

Anisotropic magnetoelastic coupling in single-crystalline CeFeAsO as seen via high-resolution x-ray diffraction

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Single-crystal synchrotron x-ray diffraction studies of CeFeAsO reveal strong anisotropy in the charge-correlation lengths along or perpendicular to the in-plane antiferromagnetic (AFM) wave vector at low temperatures, indicating an anisotropic two-dimensional magnetoelastic coupling. The high-resolution setup allows to distinctly monitor each of the twin domains by virtue of a finite misfit angle between them that follows the order parameter. In addition, we find that the in-plane correlations, above the orthorhombic (O)-to-tetragonal (T) transition, are shorter than those in each of the domains in the AFM phase, indicating a distribution of the in-plane lattice constants. This strongly suggests that the phase above the structural O-to-T transition is virtually T with strong O-T fluctuations that are probably induced by spin fluctuations.

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Understanding the strong magnetoelastic coupling observed in parent ferropnictides is pivotal to unraveling the mechanism that makes these systems superconducting (SC) by chemical doping.^{1–3} This strong coupling is manifested in the emergence of a stripe-type antiferromagnetic (AFM) phase that is stable in an orthorhombic (O) phase that results from shearing distortions of a high-temperature tetragonal (T) phase, as illustrated in Fig. 1. The magnetic and structural transitions occur simultaneously in the AFe_2As_2 ($\text{A} = \text{Ca}, \text{Sr}, \text{and Ba}$, “122”) systems,^{4–6} while the magnetic ordering transition temperature (T_N) is in a range of $\sim 6\text{--}18$ K below the structural transition temperature (T_S) and is strongly dependent on the sample quality in the LnFeAsO ($\text{Ln} = \text{lanthanide element}$, “111”) family.⁷ The spin-lattice coupling has also been implied in the interpretation of spin dynamics as well as its influence on the structures of these systems. Inelastic neutron-scattering studies from single-crystal CaFe_2As_2 (Refs. 8 and 9) and polycrystalline LaFeAsO (Ref. 10) showed that spin fluctuations persist above T_N up to at least room temperature. It was argued that such fluctuations introduce dynamic disorder of the O and T phases, so that finite orthorhombicity and tetragonality may exist above and below T_S , respectively.^{6,7,11} Furthermore, the magnetic and structural transitions can be simultaneously tuned by chemical substitutions, e.g., suppressing the AFM and O phases and setting in the SC state.¹²

The magnetic structure as schematically shown in Fig. 1 and spin dynamics indicate a strong magnetic frustration due to the competition between the nearest-neighbor (NN) and the next-NN (NNN) AFM exchange couplings (referred to as the J_1 - J_2 model) that produces a stripe-type AFM structure for $J_2 > J_1/2$ in the parent pnictides. One of the consequences of the magnetic frustration is the possible emergence of nematic degrees of freedom¹³ that can give rise to a short-range O order

above T_S .^{14–16} However, directly probing nematic fluctuations is nontrivial due to the difficulties in unequivocally decoupling their effects from that of the twin domains,¹⁷ as well as the fact that the magnetic field fluctuations associated with them average to zero. The roles of nematic fluctuations in shaping the magnetic and structural transitions and in the pairing mechanism that leads to superconductivity in the iron arsenides are under intense debate.^{16,18}

Here we report on high-resolution synchrotron x-ray diffraction studies of a CeFeAsO single crystal that enable us to separately monitor the development of each of the twin domains in this system. In particular, we examine the in-plane charge correlations as a function of temperature to gain insight on the two-dimensional (2D) coupling between the structural and magnetic properties of this typical ferropnictide.

