

Self-consistent calculations within the Extended Theory of Finite Fermi Systems

A.Avdeenkov,^{1,2} F.Grümmer,¹ S.Kamerdzhev,² S.Krewald,¹ N.Lyutorovich,^{1,3} and J.Speth^{1,4}

¹*Institut für Kernphysik, Forschungszentrum Jülich, 52425 Jülich, Germany*

²*Institute of Physics and Power Engineering, 249020 Obninsk, Russia*

³*Institute of Physics S.Petersburg University, Russia*

⁴*Institute of Nuclear Physics, PAN, PL-31-342 Cracow, Poland*

The *Extended Theory of Finite Fermi Systems*(ETFFS) describes nuclear excitations considering phonons and pairing degrees of freedom, using the effective Landau-Migdal interaction and nuclear mean fields obtained from experimental data. Here we employ the nuclear mean field derived from Skyrme interactions and the corresponding particle hole interaction. This allows to extend the range of applicability of the ETFFS to experimentally not yet investigated short-lived isotopes. We find that Skyrme interactions which reproduce at the mean field level both ground state properties and nuclear excitations are able to describe the spreading widths of the giant resonances in the new approach, but produce shifts of the centroid energies. A renormalization of the Skyrme interactions is required for approaches going beyond the mean field level.

I. INTRODUCTION

The renewed interest in nuclear structure physics is due to its application to astrophysics. Theories for the element synthesis in stars need both bulk properties of unstable isotopes like masses, and more specific information, such as level densities and neutrino-nucleus cross sections. Eventually, the compression modulus and the symmetry energy of neutron-rich nuclei will be determined. For a recent review, see e.g. Ref. [1].

Most calculations in this field are restricted to the mean field approximation which describes excited states by the one-particle one-hole random phase approximation (1p1h RPA). In this approximation one is able to describe two integral characteristics of giant resonances: the mean energy and the integrated strength. The width of the giant resonances is underestimated in mean field approaches, however, even if one includes the single particle continuum.

In unstable nuclei, the analysis of, for example, the breathing mode and the giant electric dipole resonance requires realistic strength distributions because these resonances are expected to be broad and various multipole resonances overlap one another. Theories beyond the mean field level are called for.

The *Extended Theory of Finite Fermi Systems* (ETFFS) is a linear response theory which incorporates phonons and pairing degrees of freedom. One essentially replaces the particle-hole propagator, which leads to the random phase approximation, by a more general propagator and changes, correspondingly, the renormalized interaction, the effective operators and the ground-state correlations. Details can be found in Refs. [3, 4, 5, 6], where the authors worked out the coupling of phonons to the single particle propagators and to the interaction in a consistent way. The present calculations are based on the *Quasi particle Time Blocking Approximation* (QTBA) that also considers pairing correlations [7]. Examples and references to previous calculations and microscopic analyses of experimental data can be found in

[8, 9]. Another model, the QRPA plus phonon coupling model, has been developed by Bortignon and coworkers [10].

Most excited states predicted by the conventional random-phase approximation depend sensitively on the level density of the single-particle spectrum, as is well known from the schematic model by Brown and Bolsterli [11]. As the TFFS, with and without pairing, uses the experimental single particle spectrum as far as it is available, e.g. [12, 13], it generally comes closer to the experimental data than self-consistent approaches do. For that reason an extended version of the TFFS in which the effects of the phonons are included is the appropriate tool for systematic investigations of both collective and non collective states [14]. Its applicability, however, is limited by the need for experimental information on the level density in the neighboring odd mass nuclei.

At present, these properties of most short-lived neutron rich nuclei are not known experimentally. Therefore one is forced to use nuclear mean fields extrapolated from the known stable nuclei. One of the most successful approaches starts with Skyrme forces, a class of *effective interactions* defined at the *mean field level*. When going beyond the mean field level, one has to expect a renormalization of the interaction strengths [15].

