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Indirect chiral magnetic exchange through Dzyaloshinskii–Moriya-enhanced RKKY interactions in manganese oxide chains on Ir(100)

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Localized electron spins can couple magnetically via the Ruderman–Kittel–Kasuya–Yosida interaction even if their wave functions lack direct overlap. Theory predicts that spin–orbit scattering leads to a Dzyaloshinskii–Moriya type enhancement of this indirect exchange interaction, giving rise to chiral exchange terms. Here we present a combined spin-polarized scanning tunneling microscopy, angle-resolved photoemission, and density functional theory study of MnO₂ chains on Ir(100). Whereas we find antiferromagnetic Mn–Mn coupling along the chain, the inter-chain coupling across the non-magnetic Ir substrate turns out to be chiral with a 120° rotation between adjacent MnO₂ chains. Calculations reveal that the Dzyaloshinskii–Moriya interaction results in spin spirals with a periodicity in agreement with experiment. Our findings confirm the existence of indirect chiral magnetic exchange, potentially giving rise to exotic phenomena, such as chiral spin-liquid states in spin ice systems or the emergence of new quasiparticles.

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The concept of the Ruderman–Kittel–Kasuya–Yosida (RKKY) interaction^{1,2} has successfully been applied to explain the magnetic properties of numerous indirectly coupled material systems which cannot be properly described by direct Heisenberg exchange. Prominent examples are the rare-earth metals with their partially filled but highly localized 4f shell^{3,4} or magnetic multilayers separated by non-magnetic metallic spacers. Since spin–orbit-related effects play no significant role in conventional RKKY, practical realizations are largely limited to collinear coupling terms, where—depending on spacer thickness and Fermi wave length—the relative magnetic orientation is either parallel or antiparallel⁵. Nevertheless, the giant magneto-resistance effect of layered magnetic materials is widely used in spin valve applications for field sensors or magnetic read heads⁶.

Theory predicts that spin–orbit scattering leads to a Dzyaloshinskii–Moriya^{7,8} type enhancement of the RKKY interaction^{9,10}, or DME-RKKY in short, giving rise to chiral exchange terms. First evidence of indirect chiral magnetic exchange in layered structures was obtained from magnetic field-dependent neutron diffraction studies Dy/Y superlattices^{11,12}. Further experimental evidence of DME-RKKY is essentially limited to non-collinear spin structures observed in surface-deposited clusters^{13–16}.

Here we report on the direct observation of chiral magnetic order between MnO₂ chains which is mediated by RKKY interaction via conduction electrons of the Ir substrate. The strong spin–orbit coupling in Ir leads to an appreciable DMI, resulting in a chiral spin spiral with a 120° rotation between adjacent MnO₂ chains.

Results

Structural and electronic properties. The growth and structural properties of self-organized transition metal oxide (TMO) chains on Ir(001) have recently been studied by means of STM and low-energy electron diffraction (LEED)¹⁷. It has been shown that many TMOs form extended (3 × 1)-ordered domains. Depending on the particular transition metal element, various intra-chain spin structures (along the chain) were predicted by DFT calculations¹⁷, ranging from a non-magnetic NiO₂, over ferromagnetically ordered (FM) CoO₂, to antiferromagnetic (AFM) FeO₂ and MnO₂ chains. In contrast, only a very weak inter-chain magnetic coupling between adjacent chains across the Ir(001) substrate was predicted¹⁷, too weak to result in spontaneous, permanent, and long-range magnetic order.

A topographic STM image of a typical MnO₂/Ir(001) surface is shown in Fig. 1a. Wide flat terraces are decorated by roughly rectangularly shaped islands of atomic height. Terraces and islands both exhibit stripes running along the [110] or the $\bar{1}\bar{1}0$ direction. These stripes originate from the self-organized growth of MnO₂ chains which leads to a (3 × 1) structural unit cell¹⁷. Some domain boundaries can be recognized which separate domains which differ either in stripe direction (left arrow in Fig. 1a) or by an incommensurate phase shift (right arrow). The higher resolution image of Fig. 1b was measured on a single (3 × 1) domain. The stripe periodicity of (840 ± 50) pm, corresponds well to the expected value of $3 \times a_{\text{Ir}} = 816$ pm, with the Ir lattice constant $a_{\text{Ir}} = 272$ pm¹⁷. Only 36 defects are observed (24 bright spots; 9 dumbbells, 2 point-like hole; 1 line defect), equivalent to a chain defect density below 0.35%.

