

Determining the scale in Lattice QCD

**V. G. Bornyakov^a, R. Horsley^{*b}, R. Hudspith^c, Y. Nakamura^d, H. Perlt^e, D. Pleiter^f,
P. E. L. Rakow^g, G. Schierholz^h, A. Schiller^e, H. Stübenⁱ and J. M. Zanotti^j**

^a *Institute for High Energy Physics, Protvino, 142281 Protvino, Russia,*

Institute of Theoretical and Experimental Physics, Moscow, 117259 Moscow, Russia,

School of Biomedicine, Far Eastern Federal University, 690950 Vladivostok, Russia

^b *School of Physics and Astronomy, University of Edinburgh, Edinburgh EH9 3FD, UK*

^c *Department of Physics and Astronomy, York University, Toronto, ON Canada M3J 1P3*

^d *RIKEN Advanced Institute for Computational Science, Kobe, Hyogo 650-0047, Japan*

^e *Institut für Theoretische Physik, Universität Leipzig, 04109 Leipzig, Germany*

^f *JSC, Forschungszentrum Jülich, 52425 Jülich, Germany*

Institut für Theoretische Physik, Universität Regensburg, 93040 Regensburg, Germany

^g *Theoretical Physics Division, Department of Mathematical Sciences, University of Liverpool,*

Liverpool L69 3BX, UK

^h *Deutsches Elektronen-Synchrotron DESY, 22603 Hamburg, Germany*

ⁱ *Universität Hamburg, Regionales Rechenzentrum, 20146 Hamburg, Germany*

^j *CSSM, Department of Physics, University of Adelaide, Adelaide SA 5005, Australia*

E-mail: rhorsley@ph.ed.ac.uk

QCDSF-UKQCD Collaborations

We discuss scale setting in the context of 2+1 dynamical fermion simulations where we approach the physical point in the quark mass plane keeping the average quark mass constant. We have simulations at four beta values, and after determining the paths and lattice spacings, we give an estimation of the phenomenological values of various Wilson flow scales.

The 33rd International Symposium on Lattice Field Theory

14 -18 July 2015

Kobe International Conference Center, Kobe, Japan

*Speaker.

1. Singlet quantities

Numerical lattice QCD simulations determine mass (or other) ratios but not the scale itself, which has to be determined from experiment. A hadron mass such as the proton mass or decay constant such as the pion decay constant are often used for this purpose. We discuss here the advantages of setting the scale using a flavour-singlet quantity, which in conjunction with simulations keeping the average quark mass constant allow $SU(3)$ flavour breaking expansions to be used. This is illustrated using $2 + 1$ flavour clover fermions, and in addition a determination of the Wilson flow scales, $\sqrt{t_0^{\text{exp}}}$ and w_0^{exp} is given.

This talk is based on [1], where further details can be found.

Dynamical simulations start with some values of the quark masses and then extrapolate along some path in (u, d, s) space¹ to the physical point. The strategy we have adopted here, [2, 3] is to start at a point on the $SU(3)$ flavour symmetric line, when all the quark masses are equal

$$(m_0, m_0, m_0) \rightarrow (m_u^*, m_d^*, m_s^*), \quad (1.1)$$

and to keep the singlet quark mass \bar{m} constant

$$\bar{m} = \frac{1}{3} (m_u + m_d + m_s) = \text{const.} \equiv m_0. \quad (1.2)$$

This allows an $SU(3)_F$ flavour symmetry breaking expansion for masses and matrix elements. The expansion parameter is naturally the distance from the $SU(3)$ flavour plane, parametrised by

$$\delta m_q = m_q - \bar{m}. \quad (1.3)$$

This has the trivial constraint

$$\delta m_u + \delta m_d + \delta m_s = 0. \quad (1.4)$$

The expansion coefficients are functions of \bar{m} only so provided \bar{m} is kept constant they remain unaltered whether we have mass degenerate u and d quarks or not. This opens the possibility of determining isospin breaking quantities from just $2 + 1$ simulations. The plane (or path) is called ‘unitary’ if we expand in both the same sea and valence quarks masses.

