

Model-free description of polymer-coated gold nanoparticle dynamics in aqueous solutions obtained by Bayesian analysis of neutron spin echo data

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We present a neutron spin echo study of the nanosecond dynamics of polyethylene glycol (PEG) functionalized nanosized gold particles dissolved in D₂O at two temperatures and two different PEG molecular weights (400D and 2000D). The analysis of the neutron spin echo data was performed by applying a Bayesian approach to the description of time correlation function decays in terms of exponential terms, recently proved to be theoretically rigorous. This approach, which addresses in a direct way the fundamental issue of model choice in any dynamical analysis, provides here a guide to the most statistically supported way to follow the decay of the intermediate scattering functions $I(Q, t)$ by basing on statistical grounds the choice of the number of terms required for the description of the nanosecond dynamics of the studied systems. Then, the presented analysis avoids from the start resorting to a preselected framework and can be considered as model free. By comparing the results of PEG-coated nanoparticles with those obtained in PEG2000 solutions, we were able to disentangle the translational diffusion of the nanoparticles from the internal dynamics of the polymer grafted to them, and to show that the polymer corona relaxation follows a pure exponential decay in agreement with the behavior predicted by coarse grained molecular dynamics simulations and theoretical models. This methodology has one further advantage: in the presence of a complex dynamical scenario, $I(Q, t)$ is often described in terms of the Kohlrausch-Williams-Watts function that can implicitly represent a distribution of relaxation times. By choosing to describe the $I(Q, t)$ as a sum of exponential functions and with the support of the Bayesian approach, we can explicitly determine when a finer-structure analysis of the dynamical complexity of the system exists according to the available data without the risk of overparametrization. The approach presented here is an effective tool that can be used in general to provide an unbiased interpretation of neutron spin echo data or whenever spectroscopy techniques yield time relaxation data curves.

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I. INTRODUCTION

Considerable scientific effort has been devoted in the last years to the study of gold nanoparticles (NP), in view of many important properties peculiar of these systems including quantum size effects and single electron transitions [1]. Due to their small size (1–20 nm), the electronic bands of the nanoparticles follow quantum-mechanical rules and exhibit a strong dependence of the physical properties on the size and shape of the particles [2]. These properties are further highly dependent upon the presence of an organic ligand coating or capping layer and upon the spatial relationship of one particle to another.

In order to contrast and control their tendency to aggregate, different methods have been developed to synthesize and stabilize such particles. Notable among these is the method developed by Brust *et al.* [3]. This method has had considerable

impact since it has allowed the production of gold NP (AuNP) of controlled size and reduced size dispersity stabilized by alkanethiols. This has enabled a series of functionalizations that allow using the AuNPs in a number of applications including molecular recognition [4,5] and specific binding to biomolecules [6] with significant impact in biological and biomedical studies [7,8].

Functionalization with thiol-terminated polymers provides great flexibility in modulating the properties of the AuNPs. It was shown that AuNPs protected with 4-(N,N-dimethylamino)pyridine (DMAP) are a convenient precursor of thiol-capped AuNPs. DMAP can be conveniently exchanged with thiol-terminated polymers to obtain AuNP that are soluble in organic solvents or water by choosing the appropriate thiol-terminated polymer including poly(ethylene oxide) (PEG) or polystyrene (PS) over a range of molecular weights [9].

The properties of the organic capping layer, which protects the AuNPs, impact the way the AuNPs interact with other

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systems such as a surrounding polymer matrix [10], other NPs in complex three-dimensional (3D) architectures [11], or a biomimetic assembly [12]. For practical applications, it is crucial to have a very precise knowledge of the molecular details and properties of this protective layer surrounding the NP.

In a previous study, a structural characterization of 5-nm PEG functionalized AuNPs was achieved via a combination of experimental methods including neutron and x-ray small angle scattering, transmission electron microscopy, density, and thermogravimetric measurements [13]. The functionalization was obtained through thiol derivatization of PEG2000 (45 units). The confinement of the PEG molecule on the highly curved surface of a AuNP was found to considerably affect the conformation and the hydration of the polymer as a function of the distance from the surface of the gold core.