The structure of the MnO₂ chains on Ir(001) as proposed by Ferstl et al.¹⁷ is schematically represented in Fig. 1c. Along the chains nearest-neighbor Mn atoms (yellow) are separated by two oxygen atoms (red). Interestingly, the MnO₂ chains sit above empty substrate rows, held in place by the oxygens atoms. DFT

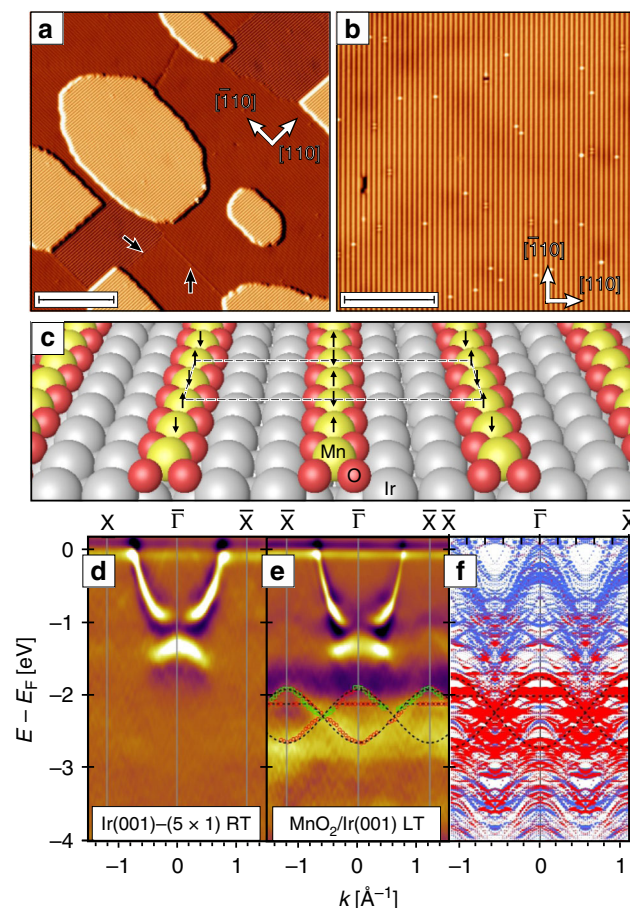


Fig. 1 Structural and electronic properties of MnO₂ on Ir(001). **a** Large scale STM image showing islands of monolayer height (scale bar: 25 nm). The entire surface including the islands are covered by chains along the [110] or $\bar{1}\bar{1}0$ direction. **b** Higher resolution STM image of the MnO₂ chains (scale bar: 15 nm). Scan parameters: $U = 1$ V, $I = 300$ pA. **c** Schematic model of the atomic structure and the AFM (6 × 2) spin arrangement predicted in ref. 17. **d, e** Second derivative ARPES data of the Ir(001)-(5 × 1) surface (at room temperature, RT) and for MnO₂/Ir(001) (at 13 K, LT) along the $\bar{\Gamma}$ – $\bar{X}$ direction ($h\nu = 150$ eV). Red and green symbols mark the dispersion of Mn-related states as determined from data taken at $h\nu = 150$ eV and 130 eV, respectively. Their periodicity is in agreement with the 2 × AFM order expected along the chains. Dashed lines are guides to the eye. **f** DFT band structure of MnO₂/Ir(001) exhibiting AFM order along the chains. Red and blue dots represent Mn and Ir states, respectively. The size of the symbols indicates the surface localization of the corresponding state. The same dashed lines shown in **e** (stretched by a factor 1.33 to consider the larger band width in DFT) agree well with surface-localized Mn states

calculations predicted an AFM coupling along the MnO₂ chains, favored by 27 meV per Mn pair with respect to a FM coupling¹⁷. Due to the large separation between adjacent chains a much weaker AFM coupling with an energy gain of 0.4 meV per Mn pair was found across the stripes, overall resulting in a rectangular (6 × 2) magnetic unit cell, sketched in Fig. 1c. ARPES measurements support the presence of an AFM intra-chain coupling. Figure 1d, e display second derivative ARPES spectra along the $\bar{\Gamma}$ – $\bar{X}$ axis of the Ir(001) surface Brillouin zone (SBZ) (corresponding to the Ir [110] direction) for the clean (5 × 1)-reconstructed substrate and the MnO₂/Ir(001) system, respectively. The second derivative is used to enhance the sensitivity to Mn-related states, which are broadened by the hybridization with the substrate.

