Impurities in inhomogeneous superconductors from Density Functional Theory (DFT)



David Antognini Silva*, Philipp Rüßmann, and Stefan Blügel

Peter Grünberg Institute and Institute for Advanced Simulation, Forschungszentrum Jülich and JARA, D-52425 Jülich, Germanys

* Corresponding author email: d.antognini.silva@fz-juelich.de

Introduction

- The juKKR code [1] is a DFT-based calculation method using Green's functions through the Korringa-Kohn-Rostoker (KKR) method
- The KKR method shows interesting features, such as:
 - single site and multiple scattering problem separation
- dealing with reduced symmetry system / impurities
- To perform calculations that include superconductivity behaviours, the juKKR code is being extended with the Bogoliubov-de Gennes (BdG) formalism [2]

At the moment, the BdG formalism has been successfully implemented for homogeneous superconductor calculations [3].

Work in progress: Implementing the BdG formalism for inhomogeneous superconductors. This will allow to investigate behaviours of exotic materials, such as Majorana qubits or Yu-Shiba-Rusinov chains.

The KKR method

In the KKR method, we define the Green's function as the resolvent of the Hamiltonian, which can be formulated as [4]:

$$G(\mathbf{r} + \mathbf{R}^{n}, \mathbf{r}' + \mathbf{R}^{n'}; E)$$

$$= \delta_{nn'} \sqrt{E} \sum_{L} H_{L}^{n}(\mathbf{r}_{>}; E) R_{L}^{n}(\mathbf{r}_{<}; E) + \sum_{LL'} R_{L}^{n}(\mathbf{r}; E) G_{LL'}^{nn'}(E) R_{L'}^{n'}(\mathbf{r}'; E)$$
single site

multiple scattering

The electronic density is obtained from the imaginary part of the Green's function:

$$\rho(\mathbf{r}) = -\frac{1}{\pi} Im \sum_{\sigma} \int_{-\infty}^{E_F} dE \ G^{\sigma,\sigma}(\mathbf{r},\mathbf{r};E)$$

Impurities within the KKR-Green function method

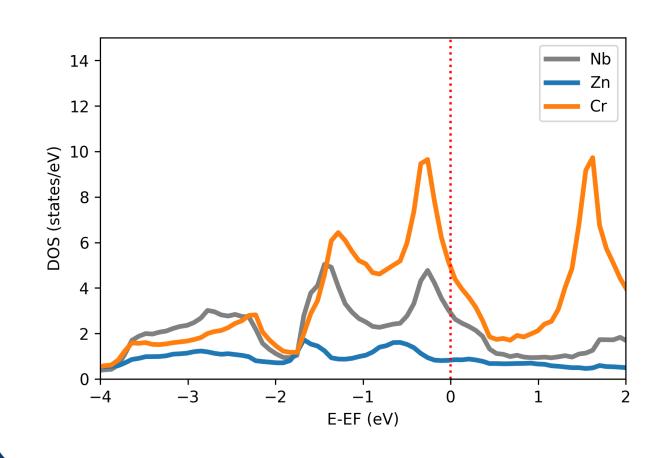
The Green's function \mathcal{G} of a perturbed system which Hamiltonian is $\mathcal{H} = \mathcal{H}_0 + \Delta \mathcal{V}$ can be obtained from the Green's function \mathcal{G}_0 of the unperturbed system through the Dyson equation [4]:

$$\mathcal{G} = \mathcal{G}_0 + \mathcal{G}_0 \Delta \mathcal{V} \mathcal{G} = \mathcal{G}_0 + \mathcal{G}_0 \mathcal{T} \mathcal{G}_0$$

The problem of calculating the Green function of the perturbed system is then translated into finding the T-matrix of the system:

$$\mathcal{T}(E) = \Delta \mathcal{V}(1 + \mathcal{G}_0(E)\mathcal{T}(E))$$

Effect of impurities on the DOS of bcc Nb



This graph shows how the DOS of homogeneous bcc Nb (grey) is affected by a Zn/Cr impurity (blue/orange).

We can see that the population at the Fermi level (red dotted-line) can either significantly increase or decrease, depending on the impurity.

The Bogoliubov de Gennes Method

Superconductivity can be described using the Bogoliubov-de Gennes (BdG) equations. These can be formulated in the effective single particle picture of density functional theory [2]:

$$H_{\text{BdG}} = \begin{pmatrix} -\nabla^2 - E_{\text{F}} + V^{\text{eff}}(\vec{x}) & \Delta^{\text{eff}}(\vec{x}) \\ \Delta^{\text{eff}*}(\vec{x}) & \nabla^2 + E_{\text{F}} - V^{\text{eff}*}(\vec{x}) \end{pmatrix}$$

$$\Psi_n(\vec{x}) = \begin{pmatrix} u_n(\vec{x}) \\ v_n(\vec{x}) \end{pmatrix} \quad \text{electron like} \quad \text{hole like}$$

Diagonalizing the BdG Hamiltonian gives the *normal* and *anomalous* densities:

$$\rho(\vec{x}) = \sum_{n} |u_n|^2(\vec{x}) f(\varepsilon_n) + |v_n|^2(\vec{x}) [1 - f(\varepsilon_n)]$$

$$\chi(\vec{x}) = \sum_{n} u_n(\vec{x}) v_n^*(\vec{x}) [1 - 2f(\varepsilon_n)]$$

This determine the effective single particle potentials

$$V^{\rm eff}(\vec{x}) = V^{\rm ext}(\vec{x}) + \int \frac{\rho(\vec{x'})}{|\vec{x} - \vec{x'}|} \mathrm{d}^3x' + \underbrace{\delta E^0_{\rm xc}[\rho]}_{\delta\rho(\vec{x})} \underbrace{\delta\rho(\vec{x})}_{\text{consistency}}$$

$$\Delta^{\rm eff}(\vec{x}) = \lambda \chi(\vec{x})$$

$$LDA, GGA, \dots$$

The coupling matrix determines the form of the pairing symmetry (e.g. s-wave) and assumes local pairing

The Coupling constant λ

In order to make accurate BdG calculations, experimental data is needed to determine the value of the coupling constant λ . This can be achieved in two ways:

1) By measuring the superconducting gap size

Variations of the coupling constant λ affect the size of the superconducting gap (cf. graphs). If the gap size is experimentally known, the value of λ is chosen such as the calculated superconducting gap size matches the measured one.

2) By measuring the critical temperature

The critical temperature T_c is linked to the coupling constant through [5]:

$$k_BT_C = 1.14\omega_D e^{-1/N_F\lambda} \qquad \omega_D = \text{Debye frequency of pnonons} \\ N_F = \text{DOS at Fermi level}$$

References

- The JuKKR code package, https://jukkr.fz-juelich.de
- G. Csire, B. Újfalussy, J. Cserti, and B. Győrffy, Phys. Rev. B 91, 165142 (2015)
- P. Rüßmann and S. Blügel, in preparation (2021)
- N. Papanikolaou, R. Zeller and P.H. Dederichs, J. Phys.: Condens. Matter 14, 2799-2823 (2002)
- J. Bardeen, L.N. Cooper, J.R. Schrieffer, Phys. Rev. **108**(5), 1175-1204 (1957)

Acknowledgements

 ω_D = Debye frequency of phonons

This work was funded by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) under Germany's Excellence Strategy - Cluster of Excellence Matter and Light for Quantum Computing (ML4Q) EXC 2004/1 - 390534769.