Here, we extend this static study to the dynamical properties of these PEG functionalized AuNPs. The study of the dynamical processes in nanosized objects like micelles, proteins, and nanoparticle dispersions presents some difficulties due to the presence of multiscale dynamics occurring at similar time and range of length scales. Methods like dynamic light scattering (DLS) are useful to measure the translational diffusion of the whole object under study. However, since the wavelength of light is larger than the dimension of these nanosized objects, internal dynamics cannot be explored with this technique. Quasielastic neutron scattering (QENS) and, in particular, neutron spin echo spectroscopy (NSE) have been shown to be an excellent tool to probe the dynamical processes in micellar dispersions [14–16]. Thanks to its wide dynamical range, NSE can probe a variety of dynamical processes, from translational diffusion of the whole micelle to the internal dynamics of the polymer chains.

The polymer-coated NPs used in this study (see Fig. 1) have sizes of the same order of those of the micellar aggregates studied in other investigations (50–200 Å) [14–16]. However, they differ significantly in various aspects. First, unlike the core of a micellar aggregate, the gold core of the NP is a solid. The internal dynamics of gold in the solid state include atomic vibrations and phonons that occur at much faster timescales than those detectable with NSE. Hence, differently from the micellar aggregate comparators, in the systems reported here the dynamical processes present in the polymer corona and in the NPs core occur at significantly

different timescales. Moreover, the nonionic surfactants, or block copolymers of the micelles, aggregate in response to the hydrophobic effect, which favors the confinement of the hydrophobic moieties in the oily bulk. These aggregates are dynamic entities, as the individual surfactant molecules have a finite solubility in the solvent phase and can leave an aggregate to become part of another one, following a kinetic equilibrium. PEG-grafted gold NPs, on the other hand, involve a strong Au-S bond (about 160–200 kJ/mol) between the thiol-derivatized PEG to the metal surface. The exchange of material between different NPs, if it in fact occurs, does so at an exceedingly lower rate under ambient condition. NSE was also used to study NPs with various functionalizations dispersed in polymeric materials to elucidate the effect of the NP dispersion on the dynamical properties of the NPs-polymer composite [17–21]. These studies differ from ours in that the NPs are dispersed in a high molecular weight polymer matrix, whereas the NPs used in our experiments were dispersed in a low molecular weight solvent.

The analysis of the NSE data was undertaken by adopting an approach composed by two building blocks:

(1) We expressed the NSE intermediate scattering function $I(Q, t)$ as a sum of exponentials following the results of general theoretical work that proves that any time correlation function can be expressed as the sum of infinite (generally complex) exponential functions [22–24]. This provided an unbiased model-free expression to interpret the experimental NSE results.

(2) The analysis of the $I(Q, t)$ was supported by a Bayesian approach implemented through the use of a Monte Carlo Markov chain (MCMC) algorithm integrated with a reversible jump (RJ) [25] option and successfully applied recently to Brillouin inelastic neutron scattering [26] and inelastic x-ray scattering [27] data.

This methodology addresses one of the most important issues in the analysis of inelastic and quasielastic neutron scattering data: the arbitrariness in the choice of the number of components needed to describe the experimental dynamic structure factor. In NSE this issue translates into how many components one chooses to describe the decay of the intermediate scattering function (the time Fourier transform of the dynamic structure factor). In this approach, the number of components is treated as a parameter to be estimated. Since inferential results from this approach are in the form of probabilities, the final choice of the number of components resides on a solid probabilistic basis. The main objective of this work is then to show that the combination of points 1 and 2 above allows a precise and theory complying representation of macromolecular polymeric systems in solutions and yields an unbiased and model-free description of the dynamics of these systems.