The photon energy is chosen such that the Ir 5*d* signal is weak, except for some bulk bands dispersing symmetrically about $\bar{\Gamma}$ within 1.6 eV below E_F . Upon formation of the MnO₂ chains new states appear between −1.9 and −2.9 eV (Fig. 1e). The peak positions of Mn-related states are marked by red ($h\nu = 150$ eV) and green symbols ($h\nu = 130$ eV). Two sinusoidal dashed lines having maxima and minima at the $\bar{\Gamma}$ and $\bar{X}$ points are guides to the eye connecting the dispersive states. Flat states below the maxima are connected by dashed segments (see Supplementary Note 1 for bare data). These lines are compared with first-principles electronic structure calculations of AFM MnO₂ chains on Ir(001) oriented along the *y*-direction (Fig. 1f). They match well with the energy position of surface-localized Mn bands, but the experimental data display a smaller band width than DFT by a factor of 1.33, probably due to correlation effects. As detailed in the Supplementary Note 2 the sinusoidal bands mainly consist of d_{yz} and $d_{x^2-y^2}$ states, whereas the flat band is dominated by states with d_{zx} and d_{z^2} character. This observation suggests the presence of a $2\times$ periodicity, which turns the $\bar{X}$ points of the original SBZ into $\bar{\Gamma}$ points of the reduced SBZ, as expected for an AFM supercell. Other features located between −1 and −1.5 eV can be interpreted as surface umklapps of the Ir bulk band near $\bar{\Gamma}$ that repeat according to the AFM supercell. We recall here that the ARPES measurements of Fig. 1e average over both directions parallel and perpendicular to the chains, as a consequence of the domain structure of the system (Fig. 1a). The inter-chain coupling which results in a $9\times$ magnetic unit cell (see below) is expected to be much weaker than the direct intra-chain AFM coupling and does not give rise to dispersive features in the ARPES data.

Spin-polarized scanning tunneling microscopy. Figure 2a shows an atomic scale STM image of MnO₂ chains on Ir(001) taken with a non-magnetic W tip. The data show a structural (3×1) unit cell (black box) and nicely reproduce earlier measurements¹⁷. The black lines in Fig. 2c show line profiles taken along the three adjacent MnO₂ chains indicated by arrows in Fig. 2a. They exhibit a periodicity (287 ± 20) pm, agreeing well with a_{Ir} , i.e., the Mn–Mn inter-atomic distance expected along the chain. Note, that within the noise level achievable in our setup the corrugation amplitude of (1.9 ± 0.1) pm remains constant.

As we will describe in the following, our spin-polarized STM experiments exhibit some additional contrasts which allow to elucidate the spin structure of the MnO₂ chains on Ir(001). Figure 2b shows an SP-STM image scanned with an in-plane sensitive Cr-coated W tip. Comparison with the spin-averaged data presented in Fig. 2a reveals two qualitative differences: (i) The periodicity measured with magnetic tips along the chains is longer and (ii) the contrast observed on different MnO₂ chains is not constant but becomes significantly smaller for every third chain. Again we analyzed line profiles taken along three adjacent MnO₂ chains in between the colored arrows in Fig. 2b. These data, which are plotted in the bottom part of Fig. 2c, immediately illustrates a doubling of the periodicity, $2a_{\text{Ir}}$. This SP-STM contrast is characteristic for alternating spins¹⁸ and consistent with the proposed AFM Mn–Mn coupling along the chains¹⁷. Furthermore, the spin-polarized data reveal a systematic variation of the corrugation. For example, the blue line trace exhibits a corrugation of 6.1 pm, in contrast to 3.2 pm (green) and 5.4 pm (red) for the two adjacent MnO₂ chains. Finally, there is also a distinct phase relation between the chains. Comparing the three colored traces plotted in Fig. 2c a π phase shift between the blue and the green trace becomes apparent, whereas no phase shift occurs between the green and the red trace.

