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# Interplay between crystallinity profiles and the performance of microcrystalline thin-film silicon solar cells studied by in-situ Raman spectroscopy

T. Fink, S. Muthmann, A. Mück, A. Gordijn, R. Carius, and M. Meier

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The intrinsic microcrystalline absorber layer growth in thin-film silicon solar-cells is investigated by in-situ Raman spectroscopy during plasma enhanced chemical vapor deposition. In-situ Raman spectroscopy enables a detailed study of the correlation between the process settings, the evolution of the Raman crystallinity in growth direction, and the photovoltaic parameters  $\eta$  (solar cell conversion efficiency),  $J_{SC}$  (short circuit current density),  $FF$  (fill factor), and  $V_{OC}$  (open circuit voltage). Raman spectra were taken every 7 nm of the absorber layer growth depending on the process settings. The Raman crystallinity of growing microcrystalline silicon was determined with an absolute error of approximately  $\pm 5\%$  for total absorber layer thicknesses  $> 50$  nm. Due to this high accuracy, inherent drifts of the Raman crystallinity profiles are resolvable for almost the entire absorber layer deposition. For constant process settings and optimized solar cell device efficiency Raman crystallinity increases during the absorber layer growth. To compensate the inhomogeneous absorber layer growth process settings were adjusted. As a result, absorber layers with a constant Raman crystallinity profile — as observed in-situ — were deposited. Solar cells with those absorber layers show a strongly enhanced conversion efficiency by  $\sim 0.5\%$  absolute. However, the highest  $FF$ ,  $V_{OC}$ , and  $J_{SC}$  were detected for solar cells with different Raman crystallinity profiles. In particular, fill factors of 74.5% were observed for solar cells with decreasing Raman crystallinity during the later absorber layer growth. In contrast, intrinsic layers with favorable  $J_{SC}$  are obtained for constant and increasing Raman crystallinity profiles. Therefore, monitoring the evolution of the Raman crystallinity in-situ provides sufficient information for an optimization of the photovoltaic parameters with surpassing depth resolution. © 2015 AIP Publishing LLC.

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## I. INTRODUCTION

Microcrystalline silicon ( $\mu c$ -Si:H) grown by plasma enhanced chemical vapor deposition (PECVD) is used as an absorber layer in thin-film solar cell stacks because of its high absorption coefficient for long wavelengths and due to the stability against light induced degradation.<sup>1–3</sup> In the last two decades, a huge effort was made to enhance the material quality of  $\mu c$ -Si:H.<sup>4–8</sup> Concerning the material quality of intrinsic absorber layers, crystalline volume fraction in the growth direction has a great impact on  $\eta$ ,  $FF$ ,  $V_{OC}$  and  $J_{SC}$ .<sup>1</sup> It is common practice to estimate the crystalline volume fraction by the Raman crystallinity.<sup>9,10</sup>

The growth of  $\mu c$ -Si:H is influenced by a large number of external and internal process parameters (substrate morphology,<sup>11</sup> seeding effect of doping layers,<sup>12</sup> plasma excitation frequency,<sup>6</sup> reactor geometry,<sup>13</sup> etc.). Therefore, the reproducibility of Raman crystallinity profiles in the growth direction is challenging and optimized Raman crystallinities throughout the absorber layer are difficult to realize. To address this issue innovative diagnostic techniques and a more profound understanding of the interdependence between the PECVD process settings, the evolution of the Raman crystallinity, and the photovoltaic parameters is required.

At low Raman crystallinities below the percolation threshold the charge carrier mobility is restricted.<sup>14</sup> For high

Raman crystallinities grain boundaries are poorly passivated and  $\mu c$ -Si:H exhibits high density of twin defects and voids.<sup>1,15,16</sup> Although the conductivity is enhanced in this case, the density of recombination centers as well as the atmospheric gas incorporation is increased.<sup>1,17</sup> The highest solar cell performances were detected for absorber layers with Raman crystallinities of about 60%–70%.<sup>1,12,15,18–20</sup> The ideal evolution of the Raman crystallinity throughout the absorber layer is still a matter of debate. Therefore, this is one of the questions that will be addressed here.

In general, PECVD-inherent process drifts<sup>19,21</sup> and the conical evolution of the crystallites<sup>1,9,22</sup> for low Raman crystallinities impede a homogeneous absorber layer growth in the narrow deposition window close to 60%–70%.<sup>21</sup> To achieve these Raman crystallinities throughout the entire absorber layer process settings have to be adapted during the deposition of  $\mu c$ -Si:H. Since this requires an adequate monitoring of the Raman crystallinity, various in-situ and ex-situ measurement techniques have been applied to characterize the  $\mu c$ -Si:H absorber layer deposition.<sup>8–10,12,13,19,21–28</sup>

To monitor the layer growth with adequate accuracy, diagnostic methods with high depth and temporal resolution are required. For example, in-situ ellipsometry features high surface sensitivity.<sup>8,22</sup> However, incident angles of  $70^\circ$ – $80^\circ$  complicate the integration of the ellipsometry setup into the

parallel plate reactor and data analysis becomes difficult for varying substrate morphology.<sup>29</sup> The plasma emission of the PECVD process is recorded with excellent temporal resolution by optical emission spectroscopy (OES),<sup>13,25</sup> but material properties of the growing layer are not probed. High depth resolution of the Raman crystallinity in the growth direction was achieved by KOH wet-etching or reactive ion etching.<sup>12,21,24</sup> Scanning the etched craters by microscopic Raman spectroscopy, the Raman crystallinity is estimated in the growth direction of the  $\mu\text{c-Si:H}$  film ex-situ.<sup>24</sup>

In the present paper, we investigate the microcrystalline silicon growth by the technique of in-situ Raman spectroscopy.<sup>30</sup> This method enables the monitoring of the Raman crystallinity during the deposition of the absorber layer. It is possible to adjust the Raman excitation laser under perpendicular incidence to the probed surface, because an optimized electric shielding of the optical feed through was achieved.<sup>31</sup> Moreover, data analysis is independent of the morphology of the growing film. The depth resolution of in-situ Raman spectroscopy is sufficient to investigate the initial deposition phase of  $\mu\text{c-Si:H}$  absorber layers as well as to study the impact of the initial Raman crystallinity on the evolution of the later layer growth in detail. Muthmann *et al.* reported on the capabilities of that technique to study the Raman crystallinity during the initial growth phase and to observe the effect of an adjusted silane to hydrogen ratio.<sup>32</sup> Here, we monitor the evolution of the Raman crystallinity throughout the entire intrinsic layer. The inhomogeneous absorber layer growth is successively compensated by an iterative readjustment of the process settings and the impact on the solar cell parameters is discussed.