CeFeAsO single crystals were synthesized in NaAs flux at ambient pressure, as described previously.¹⁹ Crystal quality was characterized by Laue backscattering, x-ray powder diffraction, heat capacity, magnetization, and resistivity measurements. A large as-grown platelike CeFeAsO single crystal (~ 5 mg) with the crystallographic c axis perpendicular to its surface was selected. The x-ray diffraction studies were carried out on the six-circle diffractometer of the 6-ID-B (XOR/MU) beamline at the Advanced Photon Source (APS) at Argonne National Laboratory. The x-ray energy throughout the experiment was kept at $E = 7.1000(5)$ keV. The incident radiation was linearly σ polarized with a spatial cross section of 1.0 mm (horizontal) \times 0.25 mm (vertical). The scattering geometry is shown in Fig. 2(a), where the χ angle represents a relative tilting of the sample with respect to \mathbf{Q} . In this configuration, charge scattering does not change the polarization of the scattered photons and occurs in the σ - σ' scattering channel. The mosaic of the single crystal is $\sim 0.09^\circ$ full width at half maximum (FWHM) as measured on the charge (205)_O Bragg

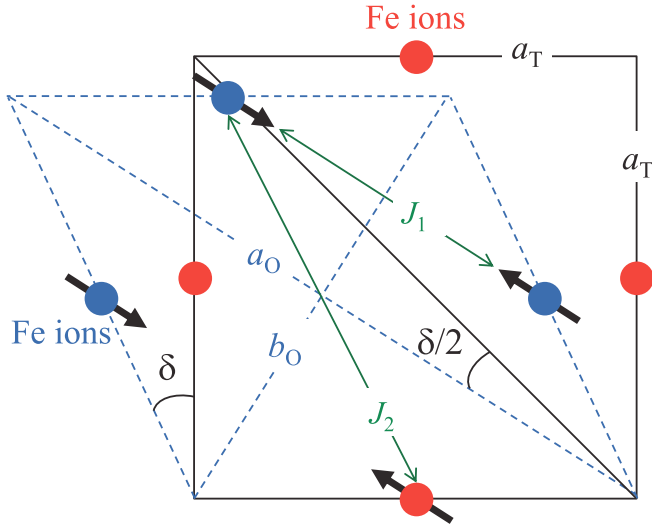


FIG. 1. (Color online) Schematic illustration of the formation of twin domains when the T structure transfers into the O phase in CeFeAsO. The O distortion in pnictides proceeds by shearing the T planar square into two rhombuses with angles $\frac{\pi}{2} \pm \delta$ (preserving the length of the square), creating twin domains. δ is the shearing angle. Here we just show one rhombus for clarity. J_1 and J_2 represent the NN and NNN AFM exchange interactions, respectively. a_T (T symmetry), and a_O and b_O (O symmetry) are the in-plane lattice constants. The arrows on the Fe ions represent the spins of iron ions.

reflection at 80 K. The sample was mounted at the end of the cold finger of a dilux cryogenic refrigerator with the ac plane as the scattering plane, and was measured at a temperature range of ~ 80 – 170 K. We note that the $(HKL)_T$ indices for the T symmetry correspond to the O reflection $(hkl)_O$ based on the relations of $h = H + K$, $k = H - K$, and $l = L$.

Upon cooling, at T_S , the T ($P4/nmm$) structure transfers into the O ($Cmma$) one. This is accompanied by an appearance of twin domains, e.g., the Bragg $(HK0)_T$ reflection in the T symmetry splits into twinned $(H + K, H - K, 0)_O / (H - K, H + K, 0)_O$ ones in the O symmetry. To obtain the angle misfit between the twin domains, we monitored the $(205)_O / (025)_O$ reflections. Figure 2(b) shows the temperature dependence of the χ angle [as illustrated in Fig. 2(a)] of both reflections. We find that for the k domain, represented by the $(025)_O$ reflection, the χ angle has no appreciable change within statistics in the investigated temperature range of ~ 80 – 164 K; whereas for the h domain, represented by the $(205)_O$ reflection, the χ angle gradually increases below ~ 130 K upon warming, and then steeply increases until merging into that of the k domain at T_S . The measured χ -angle difference $\Delta\chi = \chi_k - \chi_h$ is $\sim 0.183^\circ$ at 80 K, as indicated in Fig. 2(b).