Consider now a flavour singlet quantity $X_S(m_u, m_d, m_s)$ which by definition is invariant under u, d, s permutations. This has a stationary point about the $SU(3)$ flavour symmetric line. For upon expanding a flavour singlet quantity about a point on the $SU(3)$ -flavour line we have

$$\begin{aligned} & X_S(\bar{m} + \delta m_u, \bar{m} + \delta m_d, \bar{m} + \delta m_s) \\ &= X_S(\bar{m}, \bar{m}, \bar{m}) + \left. \frac{\partial X_S}{\partial m_u} \right|_0 \delta m_u + \left. \frac{\partial X_S}{\partial m_d} \right|_0 \delta m_d + \left. \frac{\partial X_S}{\partial m_s} \right|_0 \delta m_s + O((\delta m_q)^2). \end{aligned} \quad (1.5)$$

However on this line all the above derivatives are equal and thus we have

$$X_S(\bar{m} + \delta m_u, \bar{m} + \delta m_d, \bar{m} + \delta m_s) = X_S(\bar{m}, \bar{m}, \bar{m}) + O((\delta m_q)^2). \quad (1.6)$$

¹Practically we consider mass degenerate u and d quarks, when $m_u = m_d \equiv m_l$ but the discussion here is more general.

There are many possibilities for singlet quantities. Using hadronic masses we have, for example

$$\begin{aligned} X_N^2 &= \frac{1}{6}(M_p^2 + M_n^2 + M_{\Sigma^+}^2 + M_{\Sigma^-}^2 + M_{\Xi^0}^2 + M_{\Xi^-}^2) = (1.1610 \text{ GeV})^2 \\ X_\pi^2 &= \frac{1}{6}(M_{K^+}^2 + M_{K^0}^2 + M_{\pi^+}^2 + M_{\pi^-}^2 + M_{\bar{K}^0}^2 + M_{K^-}^2) = (0.4116 \text{ GeV})^2 \\ X_\rho^2 &= \frac{1}{6}(M_{K^{*+}}^2 + M_{K^{*0}}^2 + M_{\rho^+}^2 + M_{\rho^-}^2 + M_{\bar{K}^{*0}}^2 + M_{K^{*-}}^2) = (0.8562 \text{ GeV})^2, \end{aligned} \quad (1.7)$$

for octet baryons, pseudoscalar octet mesons and vector octet mesons respectively. Another baryon octet possibility is $X_\Lambda^2 = \frac{1}{2}(M_\Sigma^2 + M_\Lambda^2)$ but other singlet quantities can be constructed using the baryon decuplet. Alternatively gluonic quantities can be used such as the ‘Force’ scale $X_{r_0}^2 = 1/r_0^2$ or the Wilson flow scales, introduced by Lüscher

$$X_{t_0}^2 = \frac{1}{t_0}, \quad X_{w_0}^2 = \frac{1}{w_0^2}, \quad (1.8)$$

(see e.g. [4, 5]). These are all ‘secondary scales’, their physical value has to be determined.

The stationary point of X_S can be checked, using the Gell-Mann–Okubo $SU(3)$ flavour breaking expansion. For example for the pseudoscalar octet mesons we have the expansion

$$\begin{aligned} M_{\pi^+}^2 (= M_{\pi^-}^2) &= M_{0\pi}^2 + \alpha_\pi(\delta m_u + \delta m_d) + O((\delta m_q)^2) \\ M_{K^+}^2 (= M_{K^-}^2) &= M_{0\pi}^2 + \alpha_\pi(\delta m_u + \delta m_s) + O((\delta m_q)^2) \\ M_{K^0}^2 (= M_{\bar{K}^0}^2) &= M_{0\pi}^2 + \alpha_\pi(\delta m_d + \delta m_s) + O((\delta m_q)^2). \end{aligned} \quad (1.9)$$

Constructing X_π^2 gives immediately the result of eq. (1.5). Another check is to use χ -PT (assuming that it is valid in the neighbourhood of the $SU(3)$ flavour plane/line). Simply choose your favourite χ -PT result and expand about a $SU(3)$ flavour symmetric line/point. For example in [6], the chiral expansion for t_0 (for mass degenerate u and d quarks) can be manipulated [1] to give

$$t_0 = T(\bar{\chi}) \left[1 + \frac{1}{(4\pi f_0)^4} \left(\frac{5}{6} k_2 + \frac{1}{4} k_5'' \right) (\chi_s - \chi_l)^2 + \dots \right], \quad (1.10)$$

where T is a (known) function of $\bar{\chi} \equiv 1/3(2\chi_l + \chi_s)$ only. As $(\chi_s - \chi_l) \propto (\delta m_s - \delta m_l)$ then this agrees with our previous assertion: there is no linear term, the first term is quadratic in $SU(3)$ flavour symmetry breaking.