Furthermore, it is common practice in the analysis of NSE data to use a stretched exponential function to model the decay of the intermediate scattering function. The use of this function is in some circumstances a heuristic way to account for an ensemble of dynamical processes occurring in the sample with a given distribution of relaxation times and the stretching parameter might measure, in some fashion, the difficulty in describing such a complex dynamics without, however, being able to give account for this complexity. It

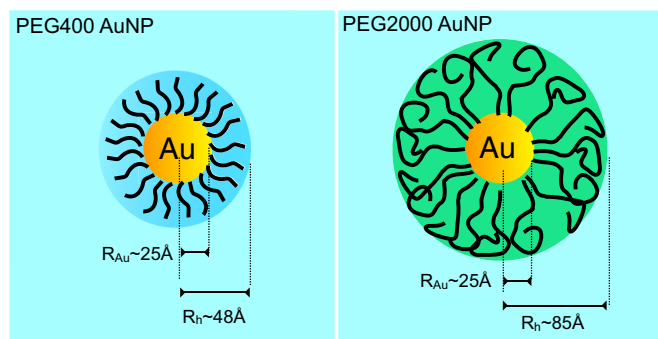


FIG. 1. Sketch of the PEG400 and PEG2000 functionalized NPs. An indication of the characteristic sizes of the Au core (R_{Au}) and of the hydrodynamic radius R_h are also shown.

is also the result of model based theoretical description of polymer dynamics such as the Zimm or the Rouse model. Nevertheless, all possible descriptions of the experimental $I(Q, t)$ of the system under study should eventually satisfy the above general property [22–24]. This principle has been exploited to give excellent descriptions of time correlation functions and spectra relevant to simpler liquids [28–30]. Here, by combining it with the Bayesian approach, we analyze the data in an alternative unbiased and model-free way by assigning a *fine structure* of relaxation processes to an otherwise undistinguished representation of the decay of the time correlation function as it is provided by the beta-stretching modeling. Furthermore and most importantly, the Bayesian method used here is able to provide the most statistically reliable description of the system supported by the experimental data without any risk of overparametrization.

We will also show that our unbiased method will provide results for the polymer corona in agreement with those predicted by coarse grained MD simulations and traditional theoretical models of polymer science.

II. MATERIALS AND METHODS

Polyethylene-glycol-coated gold nanoparticles (PEG AuNPs) were synthesized with the same procedure described in our previous work [13]. Briefly, tetraoctylammonium bromide-Au and DMAP-AuNPs were synthesized using Gittins and Caruso's procedure [31,32]. The Au-NPs were subjected to place-exchange reaction of presynthesized DMAP stabilized NPs with thiol-functionalized PEG2000SH and PEG400SH (see Fig. 1). The gold core had a radius around 25 Å with a polydispersion of 20% as a result of TEM analysis [9]. However, the indetermination in the Au core radius as a result of small angle neutron scattering was around 2% [13]. The detailed synthesis and characterization are described elsewhere [9]. PEG2000SH was purchased from Polymer Source Inc. Dorval, QC, Canada ($M_w \sim 2000$) and PEG400SH ($M_w \sim 400$) was received from Polypure Oslo, Norway. PEG2000 and PEG400 homopolymers were purchased from Sigma and used without further purification. The translational diffusion coefficients of the NPs were obtained by DLS. These measurements were carried out using an ALV CGS-3 compact goniometer equipped with a HeNe laser with a wavelength of 632.8 nm, a 22-mW output power, and an ALV LSE-5004 correlator. Samples were measured at a scattering angle of 90°, while the sample temperature was controlled via an external water bath circulator. The small angle neutron scattering profile for the PEG2000 AuNP was measured on the D11 SANS instrument at the Institut Laue-Langevin with a standard configuration. The neutron wavelength was 6 Å and two detector distances (10.5 and 1.5 m) were used. This provided a Q range between 6×10^{-3} and 0.43 Å^{-1} . The PEG2000 AuNPs were dissolved in D₂O at a concentration of 1% in weight.

A. Neutron spin echo spectroscopy

NSE experiments were performed at the IN15 and IN11 neutron spin echo spectrometers at the Institut Laue-Langevin in Grenoble. NSE measurements were carried out in solu-

TABLE I. Instrumental setups used for the various samples.