The separation of the twin domains below T_S allows us to practically focus on an untwinned single crystal and follow the evolutions of the lattice constants and the correlations along the a and b axes. Figures 3(a) and 3(b) show the temperature dependence of the $(h05)$ and $(0k5)$ reflections from the h and k domains, respectively, indicating $T_S = 145(1)$ K. The integrated intensities of the longitudinal h and k scans from both domains shown in Fig. 3(c) exhibit two features associated with T_S and T_N . The dramatic jump in the intensity at T_S is clear evidence that the structural transition is first order

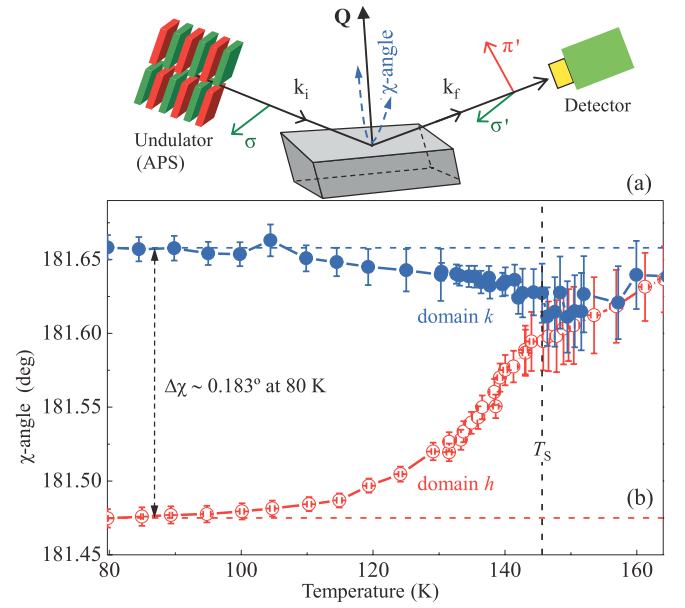


FIG. 2. (Color online) (a) Illustration of the x-ray diffraction setup at the 6-ID-B beamline at APS at Argonne National Laboratory. The χ angle indicates the directions for tilting the CeFeAsO single crystal with respect to \mathbf{Q} . (b) Temperature evolution of the χ angle of the h and k domains. The vertical dashed line implies the structural O-T transition temperature T_S . The horizontal dashed lines indicate the χ -angle difference between the h and k domains at 80 K. The solid lines are guides to the eye.

in character. This is also manifested by a ~ 2 K hysteresis effect as shown in Fig. 3(a'), consistent with similar observations in other pnictides.^{5,6} The other feature in Fig. 3(c), labeled by A and B, is more subtle and may result from the magnetic transition as discussed below, providing evidence for the coupling between lattice and spin orders.

Figure 3(d) shows the temperature dependence of the O strain S and the misfit-angle difference $\Delta\chi = \chi_k - \chi_h$. The value of the strain $S \sim 2.7 \times 10^{-3}$ at 80 K in CeFeAsO [$S \equiv (a_O - b_O) / (a_O + b_O)$, where a_O and b_O are the O lattice constants] is almost half the corresponding values in SrFe₂As₂ (Ref. 6) and CaFe₂As₂ (Ref. 5) compounds. This indicates that the out-of-plane coupling in the “1111” system is much weaker than that in the “122” family. The temperature dependence of the misfit-angle difference coincides well with that of the O strain, indicating a close relationship between microscopic (S) and macroscopic ($\Delta\chi$) parameters. This demonstrates that $\Delta\chi$ can serve as yet another probe to monitor the order parameter of the O-T structural transition.

The most remarkable observation in this paper is the temperature variation in the peak linewidth (i.e., the FWHM = κ) for both twin domains, as shown in Fig. 4(a). This linewidth obtained from the longitudinal h and k scans extends beyond the instrumental resolution, as demonstrated by comparing it with that of the corresponding l scans of the $(004)_O$ reflection [Fig. 4(a')], and is thus inversely proportional to the intrinsic in-plane charge-correlation lengths. At low temperatures ($< T_N$), the charge-correlation length along the h direction is significantly larger than that along the k direction, indicating that the h domain is majority consistent with the stronger