2. Lattice matters

We have generated $2+1$ flavour gauge configurations using an action consisting of tree level Symanzik glue and a mildly stout smeared $O(a)$ non-perturbatively improved clover action, [7], at four- β values, $\beta = 5.40, 5.50, 5.65, 5.80$ on a variety of lattice sizes $24^3 \times 48$, $32^3 \times 64$ and $48^3 \times 96$. All box sizes have $L \gtrsim 2$ fm. All the pion masses used have $M_\pi L > 4$ and range from about 500 to 220 MeV. They are either at points on the $SU(3)$ flavour symmetric line or along lines of constant \bar{m} . This gives 21 data sets at our disposal.

The quark mass m_q and δm_q are given by

$$m_q = \frac{1}{2} \left(\frac{1}{\kappa_q} - \frac{1}{\kappa_{0c}} \right), \quad \delta m_q = m_q - \bar{m} = \frac{1}{2} \left(\frac{1}{\kappa_q} - \frac{1}{\kappa_0} \right), \quad (2.1)$$

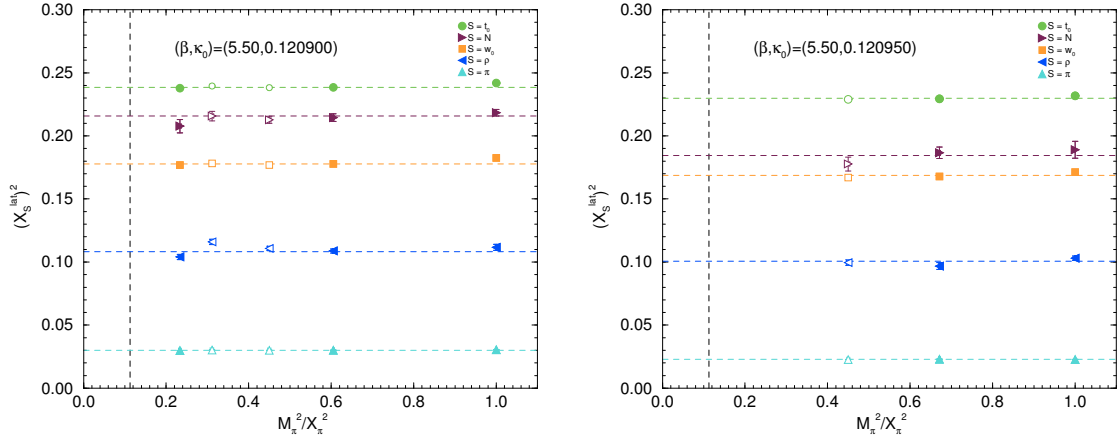


Figure 1: Top to bottom $(X_S^{\text{lat}})^2$ for $S = t_0$ (circles), N (right triangles), w_0 (squares), ρ (left triangles) and π (up triangles) for $(\beta, \kappa_0) = (5.50, 0.120900)$ (left panel) and $(\beta, \kappa_0) = (5.50, 0.120950)$ (right panel) together with constant fits. The opaque points have $M_\pi L < 4$ and are not included in the fits. The vertical line represents the physical point.

where κ_{0c} is chiral limit along symmetric line. (Note that this cancels in δm_q .)

We first investigate the constancy of singlet quantities, as given in eq. (1.6). In Fig. 1 we plot $(X_S^{\text{lat}})^2$ for $S = t_0, N, w_0, \rho$ and π for $(\beta, \kappa_0) = (5.50, 0.120900), (5.50, 0.120950)$. As expected, in agreement with the discussion of section 1, the X_S^2 singlet quantities are constant.

We now take $X_S = \text{const.}$ to determine the scale

$$a_S^2(\kappa_0) = \frac{X_S^{\text{lat}2}(\kappa_0)}{X_S^{\text{exp}2}}. \quad (2.2)$$

This is a function of m_0 or here κ_0 . So if we vary κ_0 (for example as in Fig. 1) – when pairs $a_S, a_{S'}$ cross this gives a common lattice spacing a . We apply this in particular here to²

$$(S, S') = (\pi, N), (\pi, \rho). \quad (2.3)$$

For $S = t_0, w_0$ we can arrange $X_{t_0}^{\text{exp}}, X_{w_0}^{\text{exp}}$ (from eq. (2.2)) so that these singlet quantities also cross at the same point. In Fig. 2 we show these crossings for $\beta = 5.50$. From the results for the four beta values we can now make the last, continuum extrapolation. This is shown in Fig. 3. A (weighted) average of these results gives our final estimates for $\sqrt{t_0^{\text{exp}}}, w_0^{\text{exp}}$ as found in [1].