In the present communication, we investigate quantitatively to which extent a renormalization of the Skyrme parameterization may be required by employing a self-consistent mean field and the corresponding particle hole interaction derived from existing Skyrme forces in the framework of the Extended Theory of Finite Fermi systems.

There are various Skyrme parameterizations that reproduce the ground state properties of the known isotopes with similar accuracy but result in different mean fields, as reviewed recently in Refs. [16, 17].

The nuclear matter properties of the Skyrme parameterizations, such as the effective mass, the compression modulus, and the symmetry energy can be used to group the effective interactions into some broad classes. The versions initiated by Vautherin and Brink were fitted to

ground state properties only. Those interactions are characterized by a large incompressibility (see SIII in Table I). Later versions of the Skyrme interactions modified the

TABLE I: Nuclear matter properties of the Skyrme parameterizations used in the present calculation; for comparison also SIII is given

	m^*/m	$\kappa(MeV)$	$a_s[MeV]$
BSk5	0.92	237	28.7
SLy4	0.70	230	32.0
SKM*	0.79	218	30
SIII	0.76	356	28.2

density dependence of the interaction in order to allow a smaller incompressibility. Those interactions reproduce the mean excitation energies of giant resonances reasonably well and are therefore a natural starting point for our investigation.

The interaction SLy4 [19] has been used in(Q)RPA [20, 21, 22] and in some first attempts to include explicitly long range quadrupole correlations (*alias quadrupole phonons*) into the nuclear ground state [23]. An investigation similar to the present one, but based on the relativistic mean field approach, has been published recently [24, 25]. Their results, calculated with and without phonons differ very little from each other, in contrast to our results. The reason is the so-called subtraction procedure of Ref. [7] which these authors apply. This is an *ad hoc* condition that is difficult to justify, as one needs a consistent description of both even even and even odd nuclei.

II. METHOD

We calculate the mean field within the HFB approach using the code HFBRAD [26] and the excited states with the the QTBA version [7] of the ETFFS. Here, as mentioned before, the authors developed a renormalization procedure that corrects for possible double counting of the phonon contributions that are already included in the Skyrme parameterization. We leave this correction out. This is also in line with the investigations by Bertsch et al. [23], who argue that an effective interaction has to be refitted if one includes additional correlations. As usual, the residual interaction for (Q)RPA and the present theory is derived as the second derivative of the Skyrme energy functional [27, 28]. In the present calculation we make some simplifications. For the non magic nuclei we solved the HFB equations and extracted the quasi particle energies, wave functions and the occupation numbers. In the second derivative, the ph-interaction, we dropped the spin-orbit term and we approximated the velocity-dependent terms by its Landau-Migdal limit [27, 28]. There are two kinds of velocity-dependent terms: the first is $\propto \mathbf{k}^2 \delta(\mathbf{r} - \mathbf{r}')$ and the second one is $\propto \mathbf{k}^\dagger \delta(\mathbf{r} - \mathbf{r}') \mathbf{k}$

(p-wave interaction in momentum space). The value of the first term averaged over the angular variables gives $\frac{1}{2} k_F^2 \delta(r - r')$, while that of the second one vanishes. Such an approximation, of course, violates self-consistency and one has to modify the parameters of the residual interaction to obtain the spurious center-of-mass state at zero energy. For this reason we changed the parameters proportional to $t_1 k_F^2 \delta(r - r')$ up to 20%.

In principle, the ETFFS treats the single-particle continuum on the RPA level for magic nuclei exactly and includes the ground state correlations caused by CC (GSC_{phon}) [8] explicitly. However, because of technical difficulties caused by the pairing correlations, we did not consider these effects in the present calculation. The ground state correlations are therefore restricted to the conventional RPA correlations. We discretized the continuum with a quasi-particle energy cutoff of 100 MeV. We did not find any noticeable differences in the results with a larger basis. We used 14-16 low-lying phonons with angular momentum $L = 2 - 6$ and natural parity.