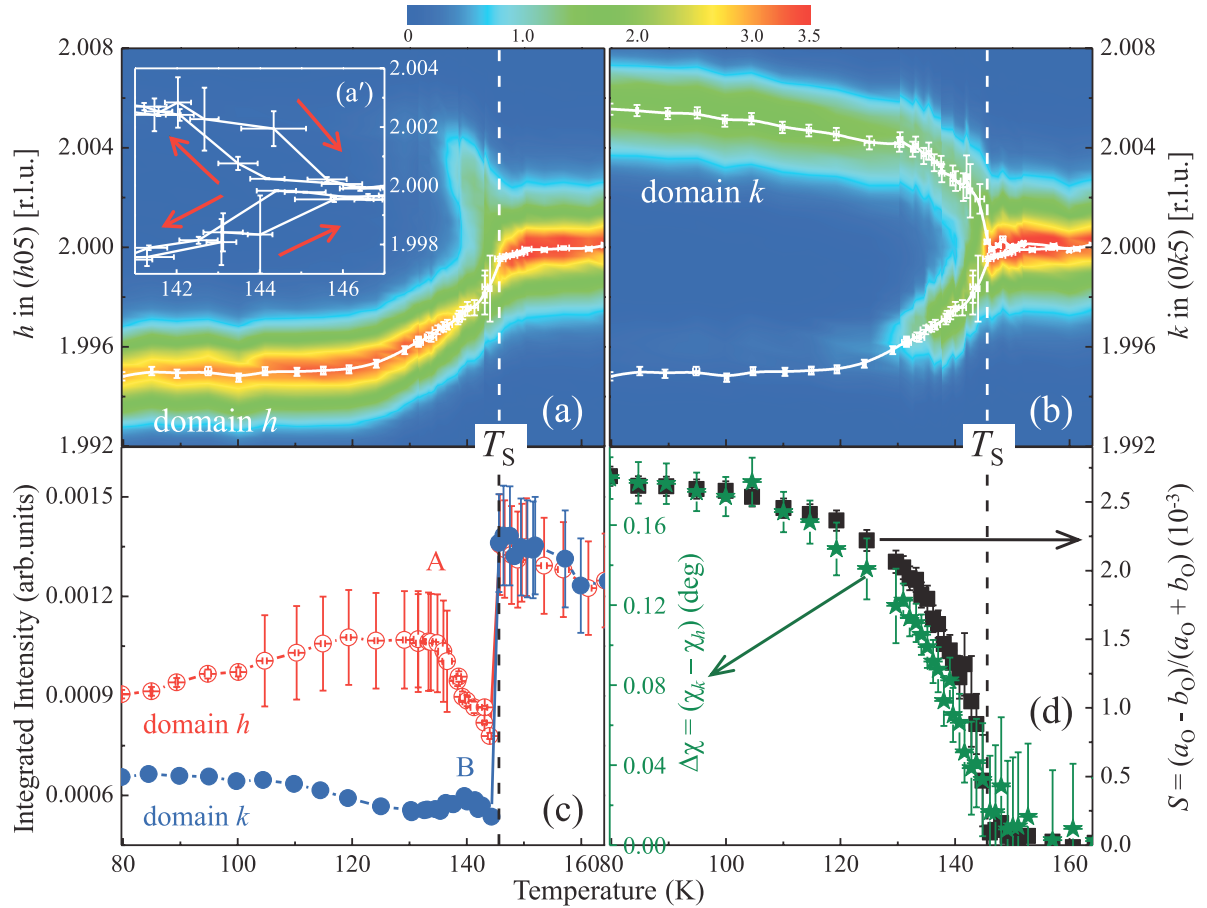


FIG. 3. (Color online) Temperature evolution of (a) the h domain and (b) the k domain. The solid line in (a) and the lower solid line in (b) indicate the Lorentzian-fit center of the longitudinal h scans at $\mathbf{Q} = (h05)_O$. The upper solid line in (b) indicates the Lorentzian-fit center of the longitudinal k scans at $\mathbf{Q} = (0k5)_O$. (c) Temperature evolutions of the integrated intensities of the longitudinal h and k scans at $\mathbf{Q} = (h05)_O$ and $(0k5)_O$, respectively, and (d) the O distortion (squares) (right) in the crystalline ab plane, namely, $S \equiv (a_0 - b_0)/(a_0 + b_0)$, and the domain-angle misfit (stars) (left), namely, $\Delta\chi = \chi_k - \chi_h$. (a') The ~ 2 K temperature hysteresis effect of the Lorentzian-fit centers indicative of a first-order O-T structural transition. Points A and B in (c) indicate the kink positions probably associated with the AFM transition, as described in the text. The vertical dashed lines in (a)–(d) indicate the structural O-T transition temperature T_S . The solid lines in (a)–(c) and (a') are guides to the eye.