Alternatively we can write

$$\frac{2M_K^2 - M_\pi^2}{X_S^2} = C - 2\frac{M_\pi^2}{X_S^2}, \quad (2.4)$$

($C = X_\pi^2/X_S^2$) for $S = N, \rho, t_0, w_0$. In Fig. 4 we plot this function for $(\beta, \kappa_0) = (5.50, 0.120900)$,

²For the β and pion mass values considered here, the ρ and K^* are stable particles.

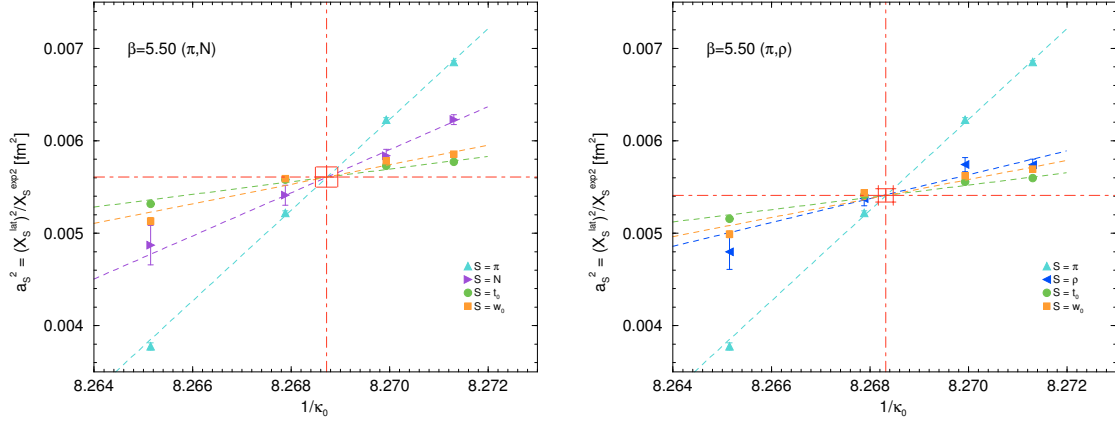


Figure 2: a_S^2 against $1/\kappa_0$ for $S = \pi, N$ and t_0, w_0 together with quadratic fits for $\beta = 5.50$. Left panel: based on (π, N) crossing; Right panel: based on (π, ρ) crossing.

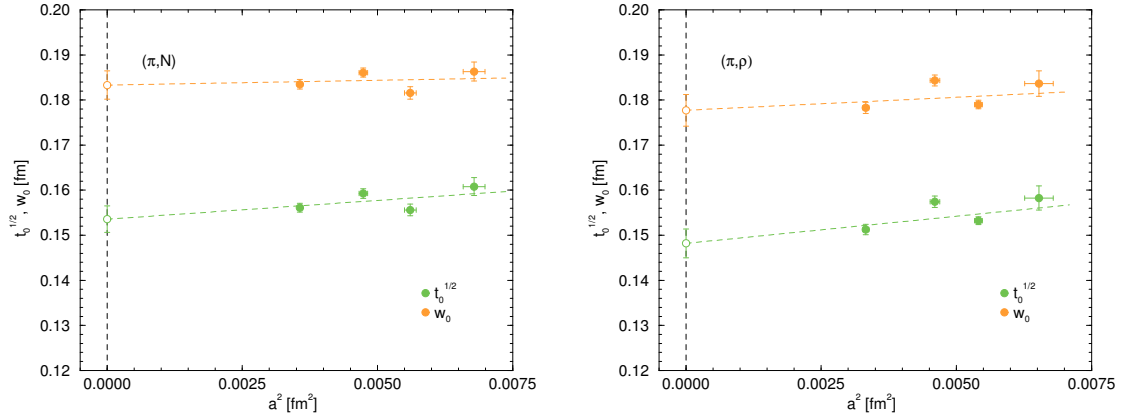


Figure 3: $\sqrt{t_0}$ and w_0 (in fm) against a^2 (in fm^2) from the (π, N) crossing (left panel) and (π, ρ) (right panel) crossing together with a linear fit.

