

ARTICLE

DOI: 10.1038/s41467-017-02524-x

OPEN

# Antiferromagnetic correlations in the metallic strongly correlated transition metal oxide $\text{LaNiO}_3$

H. Guo<sup>1</sup>, Z.W. Li<sup>1</sup>, L. Zhao<sup>1</sup>, Z. Hu<sup>1</sup>, C.F. Chang<sup>1</sup>, C.-Y. Kuo<sup>1</sup>, W. Schmidt<sup>2</sup>, A. Piovano<sup>3</sup>, T.W. Pi<sup>4</sup>, O. Sobolev<sup>5</sup>, D.I. Khomskii<sup>6</sup>, L.H. Tjeng<sup>1</sup> & A.C. Komarek<sup>1</sup>

The material class of rare earth nickelates with high  $\text{Ni}^{3+}$  oxidation state is generating continued interest due to the occurrence of a metal-insulator transition with charge order and the appearance of non-collinear magnetic phases within this insulating regime. The recent theoretical prediction for superconductivity in  $\text{LaNiO}_3$  thin films has also triggered intensive research efforts.  $\text{LaNiO}_3$  seems to be the only rare earth nickelate that stays metallic and paramagnetic down to lowest temperatures. So far, centimeter-sized impurity-free single crystal growth has not been reported for the rare earth nickelates material class since elevated oxygen pressures are required for their synthesis. Here, we report on the successful growth of centimeter-sized  $\text{LaNiO}_3$  single crystals by the floating zone technique at oxygen pressures of up to 150 bar. Our crystals are essentially free from  $\text{Ni}^{2+}$  impurities and exhibit metallic properties together with an unexpected but clear antiferromagnetic transition.

<sup>1</sup>Max-Planck-Institute for Chemical Physics of Solids, Nöthnitzer Street 40, 01187 Dresden, Germany. <sup>2</sup>Forschungszentrum Jülich GmbH, Jülich Centre for Neutron Science at ILL, CS 20156, 71 avenue de Martyrs, 38042 Grenoble, France. <sup>3</sup>Institut Laue-Langevin (ILL), 71 avenue des Martyrs, 38042 Grenoble, France. <sup>4</sup>National Synchrotron Radiation Research Center (NSRRC), 101 Hsin-Ann Road, 30077 Hsinchu, Taiwan. <sup>5</sup>Georg-August-Universität Göttingen, Institut für Physikalische Chemie, Tammannstrasse 6, 37077 Göttingen, Germany. <sup>6</sup>Physics Institute II, University of Cologne, Zùlpicher Street 77, 50937 Cologne, Germany. H. Guo and Z.W. Li contributed equally to this work. Correspondence and requests for materials should be addressed to A.C.K. (email: [Alexander.Komarek@cpfs.mpg.de](mailto:Alexander.Komarek@cpfs.mpg.de))

The rare earth nickelates  $R\text{NiO}_3$  ( $R$ =rare earth, Y) with the high  $\text{Ni}^{3+}$  oxidation state have continued to attract enormous interest due to the famous bandwidth controlled metal-insulator (MI) transition and associated unusual charge and spin-order phenomena occurring in this system<sup>1,2</sup> with even the possibility for multiferroicity<sup>3</sup>. The description of the underlying physics of these phenomena turned out to be a true intellectual challenge and it is not a surprise that important theoretical concepts have been (and still need to be) developed along the way<sup>4–12</sup>. More recently, the prediction of high- $T_c$  superconductivity in  $\text{LaNiO}_3$ -based heterostructures<sup>13</sup> has triggered a flurry of new activities on  $\text{LaNiO}_3$ - $\text{LaMO}_3$  superlattices ( $M$ =other metal ion)<sup>14,15</sup>.

Apart from  $\text{LaNiO}_3$ , all  $R\text{NiO}_3$  ( $R$ =Pr–Lu, Y) compounds exhibit a MI transition at lower temperatures<sup>1,2</sup> with an insulating antiferromagnetic ground state. With decreasing  $R$ -ionic radius, the octahedral tilts become larger. Thus, the Ni–O–Ni bond angles become smaller, which alters the electronic bandwidth and the magnetic exchange interactions. Whereas an enhancement of the insulating properties and an increase of the MI transition temperature  $T_{\text{MI}}$  can be observed for decreasing  $R$ -ionic radius, the antiferromagnetic transition temperature and the strength of the exchange interactions decrease with decreasing  $R$ -ionic radius and Ni–O–Ni bond angles. With La having the largest ionic radius of the series, one might expect the strongest antiferromagnetic properties for  $\text{LaNiO}_3$ . However, so far it was reported that  $\text{LaNiO}_3$  does not show any magnetic order, and thus, violates the trend. It remains a paramagnetic metal down to lowest temperatures with an enhanced effective mass<sup>1,2,16,17</sup>.