## II. EXPERIMENTAL

For the preparation of p-i-n single junction silicon solar cells in superstrate configuration with a  $\mu\text{c-Si:H}$  intrinsic absorber layer a multi chamber system for the deposition on  $10 \times 10 \text{ cm}^2$  corning glass was used. The approximately 30 nm thick microcrystalline boron doped p-type layers were grown on aluminum doped wet-etched zinc oxide substrates<sup>33</sup> (front contact). Although 30 nm thick p-layers are optically unfavorable, stable and reproducible starting conditions for the absorber layer deposition were provided. To complete the p-i-n single-junction solar cells ( $1300 \pm 70$ ) nm thick intrinsic absorber layers and amorphous phosphorous doped n-type layers were deposited. The back contact consists of thermally evaporated silver. An active solar cell area of  $1 \times 1 \text{ cm}^2$  is defined by shadow masks for each single solar cell. The I-V-parameters were measured under AM 1.5 illuminations in a class A solar simulator. Each data point of the photovoltaic parameters was calculated by the average of the three best  $1 \times 1 \text{ cm}^2$  solar cells. For this calculation absolute measurement errors of the photovoltaic parameters are  $\Delta\eta = \pm 0.2\%$ ,  $\Delta J_{\text{SC}} = \pm 0.5 \text{ mA/cm}^2$ ,  $\Delta FF = \pm 0.5\%$ , and  $\Delta V_{\text{OC}} = \pm 5 \text{ mV}$ . Solar cell conversion efficiency  $\eta$  and short circuit current density  $J_{\text{SC}}$  shown in this study are smaller than for recently reported record devices.<sup>34–36</sup> Low  $J_{\text{SC}}$  most likely results from parasitic absorption in the comparably

thick p-layer and from plasmonic losses at the silver back contact.<sup>37,38</sup>

For the plasma deposition of the intrinsic absorber layer an excitation frequency of 13.56 MHz in parallel plate electrode geometry with a power density of  $\sim 0.4 \text{ W/cm}^2$  was applied. For the plasma etching experiments in Section III B, power density of  $\sim 0.8 \text{ W/cm}^2$  differs from the other experiments of the manuscript. The substrate heater temperature was set to  $200^\circ\text{C}$ . The deposition pressure was kept constant at 10 Torr and the deposition gases were silane and hydrogen. From the gas flows of silane ( $\Phi_{\text{SiH}_4}$ ) and hydrogen ( $\Phi_{\text{H}_2}$ ) the silane concentration  $SC$  is defined as

$$SC = \frac{\Phi_{\text{SiH}_4}}{\Phi_{\text{SiH}_4} + \Phi_{\text{H}_2}}. \quad (1)$$

Silane concentration  $SC$  was changed significantly for the two-step process control. Nevertheless, the constant gas flow of hydrogen  $\Phi_{\text{H}_2}$  (360 sccm) is roughly 2 orders of magnitude higher than the varied silane gas flow  $\Phi_{\text{SiH}_4}$  ( $\sim 3$  sccm). Accordingly, the total gas flow rate was approximately constant during the absorber layer deposition. For the plasma etching experiments in Section III B, silane concentration was constant  $SC = 1.4\%$ .

To measure the Raman crystallinity in-situ, an optical feed through was integrated into the PECVD electrode.<sup>30,31</sup> The Raman crystallinity was measured during the growth of the i-layer with a solid state Nd:YAG 532 nm laser. The information depth, which is probed by the laser, can be defined as the depth from which more than  $2/e$  of the scattered Raman light originates. In our case, the information depth for a-Si:H and  $\mu\text{c-Si:H}$  can be estimated to 56 nm and 125 nm for 532 nm excitation wavelength,<sup>39</sup> respectively. This implies that the Raman scattered light of the microcrystalline p-layer contributes to the Raman signal during the initial growth phase of the i-layer. However, for thin p-layer thicknesses of about 30 nm, the contribution is small. To reduce the heat input on the growing sample, the laser beam was chopped to 200  $\mu\text{s}$  pulses (laser light exposure time) and the focus spot on the sample was broadened by a cylindrical lens. A temporal resolution of 25 s allows acquiring Raman spectra every  $\sim 7$  nm of silicon film thickness when growth rates between 0.25 nm/s and 0.30 nm/s are assumed. For the analysis of the Raman spectra the reference spectrum of a pure amorphous sample was fitted to the data to determine the amorphous intensity contribution  $I_a$  of the Raman scattered light in the microcrystalline material.<sup>10</sup> Accordingly, the Raman crystallinity is calculated as

$$I_{\text{C}}^{\text{RS}} = \frac{I_{\text{c}}}{I_{\text{c}} + I_{\text{a}}}, \quad (2)$$

in which  $I_{\text{c}}$  stands for the intensity contribution of the light scattered in the crystalline parts of the  $\mu\text{c-Si:H}$  layer. To analyze the recorded spectra, different background signals were considered: On the one hand, the Raman scattered light of the optical components in the experimental setup, which stays constant during the deposition, was subtracted. On the other hand, the Raman scattered light of the corning glass

carrier material and the ZnO:Al front contact layer contributes to the i-layer signal during the initial growth phase. For the subtraction of this contribution the Raman intensity of the glass-ZnO:Al substrate was recorded prior to the p-layer deposition. From the data analysis of the Raman spectra, an absolute error of  $I_C^{RS}$  is estimated by  $\Delta I_C^{RS} \leq \pm 5\%$  for silicon layer thicknesses  $d > 50$  nm.

To compensate the inhomogeneous absorber layer growth, SC was readjusted in two steps (Fig. 1).