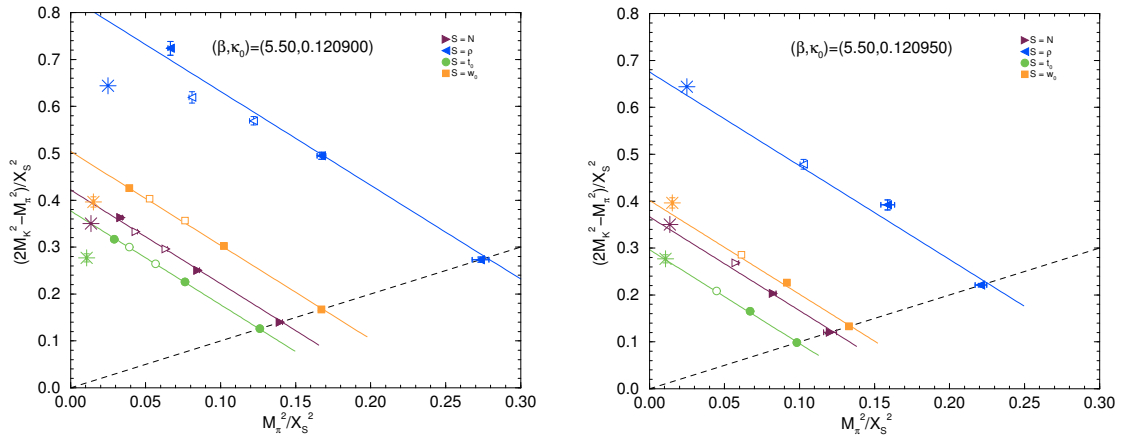


Figure 4: $(2M_K^2 - M_\pi^2)/X_S^2$ against M_π^2/X_S^2 , together with the fit from eq. (2.4) for $(\beta, \kappa_0) = (5.50, 0.120900)$ (left panel) and $(5.50, 0.120950)$ (right panel). The stars correspond to the phenomenological values.

(5.50, 0.120950). This represents the path in the quark mass plane. Also shown are the experimental values (now including those of $S = t_0, w_0$). We see that these κ_0 values straddle the optimum κ_0^* – it is clear that κ_0^* lies closer to 0.120950 than 0.120900.

Finally we comment on our results. In the left panel of Fig. 5 we plot a^2 against g_0^2 . The curve

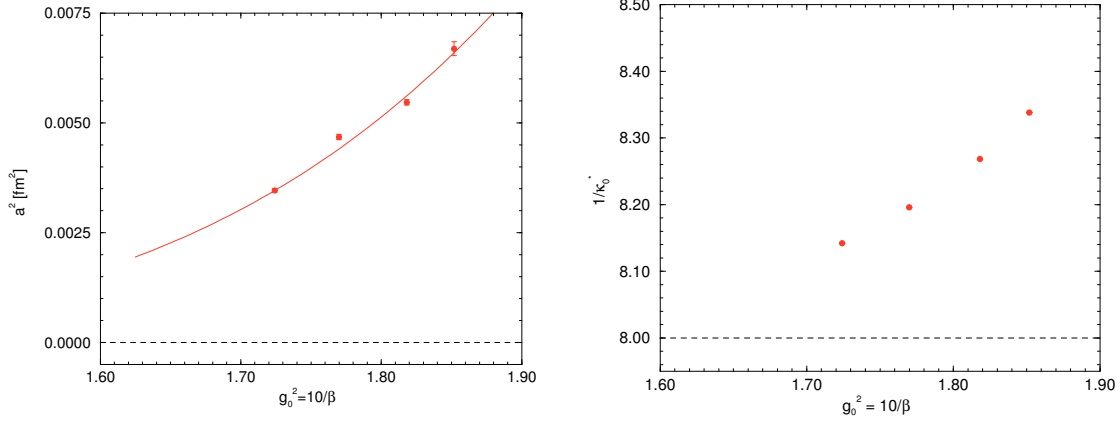


Figure 5: Left panel: a^2 versus g_0^2 . The curve is from the running coupling constant, using the 2-loop beta function normalised to the $\beta = 5.80$ result. Right panel: $1/\kappa_0^*$ versus g_0^2 . The horizontal dashed lines represents the value in the continuum limit.

is the running coupling constant, $g_0^2 = 10/\beta$ using the 2-loop QCD beta function, normalised to $\beta = 5.80$, namely

$$\frac{a^2(\beta)}{a^2(\beta_0)} = \left(\frac{\beta_0}{\beta}\right)^{-\frac{b_1}{b_0}} \exp\left(-\frac{1}{10b_0}(\beta - \beta_0)\right), \quad \beta_0 = 5.80, \quad (2.5)$$

(b_0, b_1 are the first two coefficients of the beta function). There seems to be reasonable agreement between the data points and the curve. The right hand panel of Fig. 5 indicates how the initial point, κ_0^* , on the $SU(3)$ flavour symmetric line changes with g_0^2 .