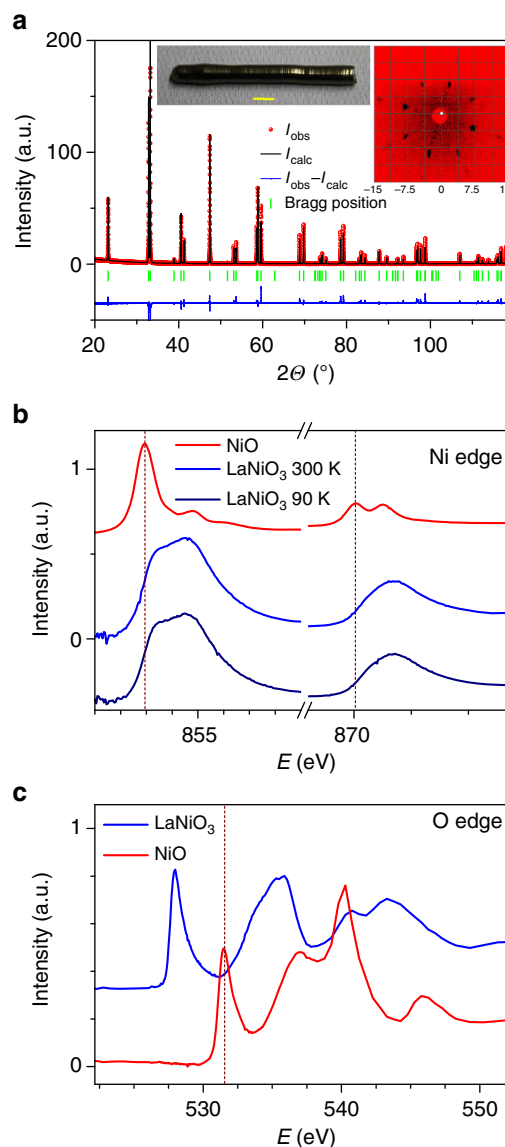
Here, we report on the successful growth of large  $\text{LaNiO}_3$  single crystals by the floating zone technique. Electrical resistivity and Hall effect measurements on our single crystals show that  $\text{LaNiO}_3$  is intrinsically not a bad metal as recently discussed for the  $R\text{NiO}_3$  ( $R$ =Pr–Lu, Y) compounds in their paramagnetic phase<sup>10</sup>. On the contrary, we found that  $\text{LaNiO}_3$  has a high conductivity. Moreover, we were able to observe bulk antiferromagnetism in magnetization, specific heat, and neutron scattering experiments. Thus,  $\text{LaNiO}_3$  appears to be a highly metallic and antiferromagnetic transition metal oxide—a rather rare combination in oxides. Special about  $\text{LaNiO}_3$  is that it is also very close to an insulating state, making  $\text{LaNiO}_3$  an intriguing quantum material, probably close to a quantum critical point, where strong local electronic correlations at the Ni sites are likely to interfere in an intricate manner with Fermi surface effects.

## Results

**Crystal growth and characterization.** Using high oxygen pressures of 130–150 bar, we were able to grow  $\text{LaNiO}_3$  single crystals with the floating zone technique, see inset of Fig. 1a. Further details can be found in the Methods section and in Supplementary Notes 1 and 2, in Supplementary Figures 1 and 2, and in Supplementary Movie 1. Powdered crystals exhibit no impurity phases within the accuracy of powder X-ray diffraction measurements, see Fig. 1a. The single crystallinity of our  $\text{LaNiO}_3$  crystal is confirmed by a series of Laue diffraction patterns from different directions and across the length of this crystal as well as by single crystal X-ray diffraction, see inset of Fig. 1a and Supplementary Figure 1, as well as Supplementary Note 3, Supplementary Figure 3, and Supplementary Table 1. Due to the occurrence of a cubic to rhombohedral phase transition ( $Pm\bar{3}m \rightarrow R\bar{3}c$ ) somewhat below 1100 K<sup>2</sup>, our floating zone grown single crystals are twinned.

Figure 1b shows the Ni- $L_{2,3}$  X-ray absorption spectroscopy (XAS) data of the  $\text{LaNiO}_3$  single crystal taken at 300 and 90 K together with that of a NiO single crystal serving as a  $\text{Ni}^{2+}$

reference compound. We have removed the La- $M_4$  white lines located at 850.6 eV from the  $\text{LaNiO}_3$  spectra using the La- $M_4$  spectrum of  $\text{LaCoO}_3$ . It is well known that XAS spectra at the  $L_{2,3}$  edge of transition metal oxides are highly sensitive to the valence state, in particular, an increase of the valence state of the transition metal ion by one causes a shift of the XAS  $L_{2,3}$  spectra by one or more eV toward higher energies<sup>18</sup>. The more than one eV higher energy shift between the spectra of NiO and  $\text{LaNiO}_3$  indicates the formal  $\text{Ni}^{3+}$  valence state in  $\text{LaNiO}_3$ <sup>18,19</sup>, with the note that the spectrum cannot be interpreted in terms of a  $\text{Ni } 3d^7$  configuration, but, rather by a coherent mixture of  $3d^8$  and  $3d^{\underline{8}\underline{L}}$  configurations<sup>4,5,9,12</sup>, where each  $\underline{L}$  denotes a hole in the oxygen ligand. Here we can exclude  $\text{Ni}^{2+}$  impurities in our  $\text{LaNiO}_3$  single



**Fig. 1** Characterization of  $\text{LaNiO}_3$ . **a** Powder X-ray diffraction pattern of a powdered  $\text{LaNiO}_3$  single crystal. The inset shows a photo of the  $\text{LaNiO}_3$  single crystal and its Laue diffraction pattern. A length scale of 1 cm is indicated by the yellow bar within the photo of the crystal and the distance between two lines in the Laue pattern amounts to 3.75 cm. **b** The Ni- $L_{2,3}$  XAS spectra of the  $\text{LaNiO}_3$  single crystal measured at 90 and 300 K together with the spectrum of a NiO reference compound. **c** The O-K XAS spectra of our  $\text{LaNiO}_3$  single crystal measured at 300 K (blue line) together with the spectra of the NiO reference compound (red line)