III. RESULTS

A. Collective isovector electric dipole states

Here we present our results for the giant dipole excitation for some of the tin isotopes from ^{100}Sn up to ^{132}Sn , where we used the SLy4 [19] parameterization of the Skyrme force. The SLy4 [19] parameterization is not only adjusted to ground state properties, but it also reproduces within the RPA the GDR and isoscalar collective states reasonably well. For comparison we used for ^{132}Sn two further parameterizations, namely BSk5 [29] and SKM* [30]. We compare our theoretical results with the data of ^{120}Sn and ^{132}Sn . In general, the experimental GDR data are analysed by fitting to a Lorentzian. For that reason we have used the same procedure to determine the calculated integral characteristics [5, 31]. The mean energies E_0 and the resonance widths Γ have been extracted from the calculated strength distribution under the condition that the first three energy-weighted moments of the Lorentzian and the theoretical distribution of the total photo absorption cross section should coincide. This method works well in the case of the ETFFS results.

In Fig. 1 the photo absorption cross sections for ^{120}Sn and ^{132}Sn are compared with the experimental ones. The full line represents the result of the present approach. The dotted line is a Lorentz parametrization of the experimental data. In both cases the shape of the calculated cross sections agrees very well with the experimental ones. The energies are in both cases, however, too low. Quantitatively, we give the same results in Table II and compare it with the RPA results. In Fig.2 the experimental cross section of ^{132}Sn (dotted line) is compared with the results of the RPA and the present theory where we used three different Skyrme parameterizations. The

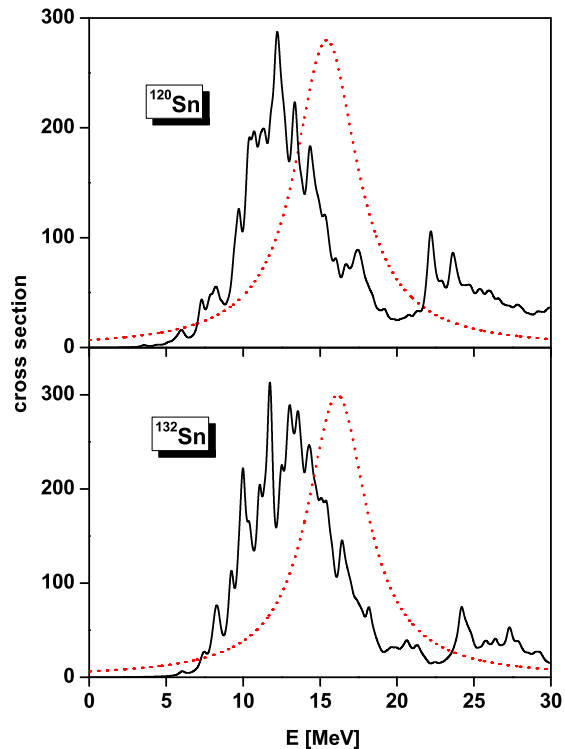


FIG. 1: E1 photo absorption cross section for ^{120}Sn and ^{132}Sn calculated within the present theory (full). The dotted lines indicate the data for ^{120}Sn [32] and ^{132}Sn [33]. The smearing parameter is $\Delta = 0.2$ MeV. The SLy4 Skyrme parametrization was used.

theoretical results look similar but deviate in detail. In all three cases the energies calculated within the ETFFS approach are too low compared to the data [33], but the strength distribution is strongly improved compared to the RPA calculation. In Table II we give a more detail comparison. The BSk5 and SLy4 parameterizations give very similar results; the widths are roughly 50% larger compared to the one calculated in RPA and are in good agreement with the data. The energies of the extended theory are more than 1.5 MeV lower than the RPA results which are already 2 MeV too low compared to the experiment.