Sample	C (wt.%)	Instrument	λ (Å)	NSE time (ns)	Q range (Å ⁻¹)
PEG2000 AuNP	5	IN15	10, 16	0.3–206	0.02–0.21
PEG400 AuNP	2	IN15	8, 16	0.18–206	0.02–0.25
PEG2000	10	IN15	8, 16.8	0.18–240	0.02–0.25
PEG2000	10	IN11	8.5	0.02–5	0.09–0.27
PEG400	5	IN15	8, 16	0.18–206	0.02–0.25

tions of PEG2000-SH and PEG400-SH coated nanoparticles dispersed in D₂O at concentrations of 5 wt.% and 2 wt.%, respectively, and in solutions of PEG2000 homopolymers (10 wt.%) and of PEG400 homopolymers (5 wt.%), also in D₂O. The NSE signal of D₂O alone was also measured in the same experimental conditions to allow for background subtraction. All measurements of $I(Q, t)$ were performed in a wide range of Q and time at two temperatures 280 and 318 K. Different setups were used depending on the sample and the spectrometer. A complete summary of the measured samples and experimental setups is given in Table I.

B. Theory

We briefly present in this section the expressions on which are based the models adopted for analyzing the neutron intermediate scattering functions (nISF) $I(Q, t)$ measured in the neutron spin echo experiments.

A number of customary approximations are generally adopted to obtain a workable expression of $I(Q, t)$. First of all, the concentration of the nanoparticles in the solutions is supposed to be low enough to neglect any correlation between atoms of different nanoparticles. Then, it is assumed that translational, rotational, and vibrational motions are decoupled. This hypothesis is strictly valid for very dilute solutions but it is frequently used for systems at higher concentrations [33] and quite plausible for our system as we will show later. Under these hypotheses, it can be shown that the nISF, after subtraction of the measured solvent contribution, can be expressed as

$$\begin{aligned}
 I(Q, t) &\simeq \exp(-DQ^2t) \frac{1}{N_{\text{at}}} \sum_{n,n'=1}^{N_{\text{at}}} \langle b_n^* b_{n'} \rangle \\
 &\quad \times \langle \exp\{i\vec{Q} \cdot [\vec{r}_n(t) - \vec{r}_{n'}(0)]\} \rangle \\
 &= \exp(-DQ^2t) I_{\text{int}}(Q, t),
 \end{aligned} \tag{1}$$

where D is the translational diffusion coefficient of the macromolecule, N_{at} is the total number of atoms in a nanoparticle, b_n is the scattering length of the n th atom, and $r_n(t)$ is the displacement of an atom from its equilibrium position due to vibration or localized diffusive motion [34], so that the last factor, labeled with the subscript “int”, describes the internal dynamics of the nanoparticle. As far as rotation is concerned, it is useful to recall that the combination of independent translational and rotational motions has been shown [35] to be satisfactorily described in the frequency domain by a single central Lorentzian line, though with a slightly increased width with respect to Fick's Lorentzian of pure translational

diffusion. This suggests that rotations can be accounted for by an effective, small, increase of the translational diffusion coefficient. However, in the specific case of the present samples, as reported in Sec. III, no such increase has been detected, suggesting that the rotational contribution to Eq. (1) is negligible.

In liquid samples, the position of a nucleus as $t \rightarrow \infty$ becomes uncorrelated with its initial value. Conversely, if the motion of the nucleus is well located in space, or confined to a particular volume (e.g., reorientations of chemical groups in a molecule), the correlation persists and tends to a nonvanishing asymptotic value at long times. The latter case applies to the internal dynamics of the polymer molecules grafted to the nanoparticles, therefore, also in the present case it is possible, at any time, to write $I_{\text{int}}(Q, t)$ as the sum of its time-independent and time-dependent parts:

$$I_{\text{int}}(Q, t) = I_{\text{int}}(Q, \infty) + I'_{\text{int}}(Q, t). \quad (2)$$

Here, the term $I_{\text{int}}(Q, \infty)$ represents a quantity proportional to the contribution of purely elastic scattering to the differential scattering cross section [36] and is related to the structure of the nanoparticle.