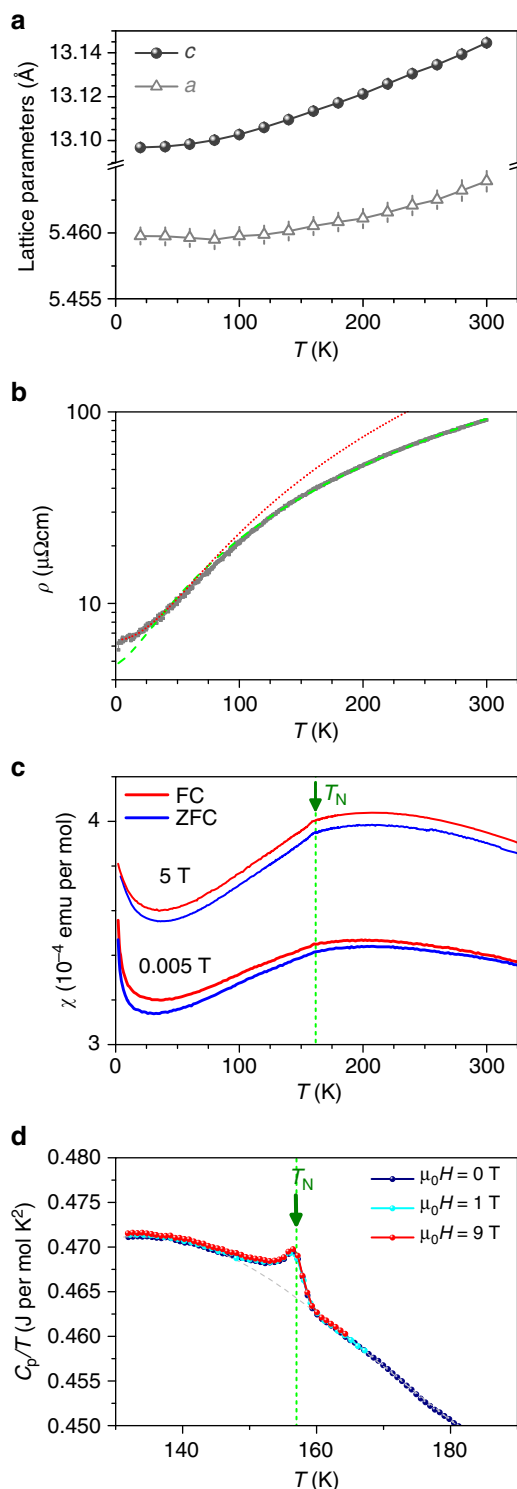
crystal—otherwise the sharp main peak of the  $\text{Ni}^{2+}$  impurity spectrum would have been visible as a sharp shoulder at the leading edge<sup>18</sup>. At the O-K edge, the pre-edge peak is shifted by about 1 eV to higher energies when going from  $\text{Ni}^{3+}$  to  $\text{Ni}^{2+}$ <sup>18,19</sup>. Figure 1c shows the O-K XAS spectra of our  $\text{LaNiO}_3$  single crystal (blue) and the  $\text{NiO}$  reference compound (red). The O-K XAS spectra demonstrate even more clearly that there is no spectral feature from  $\text{Ni}^{2+}$  impurities and thus that our as-grown  $\text{LaNiO}_3$  single crystals are highly stoichiometric. This is further confirmed by thermogravimetric and inductively coupled plasma optical

emission spectroscopy (ICP-OES) measurements, see Methods and Supplementary Note 2.

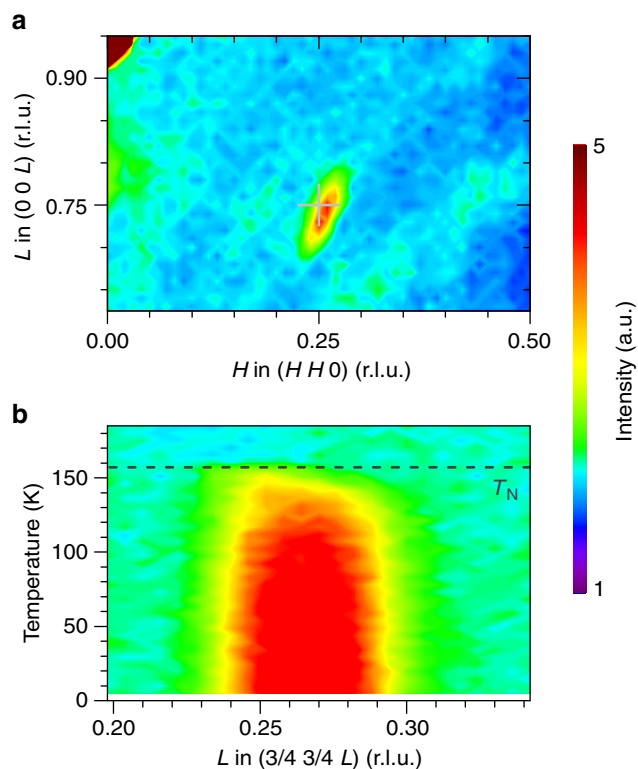
**Temperature dependence of physical properties.** In Fig. 2a, we show the temperature dependence of the lattice parameters of our  $\text{LaNiO}_3$  single crystal that has been powdered and measured by means of powder X-ray diffraction. There is no clear indication for the presence of a structural anomaly which otherwise occurs readily in the other nickelates  $\text{RNiO}_3$  with smaller rare earths ( $R = \text{Pr-Lu, Y}$ ) when cooling through the MI transition<sup>1,2</sup>.

The electrical resistivity  $\rho$  of our  $\text{LaNiO}_3$  single crystal that is shown in Fig. 2b is in the  $\mu\Omega \text{ cm}$  range, reaching  $\sim 6 \mu\Omega \text{ cm}$  at low temperatures, i.e., distinctly lower than that of  $\text{LaNiO}_3$  powder samples<sup>1,2,16,17</sup> or  $\text{LaNiO}_3$  thin films<sup>14</sup> or of crystals grown under 50 bar oxygen pressure<sup>20</sup>. Thus, our  $\text{LaNiO}_3$  single crystals are more metallic, see also Supplementary Note 4 and Supplementary Figure 4(a). Also the measurement of the Hall effect on the single crystal  $\text{LaNiO}_3$  (for temperatures between 2 and 300 K) yields high carrier densities of about  $4.8 \times 10^{28} \text{ m}^{-3}$ , corresponding to  $\sim 2.7$  (close to 3) electrons per formula unit, consistent with the 3+ valence of the Ni ions in  $\text{LaNiO}_3$ . Up to  $\sim 35 \text{ K}$ , the low-temperature behavior of the resistivity is Fermi-liquid-like,  $\rho(T) = \rho_0 + AT^n$ , with  $\rho_0 \sim 6.45 \mu\Omega \text{ cm}$ ,  $A = 1.62 \cdot 10^{-3} \mu\Omega \text{ cm K}^{-2}$ , and  $n \sim 2.0$ , see Fig. 2b. The value of  $A$  is of the same order as reported in literature<sup>21</sup>. A fit over the entire temperature range gives an exponent  $n \sim 1.50(1)$  with the residual resistivity  $\rho_0 \sim 4.83 \mu\Omega \text{ cm}$ , similar to the value of the fit at low temperatures. Most probably this behavior is a signature that  $\text{LaNiO}_3$  is close to a quantum critical point (cf. the results for  $\text{PrNiO}_3$  under pressure<sup>21</sup>).