intensity observed for the h domain in Fig. 3(c). This agrees well with the fact that the AFM interaction (interstripe) in this system is much stronger than the effective FM one (intraStripe), which results from the competing NNN interactions (J_2 , as illustrated in Fig. 1) that introduce frustration in the magnetic system. This charge-correlation-length anisotropy is also consistent with recent observations in an inelastic neutron-scattering study of SC $\text{Ba}(\text{Fe}_{0.926}\text{Co}_{0.074})_2\text{As}_2$ that show a similar anisotropy in the spin-correlation lengths.²⁰ The much smaller FWHM of the $(004)_O$ reflection compared to that of the h scans of the $(205)_O$ reflection at 80 K [Fig. 4(a')] in turn demonstrates that the spin-lattice coupling is 2D, and thus the magnetic exchange along the l direction is very weak. With the increase in temperature, the anisotropy becomes more and more pronounced, and two prominent peaklike features are observed at ~ 135 and ~ 145 K, which we associate with the stripe-type AFM and the O-T transitions, respectively. The relatively large broadening at T_S is due to the genuine increase in linewidth with temperature and the fact that the splitting may not be resolvable, and thus was treated as a

single peak. Figure 4(b) shows the temperature dependence of the normalized in-plane resistivity and its first derivative, clearly showing anomalies at T_S and T_N , consistent with the observations in Fig. 4(a).

In the T symmetry, above T_S , the linewidths are significantly larger than those at low temperatures. Similar observations have been reported in powder-diffraction studies of LaFeAsO .^{21,22} This indicates that what is measured is in fact a distribution in the d spacings due to the fluctuating O/T structures, which may not be resolved on the time scale and precision of the instrument. This seems to be an averaged-T phase with O/T fluctuations that are probably induced by the strong spin fluctuations typical in these systems.⁸ It should be noted that in general the correlation length diverges close to the O-T transition. However, the correlation lengths observed here are at their local minima close to the transition, characteristic of a martensiticlike transition, due to the shearing distortion in these systems.¹¹

The behavior of the asymmetric FWHM as shown in Fig. 4(a) may imply a remnant orthorhombicity above T_S ,

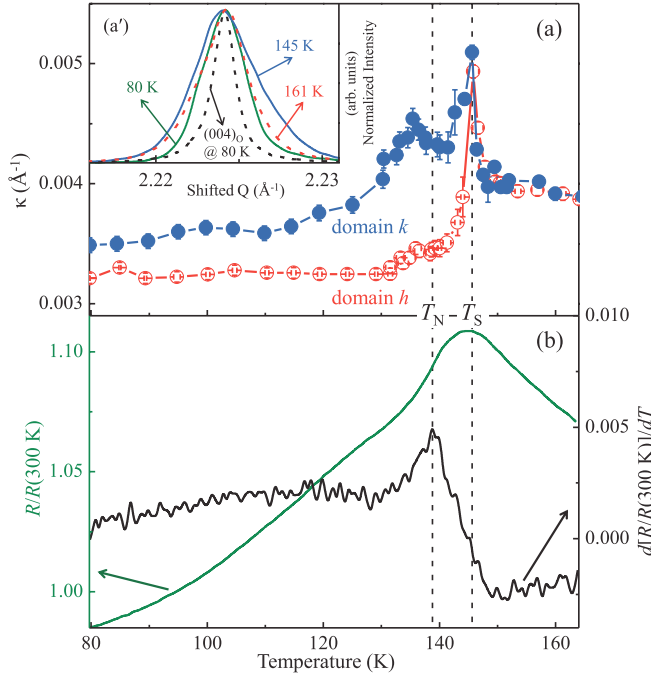


FIG. 4. (Color online) Temperature evolutions of (a) the FWHM κ of the longitudinal h and k scans at $\mathbf{Q} = (h05)_O$ and $(0k5)_O$, respectively, and (b) the resistivity measurement (left) and its first derivative (right) indicate the temperatures of the AFM transition of iron ions (T_N^{Fe}) as well as the O-T structural transition (T_S). (a') shows the longitudinal h scans at $\mathbf{Q} = (h05)_O$ and at temperatures of 80, 145, and 161 K. To estimate the instrumental resolution effect, we also showed the corresponding l scan of the $(004)_O$ reflection at 80 K. The \mathbf{Q} values and observed intensities were shifted and normalized, respectively, for comparison. The vertical dashed lines in (a) and (b) indicate the AFM and the O-T structural transition temperatures T_N and T_S , respectively. The solid lines in (a) are guides to the eye.

