghighi &  
541 Or, 2013; Lim et al., 2012). In other words, the thickness of the boundary layer  $\delta$  in the vapor  
542 diffusion equation should be carefully considered instead of defined  $\delta$  only as a function of  
543 stream velocity.

544

#### 545 **4.3 Evaluation of PE estimation methods: the controlling factors**

546 The results of the present study indicated that the energy balance method showed  
547 advantages for estimating PE dynamics. This method is recommended for applications which  
548 require a high temporal resolution of PE dynamics, such as for precision agriculture. Given  
549 energy considerations and the results of the present study, we concluded that differences in PE  
550 and temporal dynamics between saturated bare soils and water are mostly driven by albedo and  
551 thermal properties (which influence  $G$  and  $N$  by influencing the surface temperature and  
552 temperature variations along depth). The results of the present study suggest that  $G$  and  $N$  are  
553 important factors explaining differences in total PE and diurnal dynamics in PE between over-  
554 saturated soil and water. Therefore, these factors should be comprehensively considered instead  
555 of ignored or simplified.

556 Moreover, the vapor diffusion equation can provide reliable estimates of nighttime PE and  
557 this method is recommended for estimating evaporation from the water surface. Compared with  
558 the energy balance equation, the vapor transfer equation provides a simpler approach. It requires  
559 common meteorological variables, surface temperature (which determines  $\Delta e$ ) as well as  
560 calibration of the boundary layer thickness  $\delta$  against historical evaporation data. In this research,  
561 choosing different  $a$  for saturated soils and water constitutes a simple calibration procedure to  
562 close the measured evaporation rate. A more comprehensive analysis of the influence of seasonal  
563 dynamics on  $a$  was beyond the scope of this paper and will be considered in future works.

564 Finally, previous research considered surface temperature  $T$  to be an independent external  
565 forcing that determine “wet surface” evaporation (Tu et al., 2021; Yang & Roderick, 2019). The  
566 results of the present study suggest that  $T$  is a significant internal forcing for both energy and  
567 vapor transfer during the evaporation process since it influences the redistribution of energy  
568 below the evaporation surface ( $G$  and  $N$ ), outgoing longwave radiation ( $R_{lu}$ ), as well as vapor  
569 gradients above the surface ( $\Delta e$ ).

570

## 571 **5. Conclusions**

572 The present study conducted high-resolution measurements of evaporation over saturated  
573 fine sand ( $PE_{\text{fine}}$ ), coarse sand ( $PE_{\text{coarse}}$ ), and water ( $PE_{\text{col}}$ ) using Markov-Bottle based lysimeters

574 and Class A Pan evaporation ( $PE_{pan}$ ) over more than a year. Measured results were  
575 comprehensively analyzed using the energy balance method and the improved vapor transfer  
576 equation. The following conclusions were derived from the results:

577 1. There were significant differences in evaporation among the different saturated bare soils  
578 and water  $PE_s$  is ~12 % higher than  $PE_w$  on a yearly scale. There were differences in both  
579 absolute values and temporal dynamics of diurnal PE curves between saturated soils and water.  
580 Annual  $PE_{fine}$  exceeded  $PE_{coarse}$  by 7.3%, with the differences more obvious during daytime in  
581 spring and summer. The cumulative evaporation over water column and Class A Pan showed  
582 minor differences. Compared with  $PE_{col}$ -curve, the  $PE_{pan}$ -curve resembles more the  $PE_s$ -curves  
583 over sub-daily scale.

584 2. Based on energy considerations, we conclude that differences in PE and PE temporal  
585 dynamics between saturated bare soils and water were mostly driven by available energy  $R_n$ - $G$   
586 or  $N$ . The vapor transfer processes over saturated soil and water were governed by different  
587 vapor pressure gradients above the surface as well as different boundary layer thicknesses.

588 3. The energy balance method showed advantages for estimating temporal dynamics in PE by  
589 reproducing the time lag of the evaporation peak. Energy transfer over different surfaces is  
590 controlled by deficits between albedo (which influences  $R_{ns}$ ) and surface thermal properties  
591 (which influences  $R_{nl}$ ,  $G$ , and  $N$  by influencing the variations in surface temperature and soil  
592 temperature along depth). This research proposes the precisely determination of  $G$  and  $N$  over  
593 different evaporation surfaces instead of ignoring these processes or their simplified  
594 representation as a fixed fraction. The vapor transfer equation is simpler and can reliably  
595 estimate PE. The vapor transfer processes over different surfaces were a function of surface  
596 temperature (which determines vapor gradients  $\Delta e$ ) and boundary layer thickness  $\delta$ .