Figure 2c displays the temperature dependence of the magnetic susceptibility  $\chi$  of our  $\text{LaNiO}_3$  single crystal, see also Supplementary Note 4 and Supplementary Figure 4(b). First of all, we notice that it shows a significantly smaller low-temperature upturn than reported previously for powder and ceramic samples<sup>16,17</sup>, which confirms unprecedented high quality of our single crystals. Surprisingly,  $\chi$  exhibits an anomalous kink at  $\sim 157 \text{ K}$ , which we take as an indication for a hitherto unknown antiferromagnetic transition in  $\text{LaNiO}_3$ . That this anomaly is not simply caused by a signal from a tiny fraction of a magnetic impurity phase (which is so small that it is not visible in our powder X-ray diffraction measurements) can be excluded by our specific heat ( $C_p$ ) measurements. As can be seen in Fig. 2d, there is a small but clearly visible anomalous peak at  $\sim 157 \text{ K}$  in  $C_p/T$ . Also the resistivity data excludes that the anomaly in the susceptibility is caused by the presence of an oxygen-deficient  $\text{LaNiO}_3$  minority phase that becomes antiferromagnetic and insulating at low temperatures<sup>22,23</sup>: we do not observe an upturn or a slowing down of the decrease in the resistivity on cooling. Moreover, our samples have resistivities in the  $\mu\Omega \text{ cm}$  range and have conductivities higher than reported so far. All these support the notion that the transition is an intrinsic bulk property and not



**Fig. 2** Temperature dependence of physical properties of  $\text{LaNiO}_3$ . **a** Lattice parameters of our  $\text{LaNiO}_3$  as a function of temperature (space group  $R\bar{3}c$ ). The error bars are the standard deviations obtained from Rietveld refinement using the Fullprof software. **b** Electrical resistivity of  $\text{LaNiO}_3$  (gray). The red line indicates a fit to  $\rho_0 + AT^n$  (for  $T < 35 \text{ K}$ ) and the dashed green line is a fit over the entire temperature range. **c** Magnetic susceptibility of  $\text{LaNiO}_3$  under field cooled (FC) and zero-field cooled (ZFC) conditions for applied fields of  $\mu_0 H = 0.005 \text{ T}$  and  $5 \text{ T}$ . The high- and low-field measurements both together indicate an antiferromagnetic transition at  $T_N$  that is an intrinsic property of  $\text{LaNiO}_3$ . **d** Specific heat measurements of  $\text{LaNiO}_3$  in fields of  $0 \text{ T}$ ,  $1 \text{ T}$ , and  $9 \text{ T}$ . These measurements further corroborate that the antiferromagnetic transition is a bulk property of  $\text{LaNiO}_3$ .

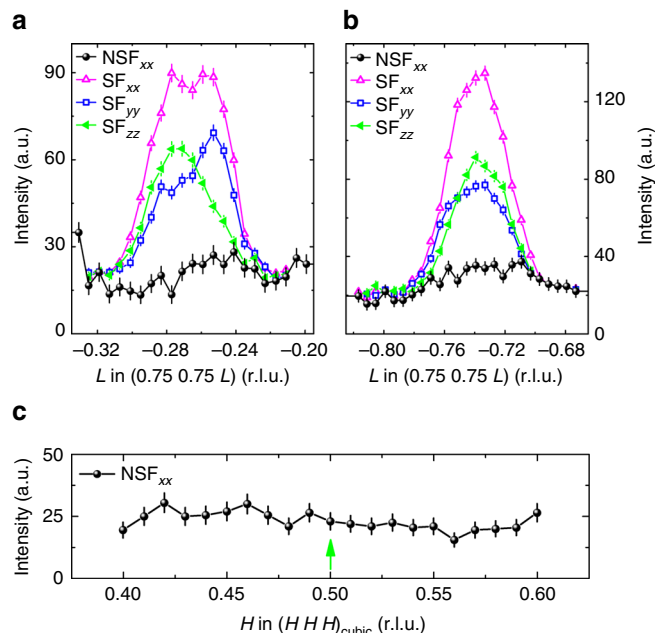


**Fig. 3** Single crystal neutron-diffraction measurements. **a** Neutron scattering intensities within the  $(H H L)$  plane of reciprocal space (in pseudocubic notation) of the  $\text{LaNiO}_3$  single crystal. These intensities were obtained at 1.6 K. **b** The temperature dependence of a quarter-integer peak as measured in  $(3/4 \ 3/4 \ L)$  scans for different temperatures. The dashed line indicates the transition temperature ( $T_N$ ) observed in the specific heat

due to an impurity phase. The low-temperature behavior of the specific heat supports a Fermi-liquid type of the ground state of  $\text{LaNiO}_3$ ,  $C_p(T) = \gamma T + \beta T^3$ , with  $\gamma \sim 17 \text{ mJ mol}^{-1} \text{ K}^{-2}$ , consistent with the value  $18 \text{ mJ mol}^{-1} \text{ K}^{-2}$  reported before<sup>17</sup>, and with  $\beta = 1.87(5) \times 10^{-4} \text{ J mol}^{-1} \text{ K}^{-4}$ , thus yielding a Debye temperature  $\theta_D = (12\pi^4 N R / 5\beta)^{1/3} = 373 \text{ K}$ , where  $N$  is the number of atoms in the chemical formula and  $R$  is the ideal gas constant.

