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ABSTRACT

We report an experimental design where the position and resonance frequency of the nitrogen vacancy (NV) in a diamond are correlated with the room temperature. A simple model trained on the interpolated correlation data predicts both quantities. The predictive tracking of the NV's location enables continuous operation of the NV quantum computer under ambient conditions for a week without recalibration.

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Over the last two decades, quantum computing has evolved from single-qubit demonstrators to systems with up to hundreds of qubits. There is intense competition between technologies and research groups have started acquiring small-scale hardware. Some of these competing technologies are cold atoms,¹ superconducting qubits,² trapped ions,³ quantum dots,^{4,5} defects in solids,⁶ and photonic systems.⁷

We focus our attention on nitrogen vacancy (NV) centers. Single defects in diamonds provide a promising platform for quantum computations.^{8–10} Electron and nuclear spins in an NV count as qubits fulfilling different roles. While electron spins offer fast control and high fidelity readout,^{11–14} nuclear spins provide additional qubits with long coherence times that can be used to store and process quantum states.^{15–19} Electron qubits play a central role because they can be used to control and connect nuclear spins as well as provide the potential to scale this technology through electron–electron and electron–photon couplings.^{19–23} The latter can help realize long-range entanglement.^{24–26}

Owing to these advantages, NVs have become a reasonable candidate for quantum computations when using optical detection of the magnetic resonance. The sensitivity of NV center's photoluminescence to temperature has been well studied.^{27,28} When operated below 10 K, spin-dependent optical transitions²⁹ have demonstrated multipartite entangled states with up to seven qubits.³⁰ High-fidelity two-qubit gates^{31–37} as well as basic quantum algorithms^{32,38} have been demonstrated. However, the higher cost and specialization required for

cryogenic equipment is a deterrent for those interested in NV quantum computing lacking a strong technical expertise.

Fortunately, the NV can also easily operate at room temperature, making it very cost competitive for devices with a few qubits. At room temperature, researchers have employed the NV's spin as a sensor for magnetic^{39,40} and electric fields,⁴¹ and thermometry.^{42,43} Operation at room temperature with typical fluctuations of environmental variables like temperature and humidity causes a thermal dislocation of the focus of the optical instruments and the NV, resulting in photoluminescence loss and computational errors. Periodically determining the NV's position within the diamond with high accuracy is therefore critical. Periodic relocation scans in three dimensions in turn interrupt quantum computations and reduce computational time.

The problem has conventionally been approached by using a large tracking window (e.g., $5 \times 5 \mu\text{m}^2$ parallel to the surface⁴⁴) but this leads to either less accurate position estimates or increased relocation time, without effectively solving the problem. Alternatively, multi-point methods can track with significantly fewer data points. The use of 8– points,^{45,46} and even 6–, 4–, 3– point methods has been demonstrated.^{42,47,48} In nanodiamond thermometry, increasing the number of points is expected to improve tracking fidelity but reduces the number of measurements, leading to reduced precision.⁴⁷ In quantum computing, any time spent on tracking the NV is lost for computation. While the multi-point methods aim to reduce the number of points scanned, they still rely on some form of active relocation measurements,

interrupting any form of computation. Overall, despite a reduction in the track-and-locate time using multi-points, these methods essentially remain in the paradigm of stop-and-fix, where the computation is stopped to fix the tracked coordinates.

In this work, we determine the NV's position without requiring active scanning, allowing for continuous operation. We report on the mechanism of placing the optical equipment in thermal isolation, leading to the correlation of the NV's position and resonance frequency with the room's temperature sensors. By training a simple model from the experimental data, we are able to accurately predict the positions and resonance frequency ν_{res} using external temperature measurements as inputs. As we show below, this design and implementation is a simple yet highly effective way to reduce the time spent to periodically track the electron within the NV at room temperature inside a tight window ($1 \times 0.7 \mu\text{m}^2$) parallel to the surface. Our work moves the recalibration into a different paradigm: a simple experimental redesign that enables modeling the drift and predicting it in real-time. Our approach shifts the paradigm from stop-and-fix to predict-and-prevent. The location of the NV is predictively tracked and the necessity for scanning is prevented for prolonged periods of time.