597

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605

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788 **Appendix A**

789 In Equation (2),  $\chi$  is the surface albedo. The value of  $\chi$  approximates constant when solar  
 790 zenith angel  $\varphi$  is greater than  $40^\circ$ .  $\chi$  at noon can be selected to calculate  $R_{ns}$ . In this study area,  $\varphi$   
 791 at local solar noon is greater than  $40^\circ$  in most days of the year, and  $\varphi$  at solar noon is smaller  
 792 than  $40^\circ$  in late autumn and winter. In this research, average local solar noon albedo (measured  
 793 by CNR4 Net Radiometers Kipp&Zonen, Netherlands) for different saturated soil textures over  
 794 March, August and October were selected, results are shown in Table .

795 **Table A Local solar noon albedo for saturated soil textures**

Sand Type	$\varphi > 40^\circ$	$\varphi < 40^\circ$
fine	0.100	0.101
coarse	0.105	0.104

796 where 0.1 and 0.14 are the typical albedo for water and Class A pan (Rotstayn et al., 2006;  
 797 H. Yang & Yang, 2012).  $\varepsilon$  is the surface emissivity, representing the reflection coefficient of  
 798 longwave radiation at the surface. For water,  $\varepsilon=0.97$  (Gianniou & Antonopoulos, 2007). For soil,  
 799  $\varepsilon$  can be determined by water content on top surface (Jensen et al., 2016):

$$800 \quad \varepsilon = \min(0.9 + 0.18 \cdot \theta_{top}; 1.0) \quad (A1)$$

801 In Equation (5), saturated soil volumetric heat capacity calculated using the following  
 802 equation:

$$803 \quad c_s = c_{vw} \cdot \theta_s + c_{si} \cdot (1 - \theta_s) \quad (A2)$$

804 where  $c_{vw}$  (4.19MJ/m<sup>3</sup>K) is the volumetric heat capacity of water, and  $\theta_s$  is the saturated  
 805 water content for the specific soil texture (0.39 for fine sand and 0.37 for coarse sand).  $c_{si}$   
 806 (MJ/m<sup>3</sup>K) is the dry soil volumetric heat capacity (0.95 for fine sand and 0.92 for coarse sand)  
 807 measured by TPS 1500 (Hot Disk, Sweden).

808 In Equation (6), aerodynamic resistance for saturated soil can be calculated as:

$$809 \quad r_{a,s} = \frac{\left( \ln \frac{z_1 - d}{z_{om}} \right) \left( \ln \frac{z_2 - d}{z_{ov}} \right)}{k^2 \cdot u} \quad (A3)$$

810 where  $k$  ( $=0.4$ ) is the von Karman constant,  $u$  ( $\text{m}\cdot\text{s}^{-1}$ ) is the wind speed,  $z_1$  (m) and  $z_2$  (m) are the  
811 measurement heights of wind speed (2m) and air temperature (1.5m),  $d$  ( $=0$ ) is the displacement  
812 height,  $z_{om}$  and  $z_{ov}$  are surface roughness lengths for momentum flux and heat flux, with values  
813 equal to 0.001 for bare soils (Simunek et al., 2005). The aerodynamic resistance for the water  
814 surface can be calculated as:

$$815 \quad r_{a,w} = \frac{4.72 \left( \ln \left( \frac{z_1}{z_0} \right) \right)^2}{1 + 0.536 \cdot u} \quad (A4)$$

816 where  $z_0$  with value equal to 0.00137 is the aerodynamic roughness factor for water surface  
817 (Shuttleworth, 1993).

818 In Equation (7),  $D_v$  ( $\text{m}^2\cdot\text{s}^{-1}$ ) is the diffusion coefficient for water vapor in air (Ben Neriah et  
819 al., 2014):

$$820 \quad D_v = 2.11 \cdot \left( \frac{T_a + 273.14}{273.14} \right) \cdot 10^{-5} \quad (A5)$$

821 In Equation (A1),  $e_s$  is the vapor pressure at surface (Pa), it is assumed to be equivalent to  
822 saturated vapor pressure at surface temperature  $T$ . Which can be calculated by (Allen et al., 1998;  
823 Zotarelli et al., 2010) :

$$824 \quad e_s(T) = 610.8 \exp\left(\frac{17.27 \cdot T}{T + 237.14}\right) \quad (A6)$$