**Neutron scattering experiments.** The availability of sizeable  $\text{LaNiO}_3$  single crystals also enabled us to study this intriguing system by means of neutron diffraction and inelastic neutron scattering. These experiments were performed at the Thales, IN12 and IN8 spectrometers at the ILL in Grenoble, France. Within elastic scans, we were able to observe quarter-integer peaks at low temperatures, see Fig. 3a. These quarter-integer peaks (in pseudocubic notation) resemble those found in powder neutron-diffraction experiments within the insulating regime of  $\text{RNiO}_3$ <sup>24</sup>. The propagation vector observed for the insulating antiferromagnetic regime of  $\text{RNiO}_3$  amounts to  $(1/4 \ 1/4 \ 1/4)$  in pseudocubic notation (or  $(1/2 \ 0 \ 1/2)$  in orthorhombic notation)<sup>2</sup>. Moreover, the study of the temperature dependence of these quarter-integer peaks indicates an onset temperature that coincides with the magnetic ordering temperature  $T_N$  that we observed in susceptibility and specific heat measurements of  $\text{LaNiO}_3$ , see Fig. 3b.

The magnetic origin of these quarter-integer peaks in  $\text{LaNiO}_3$  could be also confirmed by polarized neutrons at the IN12 spectrometer. In Fig. 4a & b  $L$ -scans across two quarter-integer peaks are shown for the three spin-flip (SF) channels and for a non-spin-flip channel. Only in the SF channels, neutron



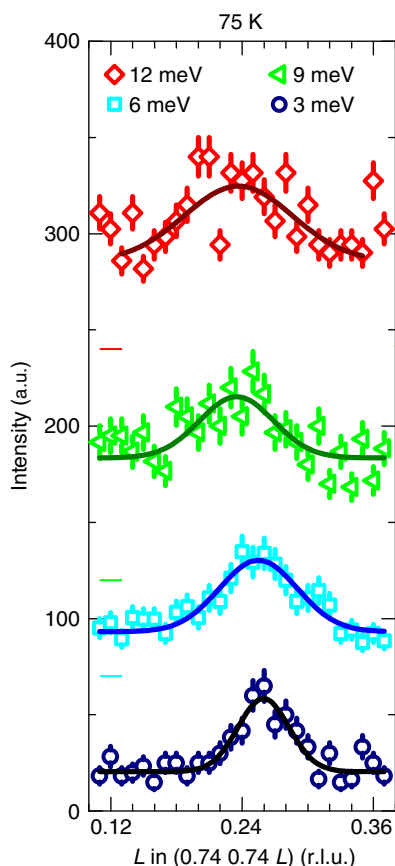
**Fig. 4** Polarized neutron-diffraction measurements. **a, b** Polarized neutron scattering experiments for the  $\text{LaNiO}_3$  single crystal at 2 K. Spin-flip (SF) and non-spin-flip (NSF) channels were measured for neutron spins parallel to the scattering vector ( $x$ -direction), perpendicular to the scattering plane ( $z$ -direction) or perpendicular to  $x$  and  $z$  ( $y$ -direction). **c** A structural scan at 2 K showing no indications for any charge ordering peak at the  $(0.5 \ 0.5 \ 0.5)$  position in pseudocubic notation (indicated by the green arrow). The intensity error bars are statistical error bars calculated by the square root of intensity

scattering intensities  $\sigma_{x_i x_i}$  can be detected. This unambiguously shows the magnetic nature of these quarter-integer reflections in  $\text{LaNiO}_3$ . Although we cannot detect a symmetry lowering from our high-resolution powder X-ray diffraction measurements—see Fig. 2a—the pseudocubic propagation vector for rhombohedral  $\text{LaNiO}_3$  is the same as the pseudocubic propagation vector for the orthorhombic insulating nickelates  $\text{RNiO}_3$  ( $R = \text{Pr-Lu, Y}$ ). Based on the magnetic symmetry analysis for the high symmetry cubic structure with space group  $Pm\bar{3}m$  and for the propagation vector  $(1/4 \ 1/4 \ 1/4)$ , a helical magnetic structure with moments spiraling perpendicular to the propagation vector is consistent with our single crystal neutron data, see Supplementary Note 2, Supplementary Tables 2 and 3, and Supplementary Figure 6. The size of the ordered moment amounts to  $\sim 0.3\mu_B$ , which indicates that magnetism in  $\text{LaNiO}_3$  is a bulk property and does not originate from a tiny (insulating) impurity phase. This small magnetic moment might explain why paramagnetic properties have been reported for  $\text{LaNiO}_3$  in the past<sup>1,2,16,17</sup>.

We also have been able to study and observe the magnetic excitations in  $\text{LaNiO}_3$  by means of inelastic neutron scattering, thereby also providing further support that the antiferromagnetism is a bulk property. As can be seen in Fig. 5, magnetic excitations are clearly visible up to at least 12 meV. With increasing energy transfer, the magnetic peaks become somewhat broader and damped, which is indicative for the appearance of fluctuations in this itinerant antiferromagnetic system. These fluctuations could be also responsible for a reduced ordered moment of  $\sim 0.3\mu_B$  in  $\text{LaNiO}_3$ .

## Discussion

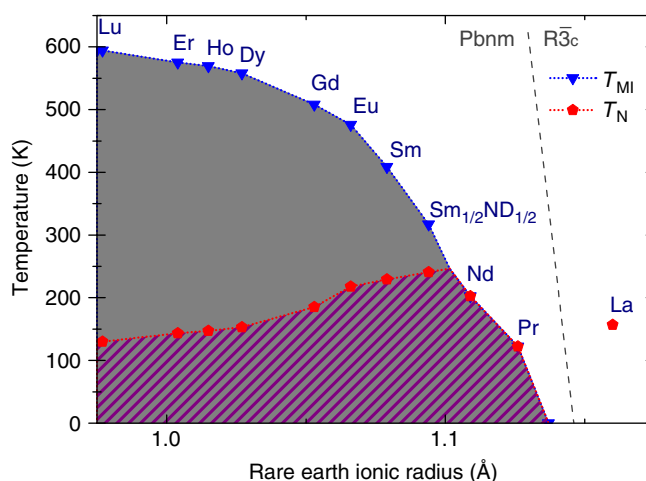
$\text{LaNiO}_3$  appears to be a rare case of an antiferromagnetic and metallic transition metal oxide with a fully three-dimensional