The experimental setup is designed such that the optical table is enclosed in a box of dimensions $45 \times 50 \times 15 \text{ cm}^3$ ($L \times W \times H$). Figure 1 shows the setup. The optics table is located in a small room of dimensions around $2 \times 2 \times 3 \text{ m}^3$ that has a weak thermal isolation, i.e., no use of heat exchange equipment. The optics table is liquid cooled and kept at a constant target temperature. A sensor is placed on the optics table to measure the temperature T_1 . The major sources of heat generation, e.g., pulse generators, are kept outside in a bigger room which is air cooled. T_2 is the room temperature measured through a single sensor mounted in the room and varies uncontrolled with the temperature outside. Both T_1 and T_2 are not to be confused with spin relaxation or coherence times and are measured every 20 s.

Our NV of interest is located within a few micrometers of the surface of an ultrapure diamond sample. The electron spin has a gyromagnetic ratio of 28.024 GHz/T and has a zero field splitting at ≈ 2.869 GHz. The external magnetic field is fixed at ≈ 0.05 T. We track the NV in a window of dimensions 1, 0.7, and $3 \mu\text{m}$ in the

X , Y , and Z orientations, respectively, every 400 s using the software Qudi.⁴⁹

We continuously track the change in location of the NV with respect to the temperature over a period of 10 days. The lowest and highest temperatures measured were 13.08 and 13.31 °C for T_1 , and 15.8 and 20.0 °C for T_2 , respectively. It is worth noting that T_1 and T_2 have different least counts, namely, 0.01 and 0.02 °C, respectively. The point of origin of the NV's coordinates is chosen arbitrarily. Data collection was interrupted for 10 min on day three by an unexpected system restart. This caused an unusually large change in Y and Z , attributed to an unknown software glitch during the abnormal shutdown. For data normalization, we added -0.227 and $0.05 \mu\text{m}$ to all Y and Z values after the third day, respectively, such that they had identical positions in the last pre- and first post-interruption measurements.

We plot the change in NV's location with respect to T_2 in Fig. 2(a). The left y-axis gives the positions in $-X$, Y , and $-Z$ coordinates, respectively, in units of 10^{-6} m . The values of X and Z are negated in order to better visualize their anti-correlation with respect to T_2 . We perform a linear interpolation of all the four variables since the temperature data and NV's location are not collected at exactly the same time but often with a difference of less than few tens of seconds. Figure 2(a) shows only a finite sample of these interpolated data points for visualization.

To better understand the relationship between the NV's location and temperature, we linearly interpolate the data points to calculate any correlation between the variables. First, we normalize the data of all five variables such that it ranges between 0 and 1. Second, we rearrange T_2 values in the decreasing order and sort all the other variables according to this order. Doing so allows us to retain the same timestamped values of others variables with reference to T_2 . Third, we perform a linear fit to all the variables to compute the respective slopes. Fourth, we use $1 - \frac{2|a| - |b|}{(|a| + |b|)}$ or $\frac{2|a| - |b|}{(|a| + |b|)} - 1$ when slopes a and b have the same signs or opposite signs, respectively, to compute the correlations. Finally, we calculate these correlations for all pairs of variables and show the results in Fig. 2(b).

Data that show the same trend will have the same slope and hence will be perfectly correlated through our formula. From Fig. 2(b), we observe that the positions X and Z (Y) are highly anti-correlated (correlated) with the temperature T_2 , and only weakly with T_1 . The weak correlation is likely an artifact of the small temperature change (0.23 °C) subjected to a limited dynamic range because of the sensor's least count (0.02 °C). Figure 2(b) quantifies the visual relations shown in Fig. 2(a). Furthermore, the positions X , Y , and Z are strongly (anti-)correlated among themselves, as is to be expected.

We now investigate the following question: To what extent are the Rabi oscillations of the electron sensitive to the driving frequency? We calculate the *Rabi contrast* as two times the amplitude of a sine fitting to the oscillations divided by the offset. We run a scan between 1.456 and 1.460 GHz and plot the percentage Rabi contrast in Fig. 2(c). The data points fit very well to a Lorentzian curve and the full width at half maxima (FWHM) is 1.55 MHz. This measure gives us a quantitative loss in Rabi contrast (detuning error) one is willing to tolerate for off resonance frequencies.