825 In Equation (A1), hourly actual vapor pressure  $e_a$  (Pa) is given by:

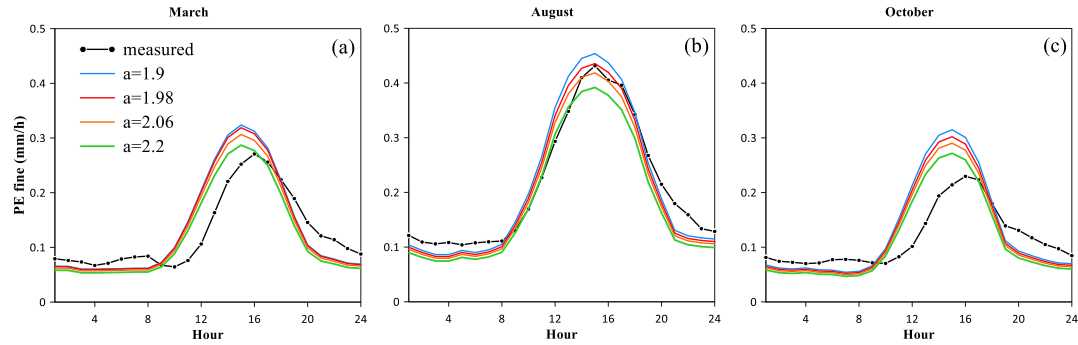
$$826 \quad e_a(T_a) = e_s(T_a) \cdot RH \quad (A7)$$

827  $RH$  is the relative humidity from meteorological station.

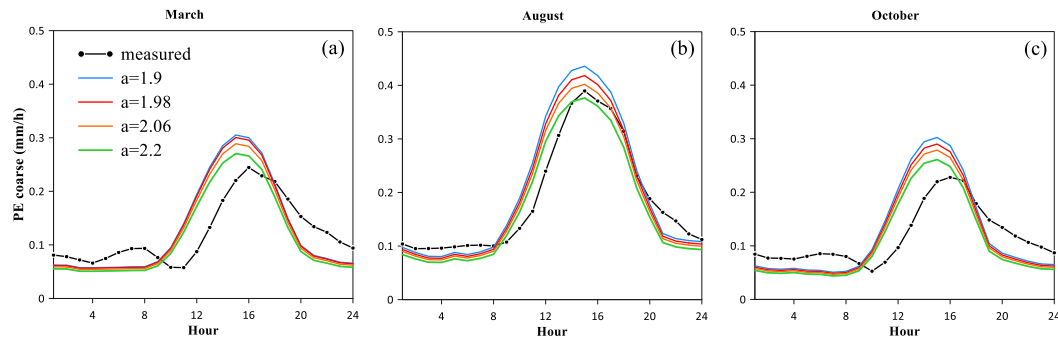
828

829 **Appendix B**

830 Measured and calculated average diurnal PE curves (with different  $a$ ) for saturated soil  
831 textures in different seasons were plotted in Figure B1 and Figure B2.



832  
833 *Figure B1 Measured and calculated average diurnal PE curves (by the vapor transfer method, with  $a=1.9$  in blue,*  
834  *$a=1.98$  in red,  $a=2.06$  in orange,  $a=2.2$  in green) for saturated fine sand in March (a), August (b) and October (c).*



835  
836 *Figure B2 Measured and calculated average diurnal PE curves (by the vapor transfer method, with  $a=1.9$  in blue,*  
837  *$a=1.98$  in red,  $a=2.06$  in orange,  $a=2.2$  in green) for saturated coarse sand in March (a), August (b) and October*  
838 *(c).*

839 Evaluation statistics between calculated and measured PE for the average diurnal cycles are  
840 shown in Table B.

841 **Table B Root mean square error (RMSE) and mean bias error (MBE) between calculated and measured**  
842 **evaporation rate (mm/h) for the saturated fine sand and coarse sand, by the vapor transfer equation for**  
843 **season-specific average diurnal cycles**

Evaluation statistics	Soil textures	Average Diurnal PE cycles (mm/h)											
		March				August				October			
		1.9	1.98	2.06	2.2	1.9	1.98	2.06	2.2	1.9	1.98	2.06	2.2
RMSE	fine	0.045	0.043	0.04	<b>0.037</b>	0.029	0.027	<b>0.027</b>	0.034	0.056	0.052	0.048	<b>0.042</b>
	coarse	0.054	0.053	0.05	<b>0.047</b>	0.043	0.037	0.033	<b>0.031</b>	0.056	0.053	0.05	<b>0.045</b>
MBE	fine	0.009	0.007	<b>0.002</b>	-0.006	<b>0.001</b>	-0.007	<b>-0.015</b>	-0.028	0.015	0.01	0.005	<b>-0.003</b>
	coarse	0.008	0.006	<b>0.001</b>	0.007	0.016	0.008	<b>0.001</b>	-0.011	0.008	0.003	<b>-0.002</b>	-0.009