Modeling the magnetic contrast. The SP-STM contrast observed on MnO₂ chains can semi-quantitatively be understood by

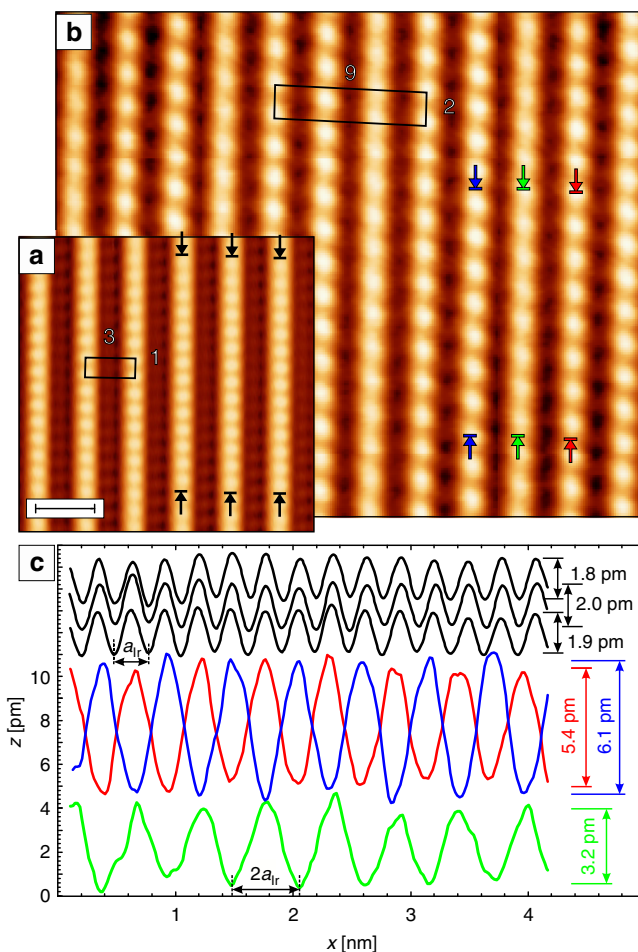


Fig. 2 Atomic resolution scans of MnO₂ chains on Ir(001). **a** A (3×1) structural unit cell is observed with a non-magnetic W tip (scale bar: 1 nm). **b** With a Cr-coated W tip the magnetic (9×2) unit cell is resolved. **c** Line profiles drawn along the stripes at the positions indicated by arrows measured with the W (black) and the Cr-coated (colored) probe tip. Spin-resolved line sections differ in periodicity, phase, and amplitude from their spin-averaged counterparts (see main text for details). Scan parameters: **a** $U = -500$ mV, $I = 3$ nA; **b** $U = 100$ mV, $I = 300$ pA

assuming an AFM coupling along the chain and a chiral 120° coupling between adjacent chains. This spin configuration which leads to a (9×2) magnetic unit cell is schematically sketched in Fig. 3a. As mentioned above the AFM Mn–Mn intra-chain coupling can directly be concluded from the doubling of the periodicity in SP-STM as compared to spin-averaged data (a_{Ir}) (cf. Fig. 2c). To also explain the corrugation amplitudes and their phase we need to consider that the magnetic corrugation in SP-STM, Δz_{SP} , depends on the cosine of the angle θ included between the magnetization directions of the tip and the sample,

$$\Delta z_{\text{SP}} \propto P_t \cdot P_s \cdot \cos\theta, \quad (1)$$

with P_t and P_s being the spin polarization of tip and sample, respectively. The expected magnetic contrast can be deduced from the scheme in Fig. 3b. It represents three sample magnetization directions which are rotated by 120° to another as symbolized by colored arrows. According to Eq. (1), Δz_{SP} is given by the projection of the sample magnetization onto the tip magnetization. Therefore, the maximum Δz_{SP} is expected for a sample magnetization which is almost collinear to the tip magnetization (represented by the black arrow). As symbolized by the lightly colored triangle this condition is fulfilled for the blue arrow in