If the NV's location changes due to a temperature change, what is the change in the resonance frequency? We continuously measure the resonance frequency for a large portion of the 10 days with only two interruptions. Each measurement takes ≈ 191 s. The results are shown

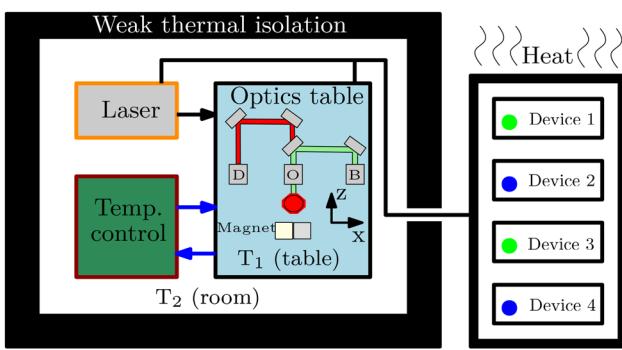


FIG. 1. Experimental optical setup of our NV center based quantum computer. The optics table, laser, and liquid cooling temperature control are placed within a weakly thermally isolated room. The optics table consists of a diamond (red) placed close to a magnet. A green beam of light originating at **B** is directed onto the diamond through the objective **O**. Emitted red light is detected at detector **D**. Control equipment consisting of several devices that generate most of the heat is kept outside.

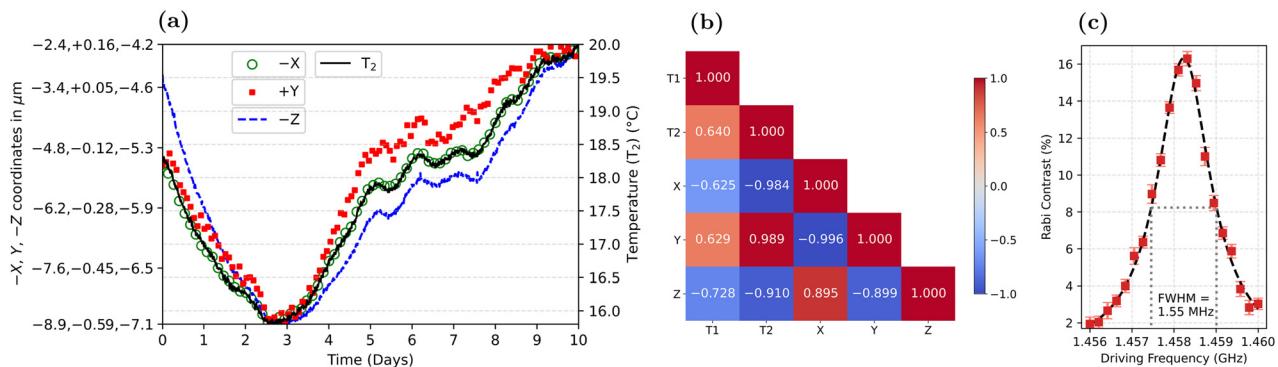


FIG. 2. (a) The relative change in NV's location (left axis) and temperature T_2 (right axis) plotted over a time period of 10 days. The values in X and Z directions are negated to better visualize their anti-correlation. Strong linear (anti-)correlations are evident despite the different scales of X , Y , and Z . The axes are not optimized for superimposition. (b) The correlation matrix between the NV's coordinates X , Y , and Z , and temperatures T_1 and T_2 shows the strength of their linear correlations (red) and anti-correlations (blue). The apparent stronger (anti-)correlation of positions with T_2 compared to T_1 is due to the measurement sensitivities of the sensors. As is also visible in (a), the measurables $-X$, Y , and T_2 have near perfect correlation. (c) The percentage difference in the crest and trough of the Rabi oscillations obtained for different driving frequencies at constant temperature is plotted in red square points. A Lorentzian fit to the data is shown in dashed black lines. The full width at half maxima (FWHM) is 1.55 MHz.

in Fig. 3. The measured resonance frequencies (in GHz) along with their standard deviations and temperature T_2 are plotted against time. We visually observe a strong first-order linear correlation between T_2 and the resonance frequency.