In region **A** the initial growth phase was investigated (Section III B). Here, a reduced initial silane concentration  $SC_i$  was applied to compensate the initially low  $I_C^{RS}$ . Therefore,  $SC_i$  was varied for an initial deposition time  $t \leq 180$  s. For  $t > 180$  s silane concentration SC was slightly increased and kept constant at 1.04% during the investigation of the initial growth phase (Fig. 1): It is well known that an increase of  $I_C^{RS}$  for the initial growth phase induces an overall increase of  $I_C^{RS}$  for the later layer growth.<sup>12</sup> Increasing SC avoids strongly deviating crystallinity from advantageous  $I_C^{RS}$  of  $\sim 60\%$ – $70\%$ <sup>1,12,15,18–20</sup> (Fig. 1). In a last step (Section III C), long-term drifts of  $I_C^{RS}$ , which occur for the later layer growth in region **B** (Fig. 1), were compensated by continuously changing SC with slope

$$\Delta SC = \frac{SC_B}{t_{\text{deposition}} - 180 \text{ s}}, \quad (3)$$

where  $SC_B$  stands for the total increase of SC in region **B** and  $t_{\text{deposition}}$  denotes the total time of the i-layer deposition ( $\Delta SC$  in units of %/h; h: hour).

### III. RESULTS AND DISCUSSION

#### A. Crystallinity profiles for constant process settings

First,  $I_C^{RS}$  profiles for different constant silane concentrations during the i-layer deposition were recorded to study the influence of the process settings (SC) on the Raman crystallinity of the  $\mu\text{c-Si:H}$  i-layer in the growth direction. Using this information it is possible to adapt the external process parameters to achieve a more homogeneous layer growth.

The crystallinity profiles during the deposition of the intrinsic  $\mu\text{c-Si:H}$  layer are shown in Fig. 2, where the Raman

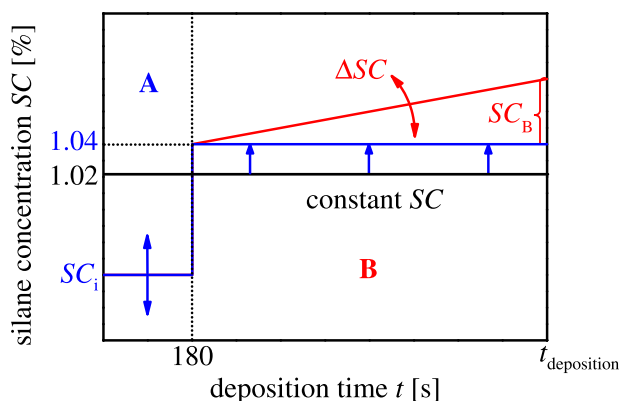


FIG. 1. Schematic of the silane concentration (SC) adjustment during absorber layer deposition: To study the influence on the initial growth phase (**A**),  $SC_i$  was varied ( $t \leq 180$  s). To compensate the long-term drifts of the Raman crystallinity during the later layer growth, the silane concentration was changed with a slope of  $\Delta SC$  (**B**).

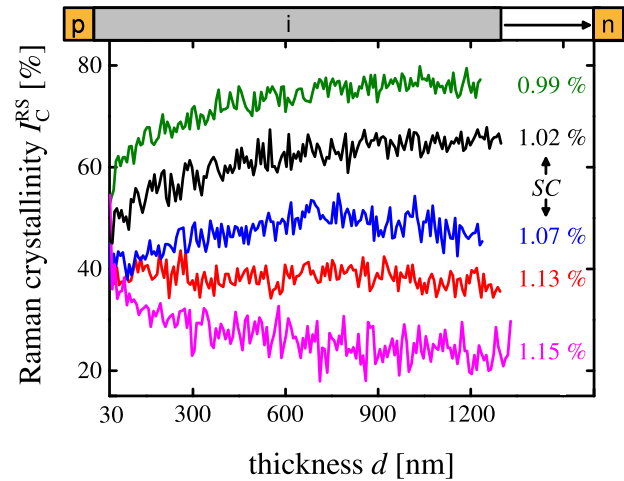


FIG. 2. Raman crystallinities  $I_C^{RS}$  of intrinsic microcrystalline silicon layers in dependence of the i-layer thickness for  $d > 30$  nm, deposited at constant silane concentrations SC and measured by in-situ Raman spectroscopy.

crystallinity is plotted as a function of intrinsic layer thickness  $d$ . All crystallinity profiles start at a Raman crystallinity which is close to that of the p-layer ( $\sim 58\%$ ). The upper two crystallinity profiles for  $SC = 0.99\%$  and  $SC = 1.02\%$  show strongly increasing  $I_C^{RS}$  at the beginning and a moderate increase of  $I_C^{RS}$  during the later layer growth. The lower two crystallinity profiles for  $SC = 1.13\%$  and  $SC = 1.15\%$  in Fig. 2 show decreasing  $I_C^{RS}$  during the deposition which is less pronounced for  $SC = 1.13\%$ . The crystallinity profile in the middle ( $SC = 1.07\%$ ) features a reversing curve shape: At the beginning  $I_C^{RS}$  increases up to a thickness of  $\sim 700$ – $800$  nm and thereafter  $I_C^{RS}$  decreases again by  $\sim 5\%$  absolute, which is in the range of the measurement error. Nevertheless, this effect was validated by the deposition of thick layers with  $d > 2 \mu\text{m}$  and additionally by KOH crater etching experiments.<sup>40</sup>

To explain the similar starting crystallinity of the Raman profiles close to that of the p-layer ( $I_C^{RS} \approx 58\%$ ) two possible reasons have to be considered: One is the seeding effect of the p-layer which leads to a continuation of the crystalline growth, if deposition conditions are favorable. Additionally, the Raman scattered light partly originates from the p-layer during the initial growth phase due to the information depth of the 532 nm excitation laser.

For silane concentrations  $SC = 0.99\%$  and  $SC = 1.02\%$  low initial  $I_C^{RS}$  is observed compared to the later layer growth (Fig. 2). Low initial  $I_C^{RS}$  results from transient depletion of the process gases which typically occurs in PECVD processes.<sup>13,23,41</sup> For the subsequent layer growth crystallites evolve conically and  $I_C^{RS}$  increases.<sup>1,9</sup>

Decreasing  $I_C^{RS}$  under comparably high SC conditions was also observed previously by TEM measurements:<sup>9</sup> In these experiments thin-film silicon was grown on crystalline silicon substrates for process conditions which promote low  $I_C^{RS}$ . In this case, the microstructure of  $\mu\text{c-Si:H}$  material differs from conically growing crystallites and resembles randomly shaped crystallites surrounded by a-Si:H.<sup>9,42</sup>

The reversing curve shape for  $SC = 1.07\%$  might be explained by different process drifts which compensate each other over time.<sup>19</sup>