Both effects are connected with the inclusion of the phonons. (I) The single particle strength is fragmented due to the coupling to the phonons and this effect gives rise to the spreading width which strongly improves the agreement with the experiments. (II) The phonons, however, also compress the single particle spectrum and for that reason the energies calculated in ETFFS are lower by about 1.5 MeV compared with the RPA. The SKM* result differs from the first two parameterizations. Already on the RPA level the strength is less fragmented compared to the previous ones. If phonons are included a fraction of the narrow RPA peak survives and gives rise to

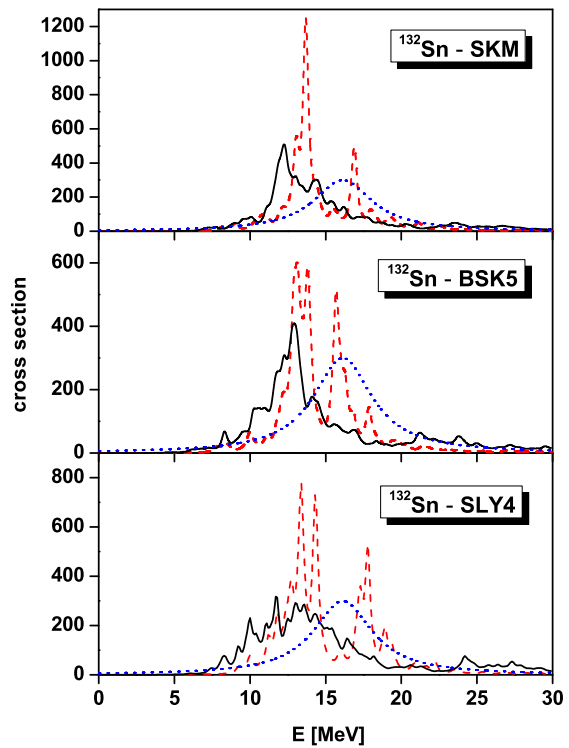


FIG. 2: E1 photo absorption cross section for ^{132}Sn calculated within the present theory with different Skyrme forces. The smearing parameter is $\Delta = 0.2$ MeV. The dotted line are the data of ref. [33].

a much smaller width than in the previous calculations. The method to approximate the strength distribution by one single Lorentzian is not well justified in the RPA case because the RPA response fragments in two clearly separated peaks. If one uses this procedure, one obtains the average of the two peaks and a Lorentzian width which happens to be large. If the analysis includes the main peak only, the RPA width reduces to 1 MeV. In Fig.

TABLE II: Mean energy and width of the GDR in ^{120}Sn and ^{132}Sn calculated within the RPA and the ETFFS. The SLy4 Skyrme parametrization was used. The E1 strength is fitted to a Lorentzian in the range of (8 - 22) MeV

	$E_0(\text{MeV})$	$\Gamma(\text{MeV})$
$^{120}\text{Sn}(\text{ETFFS})$	12.3	5.2
(RPA)	15.1	3.6
(Exp)	15.4	4.9
$^{132}\text{Sn}(\text{ETFFS})$	12.5	5.2
(RPA)	14.3	3.8
(Exp)	16.1	4.7

3 the E1 strength distributions for ^{100}Sn , ^{110}Sn , ^{120}Sn and ^{132}Sn are shown. In Table III we compare (Q)RPA results with the ones obtained with the present theory.

In all cases the energy of the present theory is shifted to lower energies by about 2 MeV, independent of the number of neutrons, and the widths are increased by 50%. As the widths are increased the maximum of the cross section is decreased up to 50%. The strength in the two closed-shell nuclei ^{100}Sn and ^{132}Sn has about the same width, despite the fact that one has added 32 neutrons. The two closed-shell nuclei show, as expected, a compact and relatively narrow GDR compared to the two open-shell nuclei, ^{110}Sn and ^{120}Sn .