We use second order polynomial regression to predict the position coordinates of the NV and the resonance frequency ν_{res} . We train our models on the data T_1 and T_2 suitably linearly interpolated across the first two and half days so that the data points have the same time-stamps. We use the model

$$\nu_{\text{res}} = \beta_{0,\gamma} + \beta_{1,\gamma} T_1 + \beta_{2,\gamma} T_2 + \beta_{3,\gamma} T_1^2 + \beta_{4,\gamma} T_1 T_2 + \beta_{5,\gamma} T_2^2,$$

where we estimate the coefficients $\beta_{i,\gamma}$ by minimizing least squares for each $\gamma \in \{X, Y, Z, \nu_{\text{res}}\}$ separately. The choice of the model is motivated by the results shown in Fig. 2(a) where some non-linear contributions are apparent to mostly linear correlations. With ordinary least squares, these parameters are optimized to best fit the interpolated observed data and capture both linear, using $\beta_{0,\gamma}, \beta_{1,\gamma}, \beta_{2,\gamma}$, and non-linear, using $\beta_{3,\gamma}, \beta_{4,\gamma}, \beta_{5,\gamma}$, contributions. The goal of the modeling is to accurately and precisely predict ν_{res} and coordinates X , Y , and Z given T_1 and T_2 . The accuracy and precision of prediction for the

coordinates are related to the size of the tracking window in which the NV's position is scanned.

The results for the coordinate predictions are shown in Fig. 4(a). The plots show the absolute difference between the actual data and the prediction data of the model. The results show that our model's accuracy is different in each coordinate. The predictions become less accurate over time, which is to be expected. While the model is able to predict locations for coordinates X and Y rather accurately, it only qualitatively captures the movement of Z . The underlying cause requires further investigation. However, the prediction error in Z is still well within the size of the tracking window.

In Fig. 4(b), we show the prediction for ν_{res} along with the actual values. We observe that our model is able to predict the change in ν_{res} within about one FWHM [see Fig. 2(c)] of Rabi contrast for slightly more than a week when temperature is left uncontrolled. The cause of the frequency drift is most likely the change of distance between the magnet and the diamond due to the thermal expansion/contraction of the optics table on which they are mounted. We emphasize that for continued normal operations of an NV center based quantum computer, it is only required that ν_{res} is predicted accurately within a few hours, not days. Our simple model is sufficient for this purpose.

Albeit highly setup-specific, it is interesting to calculate the shift in NV's position per $^{\circ}\text{C}$ by dividing the maximum change in each coordinate by the maximum change in T_2 across the entire data. The position changes by $\Delta X = 1.533 \mu\text{m}/^{\circ}\text{C}$, $\Delta Y = 0.178 \mu\text{m}/^{\circ}\text{C}$, and $\Delta Z = 0.682 \mu\text{m}/^{\circ}\text{C}$. Similarly, $\Delta \nu_{\text{res}} = 517 \text{ kHz}/^{\circ}\text{C}$. Thus, for a temperature change of >0.33 , >1.98 , and >2.21 $^{\circ}\text{C}$, the brightest spot of the NV's point spread function (PSF) will be outside our tracking window in the X , Y , and Z directions, respectively. Similarly, a temperature change of 1.41 $^{\circ}\text{C}$ reduces the Rabi contrast by half. These values give us a reference for our experiments to which, if the external temperature T_2 is optimally controlled, and everything else remains unchanged, the NV will always be found within the tracking window and reachable at the same resonance frequency with a high contrast.

If T_2 varies uncontrolled, as in our experiment, the trained model allows one to predict the NV's movement without requiring periodic active position measurements. The NV's position can then be updated,

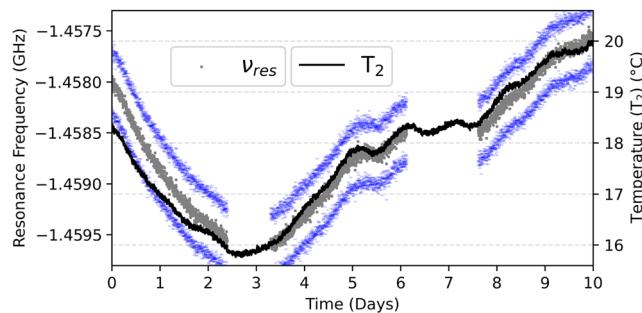


FIG. 3. The negative of measured resonance frequencies (gray points) in GHz (left axis) and temperature T_2 (right axis) in $^{\circ}\text{C}$ (black line) are plotted against time. Each blue point represents one standard deviation of the resonance frequencies.