NSE experiments measure a normalized intermediate scattering function (nISF), which is normally written as

$$\frac{I(Q, t)}{I(Q, 0)} = \frac{I^{\text{coh}}(Q, t) - \frac{1}{3}I^{\text{incoh}}(Q, t)}{I^{\text{coh}}(Q, 0) - \frac{1}{3}I^{\text{incoh}}(Q, 0)}, \quad (3)$$

where, in general,

$$I^{\text{coh}}(Q, t) = \frac{1}{N_{\text{at}}} \sum_{n, n'=1}^{N_{\text{at}}} \langle b_n^{\text{coh}} b_{n'}^{\text{coh}} \rangle \langle e^{-i\vec{Q} \cdot \vec{r}_n(0)} e^{+i\vec{Q} \cdot \vec{r}_{n'}(t)} \rangle, \quad (4)$$

$$I^{\text{incoh}}(Q, t) = \frac{1}{N_{\text{at}}} \sum_{n=n'=1}^{N_{\text{at}}} \langle (b_n^{\text{incoh}})^2 \rangle \langle e^{-i\vec{Q} \cdot \vec{r}_n(0)} e^{+i\vec{Q} \cdot \vec{r}_n(t)} \rangle. \quad (5)$$

Both these terms are implicitly contained in Eq. (1) and contribute to $I(Q, t)$. In the experiment performed here, the coherent term prevails especially at low Q ($\leq 0.1 \text{ \AA}^{-1}$). In the region of higher Q the incoherent terms reached 40%–50% of the total contribution as it is shown in Fig. S1 in the Supplemental Material [37].

To further model $I'_{\text{int}}(Q, t)$, we refer to the already quoted general theoretical results [22–24], showing that any correlation function that decays to zero at long times is exactly represented by an infinite series of exponentials of generally complex argument, each exponential being weighted, in the series, by a generally complex amplitude. When the amplitude and time constant of an exponential term are complex, then the sum of such term with its complex conjugate (also present in the series) describes an exponentially damped oscillation. Differently, if amplitude and time constant are real, then the exponential term depicts a pure exponential decay [38]. It is important to note that the possibility of expressing the $I'_{\text{int}}(Q, t)$ as a sum of exponential terms is valid irrespectively of the fact that it represents, or is obtained from a coherent, incoherent scattering or a combination of the two.

Since we are dealing with quasielastic spectra, we can rather confidently assume that the normalized time correlation function $I'_{\text{int}}(Q, t)$ does not require complex modes in the series representation, and that it can be reliably described by real modes only. Moreover, in Refs. [28–30] it was shown that time correlation functions are most typically described with high accuracy by using only a small number (let us say k) of terms in the series, thus, finally we modeled the internal dynamics as

$$\frac{I_{\text{int}}(Q, t)}{I_{\text{int}}(Q, 0)} = \frac{I_{\text{int}}(Q, \infty)}{I_{\text{int}}(Q, 0)} + \left(1 - \frac{I_{\text{int}}(Q, \infty)}{I_{\text{int}}(Q, 0)}\right) \times \sum_{j=2}^k I_j \exp\left(-\frac{t}{\tau_{\text{int},j}}\right). \quad (6)$$

The reason why the index j starts at 2 will be clear immediately below. In Eq. (6), $\tau_{\text{int},j}$ and I_j are, respectively, the time constant and the amplitude of the j th decay channel considered in the series (we omitted their obvious dependence on Q), and $\sum_{j=2}^k I_j = 1$. The factor in front of the exponential series is imposed by the normalization condition $[I_{\text{int}}(Q, t)/I_{\text{int}}(Q, 0)]_{t=0} = 1$. Thus, by substituting Eq. (6) in Eq. (1), our system can actually be modeled as

$$\frac{I(Q, t)}{I(Q, 0)} \simeq \exp(-DQ^2t) \left[A_1 + \sum_{j=2}^k A_j \exp\left(-\frac{t}{\tau_{\text{int},j}}\right) \right], \quad (7)$$

with $A_1 = I_{\text{int}}(Q, \infty)/I_{\text{int}}(Q, 0)$, $A_j = (1 - A_1)I_j$ for $j \geq 2$, and the sum rule $\sum_{j=1}^k A_j = 1$. Defining $1/\tau_1 = DQ^2$ and $1/\tau_j = 1/\tau_1 + 1/\tau_{\text{int},j}$ for $j \geq 2$, Eq. (7) can be rewritten as

$$\frac{I(Q, t)}{I(Q, 0)} \simeq A_1 \exp\left(-\frac{t}{\tau_1}\right) + \sum_{j=2}^k A_j \exp\left(-\frac{t}{\tau_j}\right). \quad (8)$$