TABLE III: Integral characteristics of the GDR in ^{132}Sn calculated within the RPA and with inclusion of phonons with three different Skyrme forces. The E1 strength is fitted to a Lorentzian in the (8 - 22) MeV interval. Data are from Ref. [33]

	$E_0(\text{MeV})$	$\Gamma(\text{MeV})$
Bsk5(ETFFS)	12.6	4.4
Bsk5(RPA)	14.0	2.9
SLy4(ETFFS)	12.5	5.3
SLy4(RPA)	14.3	3.8
SKM*(ETFFS)	13.1	3.4
SKM*(RPA)	14.0	3.3
exp.	16.1(7)	4.7

We also mention that for ^{132}Sn we solved the continuum RPA. The nearly complete overlap with the discretized RPA shows that our discretization procedure, together with the smearing parameter $\Delta = 200$ keV, is well justified for the present energy range. As the two distributions are hard to distinguish, we do not show the continuum result in Fig. 3.

Our results of ^{120}Sn and ^{132}Sn for the GDR is qualitatively similar to the calculations by Bortignon et al. [35]. They use for their comparison the SIII parametrization which gives on the RPA level a centroid energy about 2 MeV higher than the SLy4 parametrization. Therefore SIII has the right tendency.

However, the far too high incompressibility (see Table I) was one of the motivations for the new generation of parameterizations with a different density dependence ($\alpha < 1$), which we use in our calculations. Note: there is a difference between the definition of the mean energy of the total E1 distribution $E_0 = m_1/m_0$ and the mean energy of the GDR which is used in Table III. With the first definition the RPA and ETFFS results are very similar e.g. for ^{132}Sn $E_0 = 14.74\text{MeV}$ (RPA) and $E_0 = 14.47\text{MeV}$ (ETFFS), whereas the mean energy of the GDR (Table III) differ up to 1.8 MeV. If one compares the theoretical results with the data one has to use the latter definition because the experimentalists extract their results in the same way from their data.

Finally we compare in Fig. 4 the original theoretical result for the GDR in ^{120}Sn and ^{132}Sn with the result where the isovector force was increased by 20%. Here we obtain good agreement with the data. The structures beyond 20 MeV are due to non collective ph components

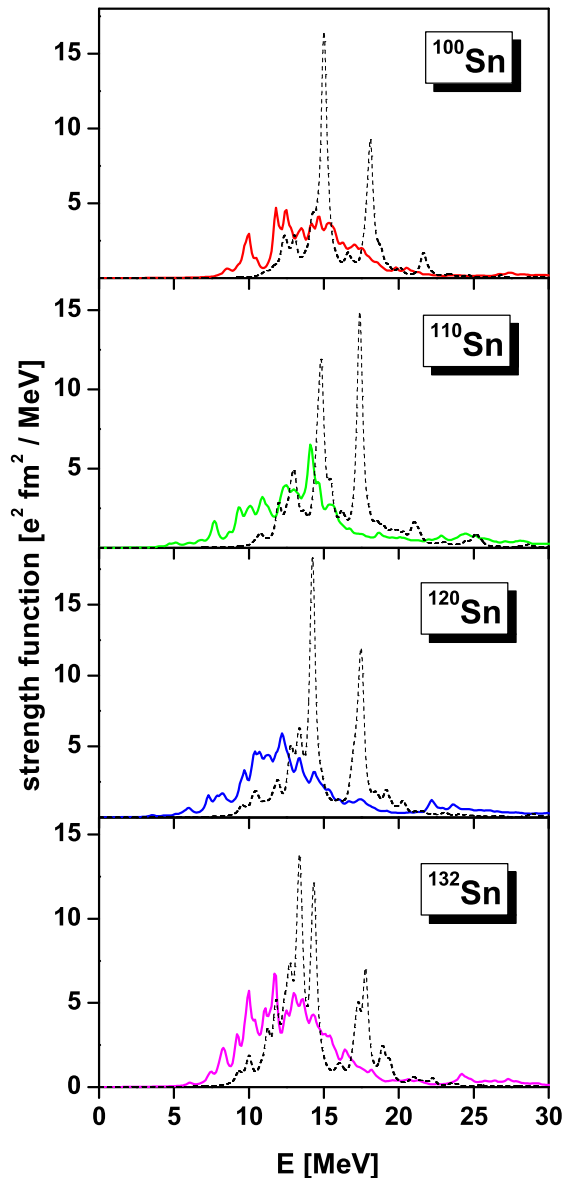


FIG. 3: Strength functions for GDR in ^{100}Sn , ^{110}Sn , ^{120}Sn and ^{132}Sn . The (Q)RPA results (dashed lines) are compared with the results of the ETFFS. The SLy4 Skyrme parametrization was used.

which will be smeared out if we include the single particle continuum.