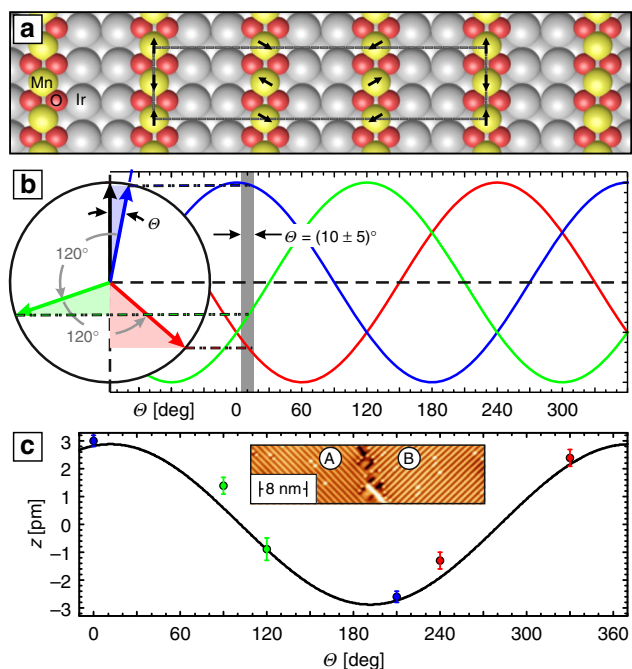


Fig. 3 Interpretation of SP-STM results by a chiral inter-chain exchange. **a** Schematic model of the chiral (9×2) spin structure of MnO₂/Ir(001). **b** Sketch of the SP-STM signal expected on a sample with three domains rotated by 120° to another (see text and Eq. 1 for details). The corrugations data of Fig. 2b are best fit with an angle $\theta = (10 \pm 5)^\circ$ between tip magnetization (black) and the nearest domain (blue). **c** Corrugation values determined by fitting line profiles measured on two adjacent domains A and B (inset) as measured with an in-plane tip. The error bar represents the residual sum of squares. The black curve represents the cosine expected for an in-plane rotating spin spiral

Fig. 3b (offset by angle θ). In this situation it is unavoidable that the projection of the other two arrows points into the direction opposite to the black arrow. This can also be verified by inspecting the right part of Fig. 3b, where we plot three cosine functions shifted by 120° .

In other words, from the fundamental principles of SP-STM it follows that (i) whenever we obtain a large contrast on one AFM spin chain the other two chains with spin quantization axes rotated by $\pm 120^\circ$ must exhibit a magnetic corrugation which is phase-shifted with respect to the high-contrast row. Furthermore, (ii) even if θ is relatively small, one of the $\pm 120^\circ$ -rows exhibits a much lower magnetic contrast since the $\cos\theta$ term in Eq. (1) is close to zero. As marked by a gray box in Fig. 3b, the corrugations measured in Fig. 2b can nicely be explained by a tip which is rotated by $\theta = (10 \pm 5)^\circ$ with respect to the (blue) domain. As discussed in detail in the Supplementary Note 3 we have performed various SP-STM measurement to identify the spin orientation of the MnO₂ chains. Figure 3c shows the magnetic corrugation measured on the three MnO₂ chains (blue, red, and green) of two domains (A and B; see inset), respectively, which are rotated by 90° with respect to another. The black line is the corrugation expected for a spin spiral rotating in the surface plane. The very good agreement with our experimental data suggests an inter-chain coupling characterized by an in-plane 120° rotation of the azimuthal spin orientation. In order to verify if this spin order is indeed chiral we determined the rotational sense of seven separate MnO₂ domains (see Supplementary Note 4). Indeed, our SP-STM measurements show that all domains exhibit the same rotational sense, a result which is highly unlikely under non-chiral conditions ($<2\%$).