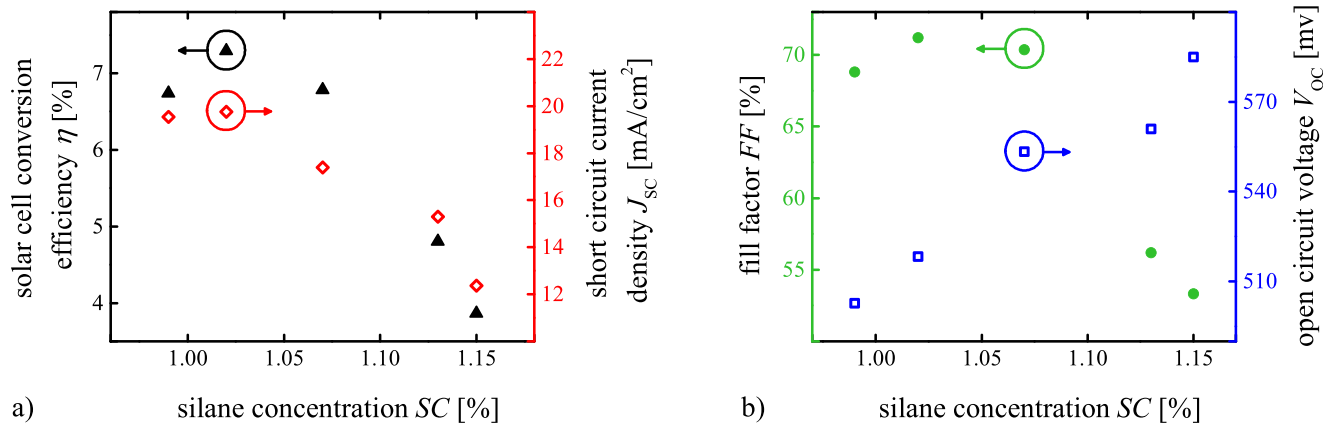


FIG. 3. (a) Solar cell conversion efficiency  $\eta$  resp. short circuit current density  $J_{sc}$  and (b) fill factor  $FF$  resp. open circuit voltage  $V_{oc}$  in dependence of the silane concentration  $SC$ . Maximum absolute measurement errors of the determined photovoltaic parameters are  $\Delta\eta = \pm 0.2\%$ ,  $\Delta J_{sc} = \pm 0.5 \text{ mA/cm}^2$ ,  $\Delta FF = \pm 0.5\%$  and  $\Delta V_{oc} = \pm 5 \text{ mV}$ .

The Raman profiles shown in Fig. 2 provide direct information about the evolution of Raman crystallinity in  $\mu\text{c-Si:H}$  absorber layers. To correlate the  $I_C^{RS}$  profiles with the device performance, conversion efficiency  $\eta$ , short circuit current density  $J_{sc}$ , fill factor  $FF$ , and open circuit voltage  $V_{oc}$  were measured (Fig. 3). As expected,<sup>1</sup>  $V_{oc}$  is increasing with increasing  $SC$ . The fill factor exhibits a maximum of 71.2% for  $SC = 1.02\%$ . At the same  $SC$  the solar cell conversion efficiency  $\eta$  has its maximum at  $\eta = 7.29\%$  and drops for higher  $SC$ . In the considered  $SC$  range the highest  $J_{sc}$  and  $\eta$  are observed for  $SC \approx 1.02\%$ .

For constant process settings, the maximum of  $\eta$  (Fig. 3(a)) is achieved for  $SC = 1.02\%$ . The crystallinity profile for  $SC = 1.02\%$  shows an average  $I_C^{RS}$  close to 60% and increases by about 15%–20% absolutely over the thickness of the intrinsic layer (Fig. 2). High depth resolution of in-situ Raman spectroscopy enables estimation of the adjustment of  $SC$  which is required to realize homogeneous layer growth.

## B. Effects of modified Raman-crystallinity during the initial growth phase

As described above, low  $SC$  was applied to increase  $I_C^{RS}$  of the i-layer close to the p-type layer. Beforehand, the growth rate of the initial layer was estimated. For this reason, a sequence of thin layers was deposited and the layer

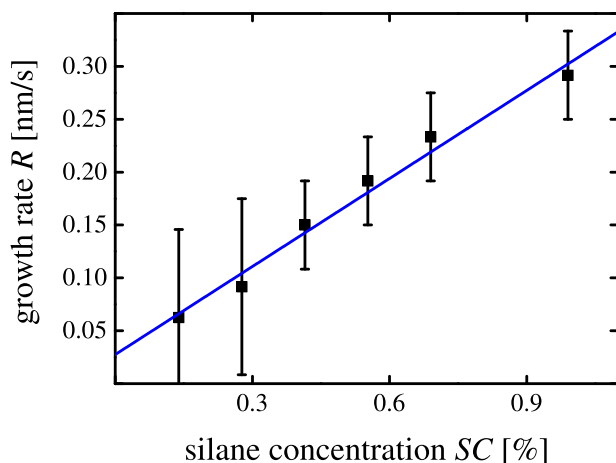


FIG. 4. Growth rate  $R$  as a function of silane concentration  $SC$ .

thickness was determined by a profilometer (Fig. 4). The growth rate increases approximately linearly with  $SC$  according to literature.<sup>43–45</sup> After a deposition time  $t_i = 180 \text{ s}$  of the  $\mu\text{c-Si:H}$  i-layer, the layer thickness for  $SC = 1.02\%$  is approximately 50 nm. Up to this thickness, the highest rate of change is observed in terms of the crystallinity evolution (Fig. 2).

An optimum of the device performance is expected when  $SC_i$  is decreased because for initially high  $SC_i$  of 1.02% unfavorable Raman crystallinity  $I_C^{RS} < 60\%$  was obtained at the beginning of the i-layer growth (Fig. 2). To find out the deposition conditions for optimum device performance,  $SC_i$  was reduced and varied.

The photovoltaic parameters  $\eta$ ,  $J_{sc}$ ,  $FF$ , and  $V_{oc}$  of solar cells deposited with and without changing the initial silane concentration  $SC_i$  are shown in Fig. 5. A maximum of the solar cell conversion efficiency  $\eta$  is observed for  $SC_i = 0.83\%$  (Fig. 5(a)). Compared to the results under constant deposition conditions ( $SC_i = 1.02\%$ ),  $\eta$  was increased by  $\sim 0.4\%$  absolute. Short circuit current density  $J_{sc}$  is improved by approximately  $1 \text{ mA/cm}^2$  for reduced  $SC_i$ . Fill factor  $FF$  as well as  $V_{oc}$  are shown in Fig. 5(b): In comparison to the reference ( $SC_i = 1.02\%$ ),  $FF$  is slightly increased in the range between  $SC_i = 0.61\%$  and  $SC_i = 0.83\%$ . In contrast,  $V_{oc}$  features no local maximum and decreases by almost 20 mV for decreasing  $SC_i$ .