B. Collective electric isoscalar states

Energetically low-lying collective states, such as the 3^- mode, are extremely sensitive to the strength of the isoscalar particle hole interaction. In Fig. 5 the E2 strength functions are presented for ^{120}Sn and ^{132}Sn . We show the results of the RPA and the present theory. The

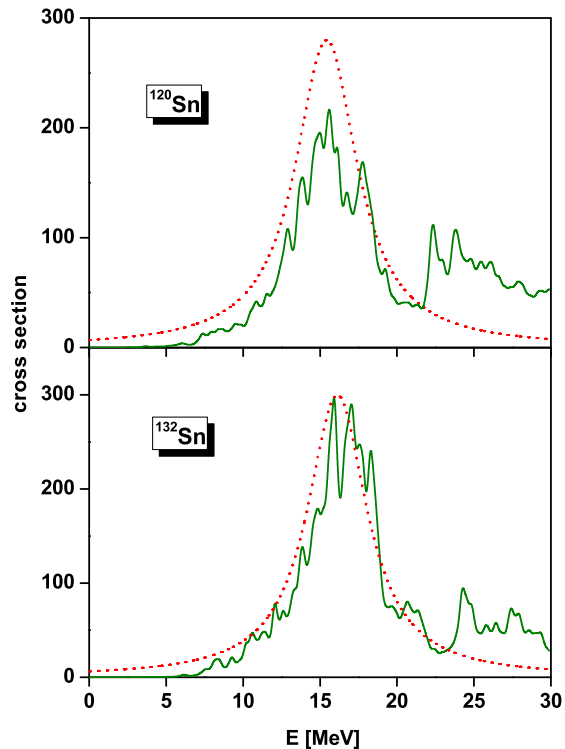


FIG. 4: E1 photo absorption cross section for ^{120}Sn and ^{132}Sn calculated within the present theory where the isovector force has been increased by 20% (full). The dotted lines indicate the Lorentzian fit to the experimental data [32, 33].

lowest 2^+ in ^{120}Sn is below zero and, as well, the isoscalar giant electric quadrupole resonance is too low compared with the data. In ^{132}Sn the isoscalar electric quadrupole resonance has not yet been measured. In Table IV the theoretical energies of the low lying collective 2^+ and 3^- states are compared with the experiments. The RPA results for the 2^+ modes are in nice agreement with the data for both isotopes. The 3^- mode in ^{132}Sn is underestimated by 1.35 MeV. In the presence of phonons, the collective 3^- modes are pulled down by approximately 1.5 MeV. The 2^+ mode in ^{132}Sn is shifted downwards by 2.36 MeV, and the 2^+ mode in ^{120}Sn is even shifted below the ground state energy. Here, however, one has to bear in mind that in ^{120}Sn the magnitude of the pairing gaps plays an important role.

The reduction of the excitation energies of the collective isoscalar modes with respect to the random phase approximation has two sources, the modified level density in the vicinity of the Fermi level, and the strength of the isoscalar particle hole interaction. The compression of the level density due to phonons can be understood qualitatively by second order perturbation theory: The intermediate states generated by the phonon degree of freedom have larger excitation energies than the single particle level obtained in the mean field approximation.