Density-functional theory calculations. To obtain some insights into the origin of the observed magnetic structures we performed DFT calculations (see Supplementary Note 5 for details). The preferred magnetic ordering along the chains was found to be AFM, in agreement with earlier calculations¹⁷ and our experimental results (see Fig. 2b). Whereas a weak AFM inter-chain coupling of 0.4 meV per Mn atom was found previously¹⁷, our calculations performed at a much denser k -point sampling (using a 24×36 Monkhorst-Pack grid) leads to a weak FM coupling of 1.7 meV. We calculated flat spin spirals with various wave vectors q , where the FM (AFM) state corresponds to $q = 0$ ($q = 0.5$) in units of $2/3a_{\text{Ir}}$. From our spin spiral calculations, Fig. 4, we can see that symmetric (Heisenberg-type; blue) exchange interactions lead to a flat dispersion, without any minimum at finite q . Our results show that the DMI is largest for a spin spiral with $q = 1/3$, i.e., the modulation vector found experimentally, lowering the total energy by about 0.3 meV. We have to note, however, that the theoretically obtained Dzyaloshinskii vector, $\mathbf{D}$, points along the chain direction, whereas experiments suggest an in-plane spin spiral, corresponding to $\mathbf{D}$ along the surface normal. To explain the experimentally observed unique rotation sense, there must be a significant out-of-plane component of $\mathbf{D}$, e.g., due to a structural distortion that removes the $[110]$ mirror plane. A similar mechanism has recently been shown to exist for zigzag Co/Ir(001)¹³. Indeed, some hints of a potential distortion of the (3×1) structural unit cell can not only be recognized in our data (Fig. 2a), but also in the data published by Ferstl and co-workers (see Figs. 1c, S2 in ref. 17, and Supplementary Note 6 of this article). In either case atomic resolution data recorded on magnetic TMO chains with non-magnetic tips show some oblique distortion of the expected rectangular surface unit cell. Although this does not directly lead to a non-vanishing effective perpendicular $\mathbf{D}$, it indicates that some structural details or relaxation effects due to the finite size of structural domain still need to be resolved (see Supplementary Note 5).

Discussion

Recent research indicates that spin-orbit coupling supports an effective spin transfer torque^{19–21} which is particularly important in applications. Different mechanisms have been proposed to explain the relatively low current thresholds necessary to drive skyrmions or chiral domain walls, including inhomogeneous spin currents¹⁹, Rashba fields, the spin Hall effect²⁰, the DMI, or a combination of the latter two²¹. We speculate that a DME-RKKY interaction-induced indirect chiral magnetic exchange may also lead to an extreme reduction of the required current density in layered magnetic structures. It remains to be investigated whether a chiral magnetic interlayer coupling as it has been observed in Dy/Y superlattices^{11,12} can also be found in other material combinations with strongly spin-orbit coupled non-magnetic spacer layers. In more general terms, DME-RKKY interaction may give rise to rather exotic phenomena, such as chiral spin-liquid states in spin ice systems^{22,23} or the emergence of new quasiparticles due to the trapping of single electrons in self-induced skyrmion spin textures²⁴.

In summary, we have investigated the intra-chain and inter-chain magnetic coupling of the quasi one-dimensional system of structurally (3×1) -ordered MnO₂ on Ir(001) by spin-polarized scanning tunneling microscopy, angle-resolved photoemission, and density functional theory. Both experimental methods confirm an antiferromagnetic order along the chains, as predicted earlier¹⁷. In addition, spin-polarized scanning tunneling microscopy reveals a chiral 120° rotation between adjacent MnO₂ chains, resulting in a (9×2) magnetic unit cell. Density functional theory finds that a Dzyaloshinskii–Moriya type enhancement of the RKKY interaction indeed leads to chiral interchain coupling

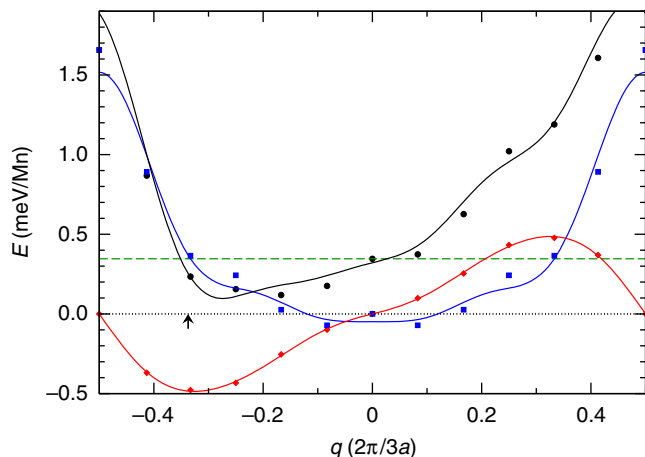


Fig. 4 Results of DFT calculations. Total energy of AFM ordered MnO_2 chains on Ir(001) where the spin directions from chain to chain rotate by an angle of $3a_{\text{Ir}}$. The spins form a cycloidal spin spiral propagating perpendicular to the chains with the rotation axis parallel to the chain direction. The contribution from symmetric exchange interaction is shown by blue squares, from the DMI by red diamonds, and from the magneto-crystalline anisotropy by the dashed line. The sum of these terms is marked by black circles. The blue line is a fit of the symmetric exchange interaction with $\cos(2nq)$, likewise the DMI was fitted (red line) with $\sin(2nq)$, both up to $n = 4$