Obviously, the local maximum of  $\eta$  results from opposing trends of increasing  $J_{sc}$  and decreasing  $V_{oc}$  for reduced initial  $SC_i$ .

Since the initial growth phase of the i-layer is changed, we investigated  $J_{sc}$  wavelength dependent by measuring the external quantum efficiency  $EQE$ . In Fig. 6,  $EQE$  is shown for constant process settings ( $SC = 1.02\%$ ) and for decreased initial  $SC_i$  ( $t \leq 180 \text{ s}$ ,  $SC(t > 180 \text{ s}) = 1.04\%$ ): In the long wavelength range for  $\lambda > 700 \text{ nm}$ , no significant modification of  $EQE$  is observed. For wavelengths  $\lambda < 550 \text{ nm}$  and  $SC = 1.02\%$   $EQE$  is not varied significantly with and without reverse bias voltage  $V_{Bias}$  (Fig. 6). For slightly decreased initial  $SC_i$  from  $SC_i = 1.02\%$  to  $SC_i = 0.91\%$   $EQE$  is enhanced and total  $J_{sc}$  is increased from  $19.6 \text{ mA/cm}^2$  to  $20.3 \text{ mA/cm}^2$  (Fig. 5(a)). The significance of this result is substantiated by an absolute error of  $\Delta EQE = \pm 2\%$  resp.  $\Delta J_{sc} = \pm 0.5 \text{ mA/cm}^2$ . A further decrease of  $SC_i$  ( $SC_i < 0.91\%$ ) results in minor changes of  $EQE$  and  $J_{sc}$  only (Figs. 5(a) and 6).

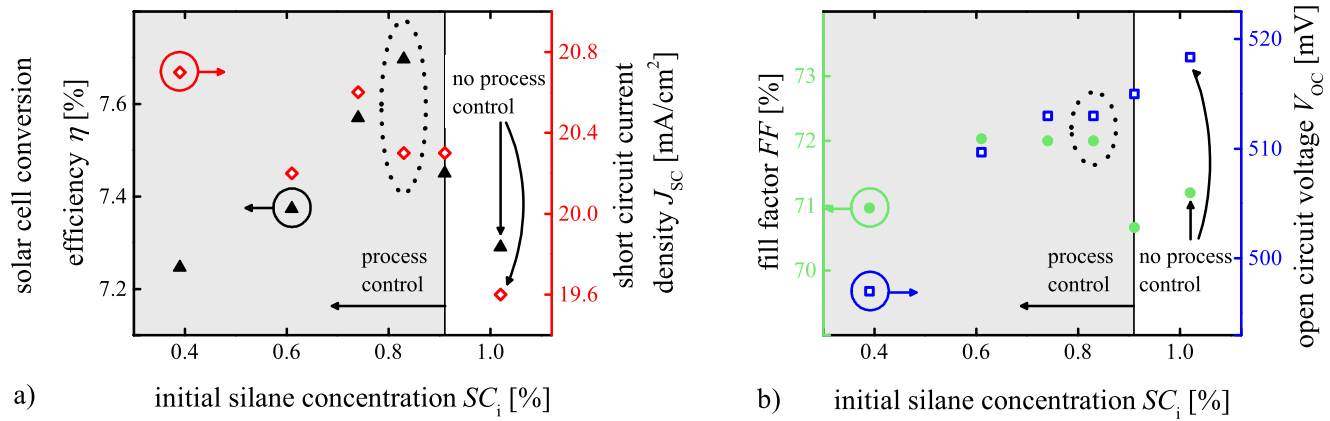


FIG. 5. (a) Solar cell conversion efficiency  $\eta$  resp. short circuit current density  $J_{sc}$  and (b) fill factor  $FF$  resp. open circuit voltage  $V_{oc}$  as a function of initial silane concentration  $SC_i$  ( $SC(t > 180 \text{ s}) = 1.04\%$ ). The data points at  $SC_i = 1.02\%$  (constant  $SC$ ) represents the reference. Maximum absolute measurement errors of the determined photovoltaic parameters are  $\Delta\eta = \pm 0.2\%$ ,  $\Delta J_{sc} = \pm 0.5 \text{ mA/cm}^2$ ,  $\Delta FF = \pm 0.5\%$  and  $\Delta V_{oc} = \pm 5 \text{ mV}$ .

After the adjustment of  $SC_i$  Raman crystallinity  $I_C^{RS}$  was monitored to study the phase mixture of the initial growth phase for the observed optimum of the device efficiency. In-situ Raman spectra are shown in Fig. 7 for  $SC_i = 1.02\%$  (constant process settings with low initial  $I_C^{RS}$ ) and  $SC = 0.83\%$  ( $SC_i$  with local maximum of  $\eta$ ). The spectra were recorded at  $t_i = 180 \text{ s}$ , for which the deposited layers exhibit thicknesses of  $\sim 45\text{--}55 \text{ nm}$ . For  $SC_i = 1.02\%$  the Raman crystallinity is estimated to  $I_C^{RS} = 50\%$ . For  $SC_i = 0.83\%$  Raman crystallinity  $I_C^{RS} = 62\%$  is observed. These measurements show that even slight changes of the initial  $I_C^{RS}$  can be monitored by in-situ Raman spectroscopy and a well-defined adaption of the Raman crystallinity is enabled by changing the external process settings.

Due to the higher  $I_C^{RS}$  during the initial growth phase ( $SC = 0.83\%$ ), increased defect density of the absorber layer material could be the reason for decreasing  $V_{oc}$  in Fig. 5(b).<sup>16</sup> However, slightly increased  $FF$  for solar cells deposited with  $SC = 0.83\%$  indicates moderately decreasing recombination loss in the absorber layer (Fig. 5). Therefore, we suppose that decreasing  $V_{oc}$  is not caused by higher defect density of the absorber layer. In addition, moderately

reduced charge carrier recombination probability cannot be responsible for the strongly enhanced  $EQE$  only (Fig. 6).