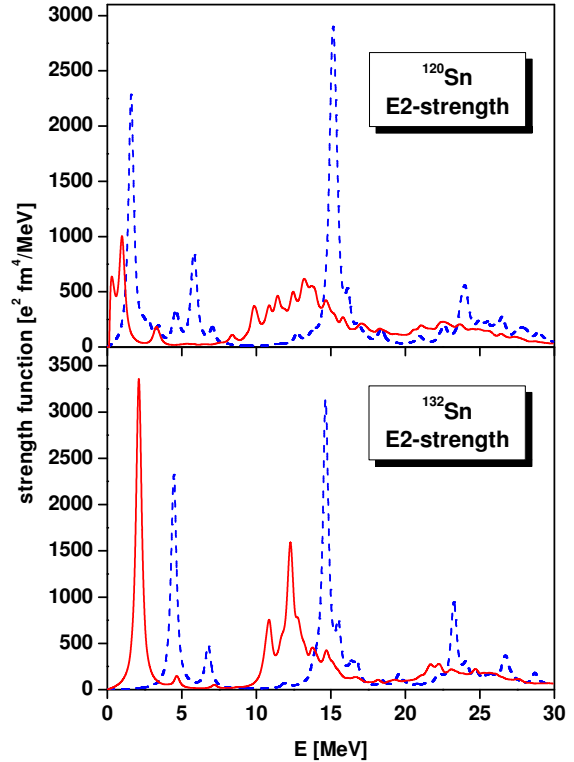


FIG. 5: The strength function for 2^+ excitation for ^{120}Sn and ^{132}Sn calculated with (full) and without (dashed) Phonons. The smearing parameter is $\Delta = 0.2$ MeV. The SLy4 Skyrme parametrization was used.

Therefore phonons introduce an additional attractive interaction, whatever the details of the phonon-particle coupling are, which increases the level density near the Fermi level. The only way to compensate the reduced level spacing between particle and hole states is to reduce the strength of the isoscalar particle hole interaction in comparison with the interaction used in a random phase calculation.

Similar arguments hold for the collective isovector modes discussed above. The repulsive isovector parts of the particle hole interaction have to be increased compared with the strength used in the random phase approximation in order to compensate the phonon effects. This conclusion is independent on details of the theoretical framework.

IV. CONCLUSION

In the present paper we report self-consistent nuclear structure calculations performed in the Extended Theory of Finite Fermi Systems. Skyrme parameterizations were used, which reproduce not only ground state properties of nuclei but also excited collective states within the RPA. In comparison with RPA calculations, the widths of the

giant resonances are brought into better agreement with the experimental data. The centroid energies of the resonances are shifted, however. This finding implies the necessity of a renormalization of Skyrme interactions for applications beyond the mean field level. Our investigation suggests two major renormalizations: a reduction of the strength of the isoscalar particle-hole interaction and an increase of the isovector strength.

V. ACKNOWLEDGMENT

We thank Victor Tselyaev for many discussion concerning the QTBA. We also thank Sergey Tolokonnikov for useful comments on the energy-density functional. One of us (JS) thanks Stanislaw Drożdż for many discussions and the Foundation for Polish Science for financial

support through the *Alexander von Humboldt Honorary Research Fellowship*. The work was partly supported by the DFG and RFBR grants Nos.GZ:432RUS113/806/0-1 and 05-02-04005 and by the INTAS grand No.03-54-6545.

TABLE IV: Low-lying collective 2^+ and 3^- states in ^{120}Sn and ^{132}Sn calculated within and without phonons. The SLy4 Skyrme parametrization was used.

	^{120}Sn		^{132}Sn	
	2^+ [MeV]	3^- [MeV]	2^+ [MeV]	3^- [MeV]
without phonons	1.60	2.60	4.46	3.00
with phonons	—	1.3	2.10	1.50
exp.	1.17	2.40	4.04	4.35