**Fig. 5** Magnetic excitations in  $\text{LaNiO}_3$ . Inelastic neutron scattering intensities measured at 75 K for 3, 6, 9, and 12 meV energy transfer. The horizontal bars indicate the shift of the data in vertical direction. Up to 12 meV, the magnetic peaks become broader and damped, which is indicative for the appearance of fluctuations in this itinerant antiferromagnetic system. The intensity error bars are statistical error bars calculated by the square root of intensity

(3D) crystal and electronic structure. Other systems like  $(\text{La}, \text{Sr})_3\text{Mn}_2\text{O}_7$  and  $\text{Ca}_3\text{Ru}_2\text{O}_7$  have a lower electronic and structural dimensionality where the antiferromagnetic order is resulting from a stacking of ferromagnetic layers<sup>25,26</sup>. So far, the chromate system  $\text{ACrO}_3$  ( $A = \text{Ca}, \text{Sr}$ ) and very recently,  $\text{RuO}_2$  are known to be the only other intrinsically antiferromagnetic and metallic transition metal oxides with such a fully 3D crystal and electronic structure<sup>27,28</sup>. The systems  $\text{CaCrO}_3$  and  $\text{LaNiO}_3$  have in common that the oxidation state of the transition metal ion is very high. Thus, the oxygen  $2p$  to transition metal  $3d$  charge-transfer energy here is apparently negative<sup>5,12,29</sup> resulting in extreme  $2p$ - $3d$  covalency where the presence of holes in the oxygen band can effectively prevent the opening of the conductivity gap and at the same time mediate strongly the magnetic exchange interactions. However, in contrast to  $\text{CaCrO}_3$ <sup>27</sup>, the rare earth nickelate  $\text{LaNiO}_3$  in single crystalline form is much more metallic and shows conductivities in the  $\mu\Omega \text{ cm}$  range. This is probably also true for  $\text{RuO}_2$ <sup>28</sup>. Unique for  $\text{LaNiO}_3$  is that it is close to the insulating phase of  $\text{RNiO}_3$  ( $R = \text{Pr-Lu}, \text{Y}$ ) indicating the importance of strong correlation effects, making it rather exceptional among all transition metal oxides.

Although charge order, which plays a very important role in nickelates  $\text{RNiO}_3$  with small rare earth ions  $R$ , cannot be observed in our present X-ray diffraction (XRD) and neutron-diffraction measurements—see Fig. 4c—the presumably associated symmetry lowering of the structure that goes along with the antiferromagnetic order may produce only very weak new peaks in diffraction experiments. Some hints for this may in fact be found in a recent pair distribution function study using neutron diffraction<sup>30</sup>. Nevertheless, with the charge order effects being so weak, one could infer that  $\text{LaNiO}_3$  is perhaps better described using Fermi surface arguments<sup>7,28</sup>, while the other insulating nickelates  $\text{RNiO}_3$  with the smaller  $R$ -ionic sizes and with higher charge ordering temperatures can be more intuitively understood in terms of local charge or bond disproportionations<sup>5,6,8,9,11,12</sup>. Note, that recent ab initio calculations<sup>31</sup> reproduced metallic and antiferromagnetic state of  $\text{LaNiO}_3$ , and also predicted charge (or rather bond) disproportionation, whose magnitude is however



**Fig. 6** Tentative  $\text{RNiO}_3$  phase diagram. The blue triangles indicate the metal-insulator transition temperatures  $T_{\text{MI}}$  associated with the onset of charge ordering and the appearance of monoclinic distortions. The red symbols indicate the Néel temperatures  $T_{\text{N}}$  associated with the onset of antiferromagnetic ordering. Data points for other  $R$  than La were taken from literature<sup>24,32</sup>. The gray dashed curve subdivides orthorhombic (monoclinic) and rhombohedral regimes



below our experimental detection limit. One needs special dedicated experiments to probe for the eventual symmetry lowering with the appearance of inequivalent Ni ions. However, our main conclusion—the existence of antiferromagnetic ordering in highly metallic single crystals of  $\text{LaNiO}_3$ , is confirmed by these calculations. According to our findings, we now also present a tentative  $\text{RNiO}_3$  phase diagram in Fig. 6.

## Methods

**Chemical synthesis.** The  $\text{LaNiO}_3$  single crystal was grown under high oxygen pressures of 130–150 bar with a growth speed of 6–7 mm per h in a mirror furnace from Scidre that was equipped with a 5000 W Xe lamp and with counter-rotation of feeding and seeding rods—see Supplementary Movie 1. The seeding and feeding rods were synthesized by pressing appropriate mixtures of  $\text{La}_2\text{O}_3$  and NiO into rods with roughly 10 cm length and 6 mm diameter. These rods were sintered at 800–1000 °C for several days. The temperature of the melting zone was measured in situ by means of a pyrometer and is close to 1800 °C at ~140 bar  $\text{p}_{\text{O}_2}$ .

**X-ray diffraction.** Laue diffraction measurements were performed on a multiwire real-time back-reflection Laue camera from Multiwire Laboratories; see also Supplementary Note 1 and Supplementary Figure 1.

Powder X-ray diffraction measurements were performed on a Bruker D8 Discover A25 diffractometer, which is equipped with a Johansson monochromator for Cu  $K_{\alpha 1}$  radiation. A closed cycle helium cryostat (Phenix of Oxford Cryosystems) was used for temperature-dependent measurements.