while the gradual decrease of the FWHM of the k domain below ~ 130 K may indicate a remnant tetragonality below T_S . It is possible that the remnant tetragonality follows the k domain because both lattice constants b_O and b_T decrease upon cooling. Therefore, we argue that both T and O phases may coexist dynamically in a certain temperature range around T_S , most likely due to the strong magnetic fluctuations, and, perhaps, due to spin nematic degrees of freedom.^{7,14,20} This scenario may explain the small temperature hysteresis effect [Fig. 3(a')] and the gradual increase of the order parameters in the first-order structural O-T transition below T_S [Fig. 3(d)].

Our discussion above assumes that the variations in the linewidths are purely from correlations. Another alternative is the domain-size effect, in which the variations indicate that in the critical fluctuation regime, magnetoelastic coupling already nucleates anisotropic domains, giving rise to elongated domains upon freezing, consistent with the observation in electron-microscopy studies.¹¹

It is interesting to note that the structural O-T transition temperature $T_S = 145(1)$ K determined in Fig. 3(d) is also reflected in the linewidths [Fig. 4(a)] and in the integrated intensities [Fig. 3(c)]. This shows that the temperature variation of integrated intensities could be useful in determining the O-T structural transition temperature in these systems.⁶ Similar observations by high-resolution x-ray diffraction were reported for TbVO_4 and TbAsO_4 ,²³ where the T-to-O structural transition is driven by the cooperative Jahn-Teller distortion rather than spin fluctuations.

To summarize, we demonstrate that with high-resolution synchrotron x-ray diffraction, the twin domains observed in single-crystal CeFeAsO can be distinct, allowing the practical study of an individual untwinned crystal. Most importantly, we find that the charge correlations show a significant anisotropy along and perpendicular to the stripe-type 2D AFM wave vector. This is consistent with the anisotropic 2D spin correlations, indicative of an anisotropic 2D magnetoelastic coupling, and implies that the AFM ordering is the driving force for the structural T-to-O transition in this system. In addition, we show that the misfit-angle difference of the O twin domains (a macroscopic variable) follows well the O strain (or shearing) order parameter, indicating that it can be used to monitor the order parameter of the O-T structural transition. Furthermore, based on the observed charge-correlation lengths of the O twin domains and the T domain below and above T_S , respectively, we argue that at high temperatures the system is virtually a T phase that exhibits strong O/T structural fluctuations. This fluctuating phase is probably induced by the strong spin fluctuations, most likely by the spin-nematic phase.

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¹Y. Kamihara, T. Watanabe, M. Hirano, and H. Hosono, *J. Am. Chem. Soc.* **130**, 3296 (2008).

²X. H. Chen, T. Wu, G. Wu, R. H. Liu, H. Chen, and D. F. Fang, *Nature (London)* **453**, 761 (2008).

³Z. A. Ren, W. Lu, J. Yang, W. Yi, X. L. Shen, Z. C. Li, G. C. Che, X. L. Dong, L. L. Sun, F. Zhou, and Z. X. Zhong, *Chin. Phys. Lett.* **25**, 2215 (2008).

⁴J.-Q. Yan, A. Kreyssig, S. Nandi, N. Ni, S. L. Bud'ko, A. Kracher, R. J. McQueeney, R. W. McCallum, T. A. Lograsso, A. I. Goldman, and P. C. Canfield, *Phys. Rev. B* **78**, 024516 (2008).

⁵A. I. Goldman, D. N. Argyriou, B. Ouladdiaf, T. Chatterji, A. Kreyssig, S. Nandi, N. Ni, S. L. Bud'ko, P. C. Canfield, and R. J. McQueeney, *Phys. Rev. B* **78**, 100506(R) (2008).