Alternatively, the hydrogen rich plasma at the beginning of the absorber layer growth could be an explanation for decreasing  $V_{oc}$  and increasing  $EQE$  in the short wavelength range. Several effects are conceivable: The growth rate at the very beginning of the intrinsic layer deposition is not known exactly (Fig. 4). Consequently, a short etching period during the very beginning of the absorber layer growth cannot be excluded. In addition the Raman crystallinity of the subsurface layers (p-layer) can be increased.<sup>46–48</sup> Since the absorption coefficient of  $\mu\text{c-Si:H}$  depends on  $I_C^{RS}$ ,<sup>25,49</sup> the transparency of the p-layer might be enhanced. In both cases, an increase of the charge carrier generation could cause the improved  $EQE$  for initially low  $SC_i$ . Furthermore, the passivation of the doping atoms by hydrogen radicals<sup>50,51</sup> could lead to a reduced built-in voltage of the solar cell.

To study the possible impact of a hydrogen plasma on the p-layer properties, a separate test experiment was performed: 30 nm p-type layers were treated with pure hydrogen plasma for a varying treatment time  $t_{H\text{plasma}}$  of 50 s, 100 s, and 200 s. Open circuit voltage  $V_{oc}$ ,  $FF$  and  $EQE$  of p-i-n

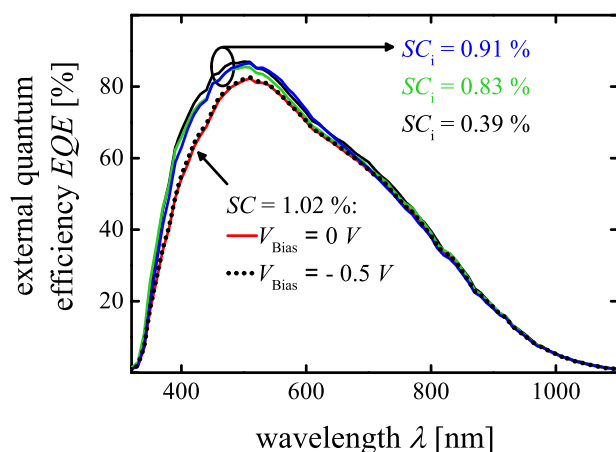


FIG. 6. External quantum efficiency  $EQE$  for solar cells deposited with decreasing initial silane concentration  $SC_i$  ( $t \leq 180 \text{ s}$ ,  $SC(t > 180 \text{ s}) = 1.04\%$ ) and for constant  $SC = 1.02\%$  in dependence of the wavelength  $\lambda$ . For  $SC_i = 1.02\%$   $EQE$  measurements with and without reverse bias voltage  $V_{Bias}$  are depicted.

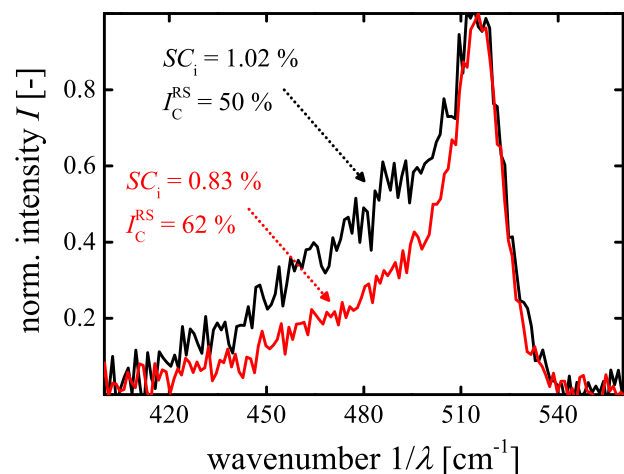


FIG. 7. Process control to increase the Raman crystallinity  $I_C^{RS}$  for the initial growth phase of the absorber layer deposition. The single in-situ Raman spectrum for 50 nm thickness of the second highest crystallinity profile in Fig. 2 is shown ( $SC = 1.02\%$ ). To increase  $I_C^{RS} = 50\%$  the initial silane concentration was decreased to  $SC_i = 0.83\%$  and  $I_C^{RS} = 62\%$  was achieved.

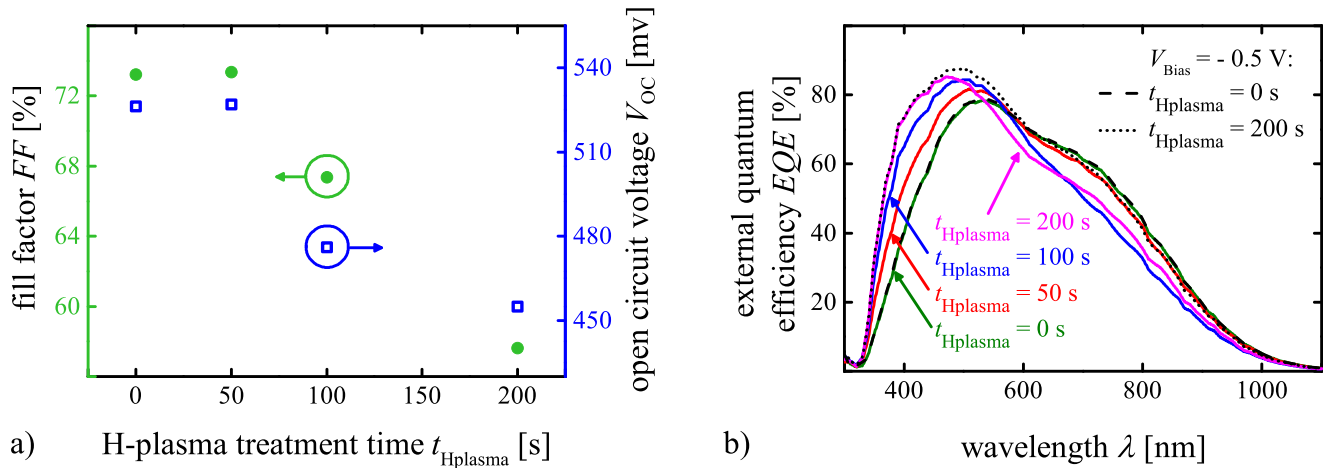


FIG. 8. (a) Open circuit voltage  $V_{OC}$  and fill factor  $FF$  as a function of the hydrogen plasma treatment time  $t_{Hplasma}$ . Maximum absolute measurement errors of the determined photovoltaic parameters are  $\Delta FF = \pm 0.5\%$  and  $\Delta V_{OC} = \pm 5$  mV. For the deposition of the i-layers a constant  $SC = 1.4\%$  was used. (b) External quantum efficiency  $EQE$  for a series of solar cells, whose p-layers were treated by hydrogen plasma, as a function of the wavelength  $\lambda$ .