Single crystal X-ray diffraction measurements have been performed on a twinned single crystal of  $\text{LaNiO}_3$  using a Bruker D8 VENTURE single crystal X-ray diffractometer equipped with a bent graphite monochromator for Mo  $K_{\alpha}$  radiation (about  $3\times$  intensity enhancement) and a Photon CMOS large area detector. A crystal with roughly 20  $\mu\text{m}$  diameter has been measured and a multi-scan absorption correction has been applied to the data (minimum and maximum transmission: 0.6184 and 0.7519, respectively). A total of 12,328 (observed) reflections ( $H$ :  $-12\rightarrow 13$ ,  $K$ :  $-13\rightarrow 9$  and  $L$ :  $-32\rightarrow 31$ ) have been collected with an internal  $R$ -value of 6.70%, a redundancy of 36.8, and with 98.85% coverage up to  $2\theta_{\text{max}} = 123.6^\circ$ . For the refinement, the Jana2006 program package was used. The goodness of fit of our crystal structure refinement amounts to 1.93, and the  $R$ - and weighted  $R$ -values amount to 2.29% and 5.79%, respectively. The refinement of the Ni and La occupancies yields an almost stoichiometric composition:  $\text{La}_{0.995}^{(12)}\text{Ni}_{1.000}^{(14)}\text{O}_3$ . The structural parameters are listed in Supplementary Table 1 and the crystal structure is visualized in Supplementary Figure 3.

**Composition determined by ICP and TG measurements.** Inductively coupled plasma optical emission spectroscopy measurements yields the following composition of  $\text{LaNiO}_3$ :  $\text{La}_{1.001(11)}\text{Ni}_{0.999(4)}\text{O}_{3.46}$ . Moreover, thermogravimetric measurements confirm an almost perfectly stoichiometric oxygen content with  $\delta = -0.002$ , see Supplementary Note 2.

**X-ray absorption spectroscopy.** X-ray absorption spectroscopy measurements have been performed at the 08B beamline of the National Synchrotron Radiation Research Center (NSRRC), Taiwan.

**Magnetization measurements.** The magnetic properties were studied using a Quantum Design Inc. MPMS-5XL SQUID magnetometer.

**Electrical conductivity.** The measurements of electrical resistivity and specific heat were carried out using a four-probe and a standard thermal relaxation calorimetric method in a Quantum Design Inc. Physical Property Measurement System.

**Neutron measurements.** Unpolarized neutron measurements have been performed at the IN8 and Thales Spectrometers at the ILL in Grenoble, France. For the elastic and inelastic measurements on the Thales spectrometer, a pyrolytic graphite (PG) (002) monochromator and analyzer as well as a velocity selector were used ( $k_i = 1.8 \text{ \AA}^{-1}$ ). On the IN8 spectrometer, very first quick inelastic scans have been made using a Si monochromator and a Flatcone analyzer array. Polarized neutron measurements have been performed at the IN12 Spectrometer at the ILL in Grenoble, France. A velocity selector was used for choosing the incident neutron wavevector  $k_i$  of  $2.25 \text{ \AA}^{-1}$ . A horizontally and vertically focussing PG (002) monochromator and Heusler (111) analyzer have been used. The incident neutron beam was polarized by a transmission polarizer (cavity) in the neutron guide and the measured flipping ratio amounts to 22.2.

**Data availability.** The data that support the findings of this study are available from the corresponding author upon reasonable request.