with a periodicity in agreement with experiment. However, the orientation of the Dzyaloshinskii vector $\mathbf{D}$ remains to be clarified. Whereas experimental results suggest a perpendicular $\mathbf{D}$, theory predicts a $\mathbf{D}$ vector oriented along the chains.

Methods

Sample preparation. Sample preparation procedures closely follow published recipes¹⁷. Initial Ir(001) preparation comprises cycles of ion-sputtering (1 keV, Ar⁺, $\approx 2\mu\text{A}$) followed by annealing to 1400 K in an oxygen atmosphere. The pressure gauge indicates a background pressure $p_{\text{O}_2} \approx 1 \times 10^{-8}$ mbar, but since the gas nozzle is located a few cm above the sample the local oxygen pressure is assumed to be about two orders of magnitude higher. We obtain the (5×1) reconstruction characteristic for clean Ir(100) by an annealing cycle without oxygen²⁵. Oxidizing this surface again in $p_{\text{O}_2} \approx 1 \times 10^{-8}$ mbar at $T_s \approx 850$ K leads to the oxygen-terminated Ir(100)-(2 × 1) reconstruction^{17,26,27}. It served as a substrate for the deposition of 0.33 monolayers (ML) of Mn at room temperature, followed by final annealing ($T_s \approx 1050$ K) under oxygen atmosphere.

Scanning tunneling microscopy (STM). STM experiments were performed in a two-chamber ultra-high vacuum (UHV) system (base pressure $p \leq 5 \times 10^{-11}$ mbar) equipped with a home-built low-temperature scanning tunneling microscope (LT-STM) (operation temperature $T = 5.5$ K). We used electro-chemically etched polycrystalline W tips which were flashed by electron bombardment and coated with Fe or Cr for SP-STM measurements²⁸.

Angular-resolved photoemission spectroscopy (ARPES). ARPES data were acquired at 130 and 150 eV photon energy at the VUV photoemission beamline (Elettra, Trieste). These photon energies are close to the Cooper minimum of the Ir 5d photoemission (PE) cross section²⁹ and enhances the PE signal of Mn 3d states (chains) with respect to the overlapping Ir 5d states (substrate). The spot of the synchrotron light on the sample ($500 \mu\text{m} \times 200 \mu\text{m}$) is much larger than the typical size of domains with parallel MnO_2 chains. Thus, ARPES provides a space-averaged signal over the two perpendicular orientations of the MnO_2 chains. The energy and momentum resolutions were set to 15 meV and $0.02 \text{ \AA}^{-1}$, respectively.

DFT calculations. Non-collinear DFT calculations were performed using the full-potential linearized augmented plane wave method as implemented in the Fleur code³⁰. We set up a seven layer film in a (3×2) unit cell as described in ref. 17, using the local density approximation³¹ with Hubbard U corrections³² on the Mn d states ($U = 2.7$ eV, $J = 1.2$ eV). We confirmed that these values put the Mn d states about 2.2 eV below the Fermi level, in good agreement with the ARPES data presented in Fig. 1d. We used the generalized Bloch theorem to calculate the spin

spiral structures and included spin-orbit coupling in first order perturbation theory to estimate the strength of the DMI in this system³³.

Data availability

The data that support these findings of this study are available on request from M.S. (STM), P.M. (ARPES), and G.B. (theory).

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Author contributions

M.S., R.C., M.V. and J.K. performed and analyzed STM measurements. P.M., P.M.S., A.K.K. and C.C. performed ARPES experiments and analyzed the data. G.B., A.B. and S. B. provided the theoretical framework. All authors discussed the results and contributed in writing the manuscript.

Additional information

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