-
- [1] F.K. Thielemann *et al.*, Prog. Part. Nucl. Phys. 46 (2001) 5.
- [2] S. Drożdż, V. Klemm, J. Speth and W. Wambach, Phys. Rep. 197 (1990) 1.
- [3] V.I. Tselyaev, Sov. J. Nucl. Phys. 50 (1989) 780.
- [4] S.P. Kamerdzhiev and V.I. Tselyaev, Bull. Acad. Sci. USSR, Phys. Ser. 55 (1991) 45.
- [5] S.P. Kamerdzhiev, J. Speth, G. Tertychny and V. Tselyaev, Nucl. Phys. A 555 (1993) 90.
- [6] S.P. Kamerdzhiev, J. Speth, G. Tertychny and J. Wambach, Z. Physik A 336 (1993) 253.
- [7] V.I. Tselyaev, arXiv:nucl-th/0505031; Phys.Rev.C (to be published); E.V. Litvinova and V.I. Tselyaev, arXiv:nucl-th/0512030.
- [8] S. Kamerdzhiev, J. Speth and G. Tertychny, Phys. Rep. 393 (2004) 1.
- [9] F.Grümmer and J.Speth, J.Phys.G: Nucl. Part. Phys. 32 (2006) R193.
- [10] G. Colo, P.F. Bortignon, Nucl. Phys. A 696 (2001) 427.
- [11] G.E.Brown and M. Bolsterli, Phys. Rev. Lett. 3(1959) 472.
- [12] P. Ring and J. Speth, Nucl. Phys. A 235 (1974) 315.
- [13] D. Zawischa, J. Speth and D. Pal, Nucl. Phys. A 311 (1978) 445.
- [14] V. Tselyaev, J. Speth, F. Grümmer, S. Krewald, A. Avdeenkov, E. Litvinova and G. Tertychny, Phys. Rev. C 75 (2007) 014315.
- [15] E. Epelbaum, Prog. Part. Nucl. Phys. 57 (2006) 654
- [16] M. Bender, P.H. Heenen and P.G. Reinhard, Rev. Mod. Phys. 75 (2003) 121.
- [17] G.F. Bertsch, B. Sabbey and M. Uusnakki, Phys. Rev. C 71 (2005) 054311.
- [18] C. Mahaux, P.F. Bortignon, R.A. Broglia and C.H. Dasso, Phys. Rep. 120 (1985) 1.
- [19] E. Chabanat, P. Bonche, P. Haensel, J. Meyer, R. Schaeffer, Nucl. Phys. A 635 (1998) 231.
- [20] J. Terasaki, J. Engel, M. Bender et al., Phys. Rev. C 71 (2005) 034310.
- [21] S. Goriely, E. Khan, M. Samyn, Nucl. Phys. A 739 (2004) 331.
- [22] J. Terasaki, Phys. Rev. C 74 (2006) 044301.
- [23] M. Bender, G.F. Bertsch and P.H. Heenen, Phys. Rev. C 73 (2006) 034322.
- [24] E. Litvinova, P. Ring and D.Vretenar, preprint (January 2007).
- [25] E. Litvinova, P. Ring and V. Tselyaev, preprint (September 2006).
- [26] K. Bennaceur and J. Dobaczewski, Comp. Phys. Comm. 168 (2005) 96.
- [27] S.O. Bäckman, A.D. Jackson, J. Speth, Phys.Lett. B 56 (1975) 209.
- [28] S. Krewald, V. Klemm, J. Speth and A. Faessler, Nucl. Phys. A 281 (1977) 166.
- [29] M. Samyn, S. Goriely, P.-H. Heenen, J.M. Pearson and F. Tondeur, Nucl. Phys. A 700 (2002) 142.
- [30] J. Bartel, P. Quentin, M. Brack, C. Guet and H.-B. Håkansson, Nucl. Phys. A 386 (1982) 79.
- [31] V.I. Tselyaev, Bull. Rus. Acad. Sci. Phys. 64 (2000) 434.
- [32] S.C. Fultz, B.L. Berman, J.T. Caldwell, R.L. Bramblett and M.A. Kelly, Phys. Rev. 186 (1969) 1255.
- [33] P. Adrich et al., Phys. Rev. Lett. 95 (2005) 132501-1.
- [34] E. Litvinova and P. Ring, Phys. Rev. C 73 (2006) 044328.
- [35] D.Sarchi, P.F. Bortignon and G.Colo, Phys.Lett. B601 (2004) 27.