3. Conclusions

Our programme is to tune strange and light quark masses to their physical values simultaneously by keeping $\bar{m} = 1/3(2m_l + m_s) = \text{const.}$. As the light quark mass is decreased then $M_\pi \searrow$ and $M_K \nearrow$. Singlet quantities, here denoted by $X_S(\kappa_0)$ remain constant starting from a point on the $SU(3)$ flavour symmetric line — the Gell-Mann–Okubo result. We can use this result and X_S^{exp} to determine the $a_S(\kappa_0)$ scale. Varying κ_0 – determines when pairs of singlet quantities such as (X_π, X_N) and (X_π, X_ρ) cross giving a common lattice spacing a . By arranging so that X_{t_0}, X_{w_0} also cross here, we are able to give a determination of the ‘secondary’ scales $\sqrt{t_0^{\text{exp}}}$ and w_0^{exp} [fm]. Finally in Fig. 6 a comparison with other results is given.

Acknowledgements

The numerical configuration generation (using the BQCD lattice QCD program [10]) and data analysis (using the Chroma software library [11]) was carried out on the IBM BlueGene/Q using

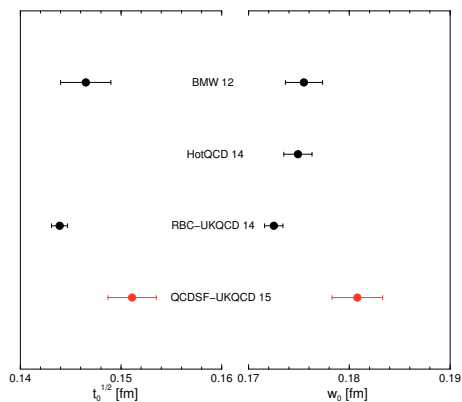


Figure 6: $\sqrt{t_0^{\text{exp}}}$, left plot and w_0^{exp} , right plot in fm for BMW 12 [5], HotQCD 14 [8], RBC-UKQCD 14 [9], together with the present results.

DIRAC 2 resources (EPCC, Edinburgh, UK), the BlueGene/P and Q at NIC (Jülich, Germany), the Lomonosov at MSU (Moscow, Russia) and the SGI ICE 8200 and Cray XC30 at HLRN (The North-German Supercomputer Alliance) and on the NCI National Facility in Canberra, Australia (supported by the Australian Commonwealth Government). HP was supported by DFG Grant No. SCHI 422/10-1. PELR was supported in part by the STFC under contract ST/G00062X/1 and JMZ was supported by the Australian Research Council Grant No. FT100100005 and DP140103067. We thank all funding agencies.

References

- [1] V. G. Bornyakov *et al.* [QCDSF–UKQCD Collaboration], arXiv:1508.05916[hep-lat].
- [2] W. Bietenholz *et al.* [QCDSF–UKQCD Collaboration], Phys. Lett. **B690** (2010) 436, [arXiv:1003.1114[hep-lat]].
- [3] W. Bietenholz *et al.* [QCDSF–UKQCD Collaboration], Phys. Rev. **D84** (2011) 054509, [arXiv:1102.5300[hep-lat]].
- [4] M. Lüscher, JHEP **1008** (2010) 071, [arXiv:1006.4518[hep-lat]].
- [5] S. Borsanyi *et al.* [BMW Collaboration], JHEP **010** (2012) 1209, [arXiv:1203.4469[hep-lat]].
- [6] O. Bär *et al.* Phys. Rev. **D89** (2014) 034505, [arXiv:1312.4999[hep-lat]].
- [7] N. Cundy *et al.* [QCDSF–UKQCD Collaboration], Phys. Rev. **D79** (2009) 094507, [arXiv:0901.3302[hep-lat]].
- [8] A. Bazavov *et al.* [HotQCD Collaboration], Phys. Rev. **D90** (2014) 094503, [arXiv:1407.6387[hep-lat]].
- [9] T. Blum *et al.* [RBC–UKQCD Collaborations], arXiv:1411.7017[hep-lat].
- [10] Y. Nakamura *et al.*, Proc. Sci. **Lattice 2010** (2010) 040, arXiv:1011.0199[hep-lat].
- [11] R. G. Edwards *et al.*, Nucl. Phys. Proc. Suppl. **140** (2005) 832, arXiv:hep-lat/0409003.