⁶H.-F. Li, W. Tian, J. L. Zarestky, A. Kreyssig, N. Ni, S. L. Bud'ko, P. C. Canfield, A. I. Goldman, R. J. McQueeney, and D. Vaknin, *Phys. Rev. B* **80**, 054407 (2009).

⁷A. Jesche, C. Krellner, M. de Souza, M. Lang, and C. Geibel, *Phys. Rev. B* **81**, 134525 (2010).

⁸S. O. Diallo, V. P. Antropov, T. G. Perring, C. Broholm, J. J. Pulikotil, N. Ni, S. L. Bud'ko, P. C. Canfield, A. Kreyssig,

- A. I. Goldman, and R. J. McQueeney, *Phys. Rev. Lett.* **102**, 187206 (2009).
- ⁹J. Zhao, D. T. Adroja, D.-X. Yao, R. Bewley, S. L. Li, X. F. Wang, G. Wu, X. H. Chen, J. P. H., and P. C. Dai, *Nat. Phys.* **5**, 555 (2009).
- ¹⁰S. Wakimoto, K. Kodama, M. Ishikado, M. Matsuda, R. Kajimoto, M. Arai, K. Kakurai, F. Esaka, A. Iyo, H. Kito, H. Eisaki, and S. Shamoto, *J. Phys. Soc. Jpn.* **79**, 074715 (2010).
- ¹¹J. C. Loudon, C. J. Bowell, J. Gillett, S. E. Sebastian, and P. A. Midgley, *Phys. Rev. B* **81**, 214111 (2010).
- ¹²J. Paglione and R. Greene, *Nat. Phys.* **6**, 645 (2010).
- ¹³P. Chandra, P. Coleman, and A. I. Larkin, *Phys. Rev. Lett.* **64**, 88 (1990).
- ¹⁴C. Fang, H. Yao, W.-F. Tsai, J. P. Hu, and S. A. Kivelson, *Phys. Rev. B* **77**, 224509 (2008).
- ¹⁵C. Xu, M. Müller, and S. Sachdev, *Phys. Rev. B* **78**, 020501(R) (2008).
- ¹⁶R. M. Fernandes, L. H. VanBebber, S. Bhattacharya, P. Chandra, V. Keppens, D. Mandrus, M. A. McGuire, B. C. Sales, A. S. Sefat, and J. Schmalian, *Phys. Rev. Lett.* **105**, 157003 (2010).
- ¹⁷J. Chu, J. Analytis, K. Greve, P. McMahon, Z. Islam, Y. Yamamoto, and I. Fisher, *Science* **329**, 824 (2010).
- ¹⁸T. Chuang, M. Allan, J. Lee, Y. Xie, N. Ni, S. Bud'ko, G. Boebinger, P. Canfield, and J. Davis, *Science* **327**, 181 (2010).
- ¹⁹J.-Q. Yan, S. Nandi, J. L. Zarestky, W. Tian, A. Kreyssig, B. Jensen, A. Kracher, K. W. Dennis, R. J. McQueeney, A. I. Goldman, R. W. McCallum, and T. A. Lograsso, *Appl. Phys. Lett.* **95**, 222504 (2009).
- ²⁰H.-F. Li, C. Broholm, D. Vaknin, R. M. Fernandes, D. L. Abernathy, M. B. Stone, D. K. Pratt, W. Tian, Y. Qiu, N. Ni, S. O. Diallo, J. L. Zarestky, S. L. Bud'ko, P. C. Canfield, and R. J. McQueeney, *Phys. Rev. B* **82**, 140503(R) (2010).
- ²¹H.-F. Li, W. Tian, J.-Q. Yan, J. L. Zarestky, R. W. McCallum, T. A. Lograsso, and D. Vaknin, *Phys. Rev. B* **82**, 064409 (2010).
- ²²N. Qureshi, Y. Drees, J. Werner, S. Wurmehl, C. Hess, R. Klingeler, B. Buchner, M. T. Fernández-Díaz, and M. Braden, *Phys. Rev. B* **82**, 184521 (2010).
- ²³K. C. Rule, M. J. Lewis, H. A. Dabkowska, D. R. Taylor, and B. D. Gaulin, *Phys. Rev. B* **77**, 134116 (2008).