Received: 21 April 2017 Accepted: 4 December 2017

Published online: 03 January 2018

## References

- Torrance, J. B., Lacorre, P., Nazzari, A. I., Ansaldo, E. J. & Niedermayer, C. Systematic study of insulator-metal transitions in perovskites  $\text{RNiO}_3$  ( $R=\text{Pr}, \text{Nd}, \text{Sm}, \text{Eu}$ ) due to closing of charge-transfer gap. *Phys. Rev. B* **45**, 8209–8212 (1992).
- Medarde, M. L. Structural, magnetic and electronic properties of  $\text{RNiO}_3$  perovskites ( $R=\text{rare earth}$ ). *J. Phys. Condens. Matter* **9**, 1679–1707 (1997).
- Giovannetti, G., Kumar, S., Khomskii, D., Picozzi, S. & van den Brink, J. Multiferroicity in rare-earth nickelates  $\text{RNiO}_3$ . *Phys. Rev. Lett.* **103**, 156401 (2009).
- Mizokawa, T. & Fujimori, A. Unrestricted Hartree-Fock study of transition-metal oxides: spin and orbital ordering in perovskite-type lattice. *Phys. Rev. B* **51**, 12880–12883 (1995).
- Mizokawa, T., Khomskii, D. I. & Sawatzky, G. A. Spin and charge ordering in self-doped Mott insulators. *Phys. Rev. B* **61**, 11263–11266 (2000).
- Mazin, I. I. et al. Charge ordering as alternative to Jahn-Teller distortion. *Phys. Rev. Lett.* **98**, 176406 (2007).
- Lee, S. B., Chen, R. & Balents, L. Landau theory of charge and spin ordering in the nickelates. *Phys. Rev. Lett.* **106**, 016405 (2011).
- Park, H., Millis, A. J. & Marianetti, C. A. Site-selective Mott transition in rare-earth-element nickelates. *Phys. Rev. Lett.* **109**, 156402 (2012).
- Johnston, S., Mukherjee, A., Elfimov, I., Berciu, M. & Sawatzky, G. A. Charge disproportionation without charge transfer in the rare-earth-element nickelates as a possible mechanism for the metal-insulator transition. *Phys. Rev. Lett.* **112**, 106404 (2014).
- Jaramillo, R., Ha, S. D., Silevitch, D. M. & Ramanathan, S. Origins of bad-metal conductivity and the insulator metal transition in the rare-earth nickelates. *Nat. Phys.* **10**, 304–307 (2014).
- Subedi, A., Peil, O. E. & Georges, A. Low-energy description of the metal-insulator transition in the rare-earth nickelates. *Phys. Rev. B* **91**, 075128 (2015).
- Green, R. J., Haverkort, M. W. & Sawatzky, G. A. Bond disproportionation and dynamical charge fluctuations in the perovskite rare-earth nickelates. *Phys. Rev. B* **94**, 195127 (2016).
- Chaloupka, J. & Khaliullin, G. Orbital order and possible superconductivity in  $\text{LaNiO}_3/\text{LaMO}_3$  superlattices. *Phys. Rev. Lett.* **100**, 016404 (2008).
- Scherwitzl, R. et al. Metal-insulator transition in ultrathin  $\text{LaNiO}_3$  films. *Phys. Rev. Lett.* **106**, 246403 (2011).
- Middey, S. et al. Physics of ultrathin films and heterostructures of rare-earth nickelates. *Ann. Rev. Mater. Res.* **46**, 305–334 (2016).
- Xu, X. Q., Peng, J. L., Li, Z. Y., Ju, H. L. & Greene, R. L. Resistivity, thermopower, and susceptibility of  $\text{RNiO}_3$  ( $R=\text{La}, \text{Pr}$ ). *Phys. Rev. B* **48**, 1112–1118 (1992).
- Zhou, J. S., Marshall, L. G. & Goodenough, J. B. Mass enhancement versus stoner enhancement in strongly correlated metallic perovskites:  $\text{LaNiO}_3$  and  $\text{LaCuO}_3$ . *Phys. Rev. B* **89**, 245138 (2014).
- Hu, Z. et al. Hole distribution between the  $\text{Ni}3d$  and  $\text{O}2p$  orbitals in  $\text{Nd}_{2-x}\text{Sr}_x\text{NiO}_{4-\delta}$ . *Phys. Rev. B* **61**, 3739–3744 (2000).
- Hu, Z. et al. On the electronic structure of  $\text{Cu(III)}$  and  $\text{Ni(III)}$  in  $\text{La}_2\text{Li}_{1/2}\text{Cu}_{1/2}\text{O}_4$ ,  $\text{Nd}_2\text{Li}_{1/2}\text{Ni}_{1/2}\text{O}_4$ , and  $\text{Cs}_2\text{KCuF}_6$ . *Chem. Phys.* **232**, 63–74 (1998).
- Zhang, J., Zheng, H., Ren, Y. & Mitchell, J. F. High-pressure floating-zone growth of perovskite nickelate  $\text{LaNiO}_3$  single crystals. *Cryst. Growth Des.* **17**, 2730–2735 (2017).
- Zhou, J.-S., Goodenough, J. B. & Dabrowski, B. Pressure-induced non-fermi-liquid behavior of  $\text{PrNiO}_3$ . *Phys. Rev. Lett.* **94**, 226602 (2005).
- Sánchez, R. D. et al. Metal-insulator transition in oxygen-deficient  $\text{LaNiO}_{3-x}$  perovskites. *Phys. Rev. B* **54**, 16574–16578 (1996).
- Gayathri, N., Raychaudhuri, A. K., Xu, X. Q., Peng, J. L. & Greene, R. L. Electronic conduction in  $\text{LaNiO}_{3-\delta}$ : the dependence on the oxygen stoichiometry. *J. Phys. Condens. Matter* **10**, 1323–1338 (1998).
- García-Muñoz, J. L., Rodríguez-Carvajal, J. & Lacorre, P. Neutron-diffraction study of the magnetic ordering in the insulating regime of the perovskites  $\text{RNiO}_3$  ( $R=\text{Pr}$  and  $\text{Nd}$ ). *Phys. Rev. B* **50**, 978–992 (1993).
- Argyriou, D. N. et al. Lattice effects and magnetic structure in the layered colossal magnetoresistance manganite  $\text{La}_{2-2x}\text{Sr}_{1+2x}\text{Mn}_2\text{O}_7$  ( $x=0.3$ ). *Phys. Rev. B* **59**, 8695–8702 (1999).
- Yoshida, Y. et al. Crystal and magnetic structure of  $\text{Ca}_3\text{Ru}_2\text{O}_7$ . *Phys. Rev. B* **72**, 054412 (2005).
- Komarek, A. C. et al.  $\text{CaCrO}_3$ : an anomalous antiferromagnetic metallic oxide. *Phys. Rev. Lett.* **101**, 167204 (2008).
- Berlijn, T. et al. Itinerant antiferromagnetism in  $\text{RuO}_2$ . *Phys. Rev. Lett.* **118**, 077201 (2017).
- Khomskii, D. I. *Transition Metal Compounds* (Cambridge University Press, 2014).
- Li, B. et al. Insulating pockets in metallic  $\text{LaNiO}_3$ . *Adv. Electron. Mater.* **2**, 1500261 (2016).

31. Subedi, A. Breathing distortions in the metallic, antiferromagnetic phase of  $\text{LaNiO}_3$ . Preprint at <https://arxiv.org/abs/1708.08899> (2017).
32. Zhou, J. S., Goodenough, J. B. & Dabrowski, B. Transition from Curie-Weiss to enhanced Pauli paramagnetism in  $\text{RNiO}_3$  ( $R=\text{La, Pr, ... Gd}$ ). *Phys. Rev. B* **67**, 020404 (2003).

## Acknowledgements

We thank O. Stockert, G.A. Sawatzky, and A. Fujimori for valuable discussions. We thank M. Schmidt and his team for thermogravimetric measurements. We thank G. Auffermann and her team for ICP measurements. We acknowledge A. Todorova and G. Ryu for support on preparatory work. We thank the team of H. Borrmann for initial room-temperature powder X-ray diffraction measurements used for phase analysis. The research in Dresden and Cologne is partially supported by the Deutsche Forschungsgemeinschaft through FOR 1346.

## Author contributions

Conceiving experiments and project management: A.C.K.; physical experiments: H.G., Z. W.L., L.Z., Z.H., C.F.C., C.-Y.K., W.S., A.P., T.W.P., O.S., and A.C.K.; chemical synthesis: A.C.K. and H.G.; interpretation and manuscript writing: A.C.K., D.I.K., and L.H.T.

## Additional information

**Supplementary Information** accompanies this paper at <https://doi.org/10.1038/s41467-017-02524-x>.

**Competing interests:** The authors declare no competing financial interests.

**Reprints and permission** information is available online at <http://npg.nature.com/reprintsandpermissions/>

**Publisher's note:** Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.



**Open Access** This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons license and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this license, visit <http://creativecommons.org/licenses/by/4.0/>.

© The Author(s) 2017