solar cells deposited on the pretreated p-layers are shown in Fig. 8 (i-layer deposition:  $SC = 1.4\%$ ). For a short treatment time of 50 s,  $FF$  and  $V_{OC}$  are not affected. However, for longer treatment times of the p-type layer with hydrogen,  $FF$  and  $V_{OC}$  are significantly reduced. External quantum efficiency  $EQE$  decreases in the long wavelength range with increasing  $t_{Hplasma}$  (Fig. 8(b)). In contrast,  $EQE$  is strongly enhanced in the short wavelength range for increasing  $t_{Hplasma}$ .<sup>52,53</sup> This effect already appears for  $t_{Hplasma} = 50$  s, although  $EQE$  of the longer wavelength range is not significantly decreased.

The loss of  $EQE$  for long wavelengths can be explained by the lowered built-in voltage of the solar cells because  $EQE$  increases again with reverse bias-voltage (Fig. 8(b) for  $t_{Hplasma} = 200$  s).  $EQE$  with reverse bias voltage is not increased in the short wavelength range (Fig. 8(b):  $t_{Hplasma} = 0$  s and  $t_{Hplasma} = 200$  s). Consequently, improved  $EQE$  for extended  $t_{Hplasma}$  is most probably caused by the optical modification of the pretreated p-layers.

In summary,  $EQE$  is increased in the short wavelength range for reduced  $SC_i$  from 1.02% to 0.91% and saturates for  $SC_i < 0.91\%$  (Fig. 6).  $EQE$  of the solar cell deposited with

constant  $SC = 1.02\%$  shows no improvement by applying reverse bias voltage  $V_{Bias}$ . This indicates that the increase of  $EQE$  for decreasing  $SC_i$  is related to changes of the optical performance of the solar cell (Fig. 6). Optical modifications were observed for solar cells with p-layers that were treated by hydrogen plasma (Fig. 8(b)). Although the outcome of this experiment cannot explain the saturation effect for decreased  $SC_i$ , it provides a qualitative explanation for the increase of  $EQE$  in the blue wavelength range (Fig. 6) by modifications of the p-layer transparency.

### C. Effects of modified Raman-crystallinity during the later layer growth

The initial growth phase of the intrinsic  $\mu c$ -Si:H layer was optimized resulting in a  $I_C^{RS}$  of  $\sim 60\%$  (Fig. 7). To optimize the growth in region B (Fig. 2),  $SC$  was changed continuously during the deposition to compensate increasing  $I_C^{RS}$ . In Fig. 9 crystallinity profiles are shown for varying  $\Delta SC$ . The adaption of the initial  $SC$  results in a constant  $I_C^{RS}$  at  $\sim 60\%$  for the first 400 nm of the i-layer deposition (Fig. 9). For small or negative  $\Delta SC$  the Raman crystallinity increases during the deposition (I:  $\Delta SC = -0.11\%/h$ ). For large positive  $\Delta SC$ ,  $I_C^{RS}$  decreases significantly (III:  $\Delta SC = 0.11\%/h$ ). An almost constant crystallinity profile is found for  $\Delta SC = 0.04\%/h$  (II). Only at the end of the 1300 nm thick i-layer the Raman crystallinity decreases slightly below 60%.

To correlate  $I_C^{RS}$  with device performance  $\eta$ ,  $J_{SC}$ ,  $FF$  and  $V_{OC}$  are studied (Fig. 10). Maximum solar cell conversion efficiencies of  $\eta \approx (7.7 \pm 0.2)\%$  are observed between  $\Delta SC = -0.06\%/h$  and  $\Delta SC = 0.07\%/h$  (Fig. 10(a)). Within this range of  $\Delta SC$  Raman crystallinities vary between approximately  $I_C^{RS} = 60\%$  and  $I_C^{RS} = 70\%$  during the later absorber layer growth (Fig. 9). Since this includes constant  $\Delta SC$  readjustment of the process settings in region B seems to be dispensable. To demonstrate the importance of the controlled deposition photovoltaic parameters  $J_{SC}$ ,  $FF$ , and  $V_{OC}$  are discussed in detail.

Highest  $J_{SC} \approx 20.8 \text{ mA/cm}^2$  were achieved for  $\Delta SC = -0.06\%/h$  and  $\Delta SC = 0\%/h$ , respectively (Fig. 10(a)). Slightly decreasing  $J_{SC}$  is detected for  $\Delta SC < -0.06\%/h$ ,

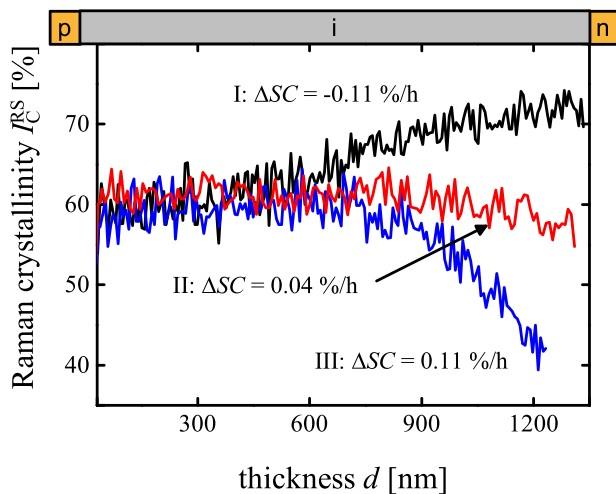


FIG. 9. Raman crystallinity  $I_C^{RS}$  as a function of thickness  $d$  for three different slopes of the silane concentration  $\Delta SC$ .