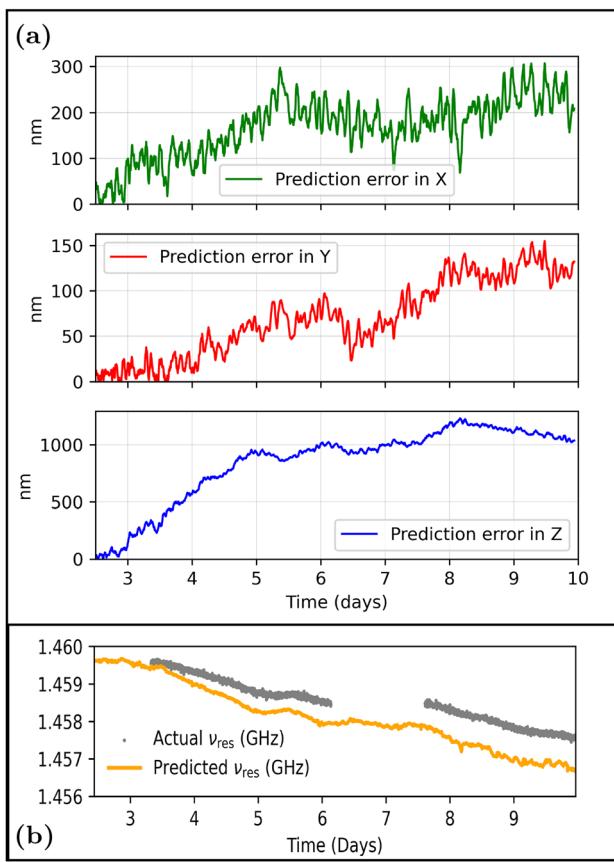


FIG. 4. (a) The absolute difference between actual and predicted X (top), Y (middle), and Z (bottom). The units of y-axes is nanometers. (b) The actual ν_{res} (gray) in GHz and the predicted ν_{res} (orange) as a function of time.

without active scanning, at regular intervals by the model's predictions based on T_1 and T_2 . This prevents the interruptions of quantum algorithms and allows for continuous operation of the quantum computer. While our model takes into account the temperature fluctuations and makes reasonable predictions for the position coordinates, temperature is not a unique factor influencing the NV's position. Factors like humidity, diamond strain, chiller flow rate, and table vibrations should be investigated and potentially added to the simple model for more precise predictions.

In addition to quantum computing, our method can be useful in nanodiamond thermometry where multi-points methods are currently used. If the nanodiamonds are immobilized, for instance by surface adsorption, polymer embedding, or covalent attachment to a substrate, their Brownian motion can be effectively suppressed, and our approach could in principle be applied in this scenario for real-time estimations of magnetic resonance shifts in biological applications.

We described a simple experimental setup where an NV center in a diamond is placed in weak thermal isolation and correlates very well with temperature sensors for measurements taken over 10 days. We demonstrated that the resonance frequency is also correlated with the room temperature. By using a simple quadratic model as a function of

the temperature of two sensors, training it for a period of two and half days, we were able to accurately and precisely predict the movement of the NV for the next one week. Similarly, our model was able to predict ν_{res} to a reasonable accuracy.

Our results can be improved in several ways. More temperature sensors with lower least counts and shorter logging intervals would help generate more accurate data for training our model, leading to more accurate predictions. Different sensors, e.g., to measure humidity and vibrations, could also be integrated in the model. The model itself could be made more complicated to include non-linear correlations more accurately.

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AUTHOR DECLARATIONS

Conflict of Interest

The authors have no conflicts to disclose.

Author Contributions

Manpreet Singh Jattana: Conceptualization (equal); Data curation (equal); Investigation (equal); Resources (equal); Visualization (equal); Writing – original draft (equal); Writing – review & editing (equal). **Thomas Lippert:** Conceptualization (equal); Investigation (equal); Resources (equal); Writing – review & editing (equal).

DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

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