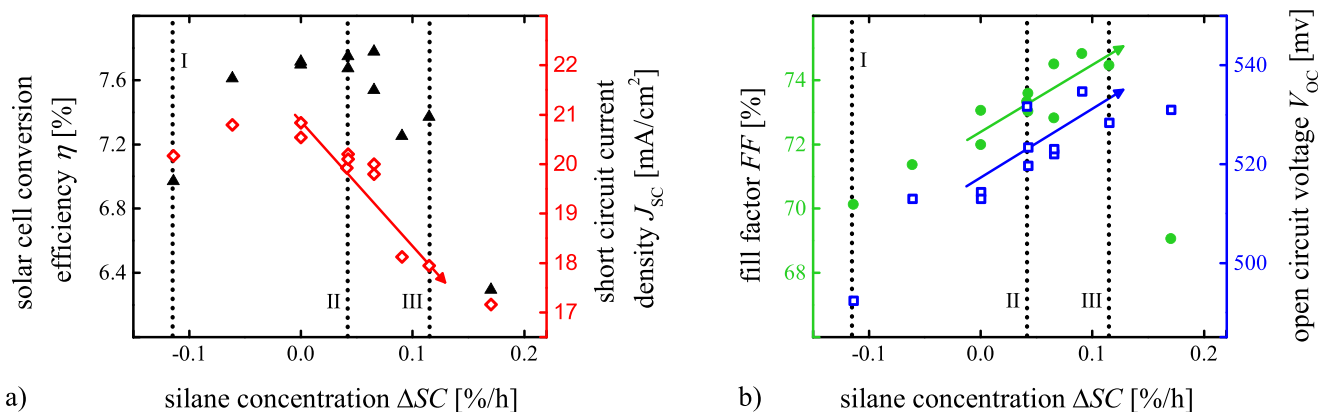


FIG. 10. Solar cell conversion efficiency  $\eta$  and short circuit current density  $J_{SC}$  in dependence of the silane concentration  $\Delta SC$ . For three selected  $\Delta SC$ , crystallinity profiles of the i-layer are shown in Fig. 9. (b) Fill factor  $FF$  and open circuit voltage  $V_{OC}$  as a function of  $\Delta SC$ . Maximum absolute measurement errors of the determined photovoltaic parameters are  $\Delta\eta = \pm 0.2\%$ ,  $\Delta J_{SC} = \pm 0.5 \text{ mA/cm}^2$ ,  $\Delta FF = \pm 0.5\%$  and  $\Delta V_{OC} = \pm 5 \text{ mV}$ .

whereas  $J_{SC}$  is reduced significantly for  $\Delta SC > 0\%/h$ . Strongly reduced  $J_{SC}$  between  $\Delta SC = 0\%/h$  and  $\Delta SC = 0.11\%/h$  is mainly caused by decreasing  $I_C^{RS}$  (Fig. 9) and decreased absorption coefficient of the intrinsic layer.<sup>54</sup> In contrast, decreasing  $J_{SC}$  for  $\Delta SC < -0.06\%/h$  cannot be explained by the enhanced absorption coefficient of the intrinsic layer.<sup>54</sup>

Fill factor  $FF$  is improved from  $FF \approx 70\%$  to  $FF \approx 75\%$  for increasing  $\Delta SC$  from  $-0.11\%/h$  to  $0.11\%/h$  (Fig. 10(b)). For  $\Delta SC = 0.17\%$   $FF$  is significantly reduced to  $FF \approx 69\%$ . Open circuit voltage  $V_{OC}$  is enhanced over the entire studied range of  $\Delta SC$  from about  $490 \text{ mV}$  to roughly  $530 \text{ mV}$ . Consequently, decreasing  $J_{SC}$  for  $\Delta SC < -0.06\%/h$  (Fig. 10(a)) is most likely caused by the lowered  $FF$ .

As a result, maximum  $FF$  resp.  $V_{OC}$  and  $J_{SC}$  are found for different process settings and different Raman crystallinity profiles of the absorber layer. Between profile I ( $\Delta SC = -0.11\%/h$ ) and II ( $\Delta SC = 0.04\%/h$ ), evolution of the Raman crystallinity is favorable for  $J_{SC}$ . For Raman crystallinities between profile II ( $\Delta SC = 0.04\%/h$ ) and III ( $\Delta SC = 0.11\%/h$ ),  $FF$  resp.  $V_{OC}$  are improved. Compared to current record  $\mu\text{c-Si:H}$  single junction solar cells,<sup>18,36</sup> fill factor and  $V_{OC}$  are quite similar ( $0.04\%/h \leq \Delta SC \leq 0.11\%/h$ ), whereas the short circuit current density of these cells is superior to our device.

For constant Raman crystallinity profiles at approximately  $I_C^{RS} = 60\%$ , solar cell conversion efficiency is improved. However, since the solar cell parameters  $FF$ ,  $V_{OC}$  and  $J_{SC}$  show maxima for different crystallinity profiles, ideal evolution of the Raman crystallinity remains unclear. Further improvement of the device performance might be achieved by a fine-tuning of the process settings. Monitoring the Raman crystallinity in-situ provides sufficient information to adjust the external process settings tailored to the requirements and enables the reproduction of high quality absorber layers.

#### IV. SUMMARY

In this paper, we investigated the effect of controlled absorber layer growth on the solar cell device performance. For this purpose, the Raman crystallinity was monitored during the absorber layer growth in-situ and growth drifts were compensated by readjusting the process settings.

First, crystallinity profiles deposited under constant process settings were studied in detail. Low Raman crystallinity at the initial growth phase and continuous drifts for the later layer deposition were detected.

To adapt the Raman crystallinity during the initial absorber layer growth, decreased initial silane concentration was applied. Due to the adjustment of the process settings the external quantum efficiency in the short wavelength range was enhanced and the solar cell conversion efficiency was improved by  $0.41\%$  absolute. For the compensation of the continuous increase of  $I_C^{RS}$  during the later layer growth, the silane concentration was gradually increased. Accordingly, a constant Raman crystallinity profile was achieved with  $I_C^{RS} \approx 60\%$  throughout the entire absorber layer. Although solar cells with those absorber layers exhibit improved device efficiency, optimum photovoltaic parameters  $J_{SC}$ ,  $FF$  and  $V_{OC}$  were obtained for different controlled Raman crystallinity profiles.

In conclusion, we demonstrated the controlled deposition of intrinsic microcrystalline absorber layers by monitoring the evolution of the Raman crystallinity in-situ, adjusting the process settings and observing the impact on the solar cell device parameters. Using this method, process drifts were compensated and beneficial crystallinity profiles were identified. Most likely, optimized crystallinity profiles as well as inherent process drifts vary dependent on the device configuration. Nevertheless, the presented iterative optimization procedure between the monitoring of the Raman crystallinity and the readjustment of the process settings is universally applicable.

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