As we will see below in the Results section, for the AuNP at the lowest values of Q in the explored range, the decay of the time correlation function is fully explained by the translational diffusive term ($A_1 = 1$) alone, therefore by what we will call from now on the relaxation time $\tau_D = \tau_1$. In the opposite limit of higher Q (i.e., $QR_h \gg 1$, being R_h the hydrodynamic radius) the weight A_1 progressively vanishes and the decay is more and more explained by the internal polymer relaxation dynamics. In the intermediate Q range, the relaxation times τ_j in the second term of Eq. (8) are somehow affected by the translational diffusive contribution of the particle.

Since, as mentioned above, the term $I_{\text{int}}(Q, \infty)$ is related to the structure of the nanoparticle, the term A_1 accounts implicitly for the form factor of the PEG AuNP. In applying our RJ-MCMC approach (see next section), we will use for the AuNP the equation model (8) where the relaxation times τ_j for $j \geq 2$ will be explicitly referred to as $\tau_{\text{pol},j}$. For the case of the homopolymer solutions of PEG2000 and PEG400, we used a more general expression of the nISF that reflects the fact that we did not make any assumption on the possible character of the first relaxation and that was instead inferred only after the

data analysis:

$$\frac{I(Q, t)}{I(Q, 0)} \simeq \sum_{j=1}^k A_j \exp\left(-\frac{t}{\tau_j}\right). \quad (9)$$

It is worth noticing that the formulation of a model describing the relaxation dynamics in terms of simple exponential functions, as in Eqs. (8) and (9), is not the standard choice in polymer physics, where stretched exponentials are often assumed. In order to demonstrate how the intermediate scattering function for the considered systems can be fully described as a linear combination of simple exponentials, we generalized the model equations (8) and (9) to sums of Kohlrausch-Williams-Watts exponential functions with stretched coefficients β_j . By means of the Bayesian approach described in Sec. II C applied to this model, we will provide a quantitative estimate of the statistical significance of stretching parameters and we will show that the choice of nonstretched ($\beta_j = 1$) exponentials is fully adequate to an accurate description of the studied system dynamics and indeed the most valid on a probabilistic basis.

C. Bayesian analysis

Equation (9) describes the measured nISF relaxation as a linear combination of pure exponential components. As explained at the end of Sec. II B, we generalized this model by replacing the simple with stretched exponential functions. Let $t = (t_1, \dots, t_n)$ denote the vector of discrete times at which data were collected, and $y = (y_1, \dots, y_n)$ the corresponding vector of the nISF measurements. Generalizing Eq. (9), the data can be assumed to be sampled from the following expression:

$$y_i = \gamma \sum_{j=1}^k A_j \exp\left[-\left(\frac{t_i}{\tau_j}\right)^{\beta_j}\right] + \epsilon_i, \quad \text{for } i = 1, \dots, n \quad (10)$$

where γ is a proportionality constant to account for a possible adjustment of the data normalization and ϵ_i , for $i = 1, \dots, n$, are independent random noises, accounting for statistical errors in the measurements. We assume these random errors to be independent and identically distributed according to a normal distribution $\mathcal{N}(0, \nu\sigma_i^2)$, where σ_i is the measurement error corresponding to the i th observation and ν is a proportionality constant.

In order to choose an adequate number of components and estimate their parameters, we adopted a Bayesian approach. The main advantage of this approach in the present context is that the number of components, i.e., the relaxations in the time correlation function decay, can be treated as an unknown parameter to be estimated, along with the other model parameters. In addition, inferential results are obtained in the form of parameter probability density functions, so that the final choice on the number of relaxations is justified on a probabilistic basis. Furthermore, the choice of the Bayesian approach intrinsically and automatically applies the Ockam's razor principle [26], thus preventing overparametrization. Finally, the Bayesian approach has already proven its superiority with respect to classical regularization techniques in the solu-

tion of linear ill-posed problems, as well as to other numerical methods in the solution of nonlinear inverse problems [39].

Bayesian inference is based on Bayes theorem [40], which allows to combine evidence from data and possible prior knowledge I , through the equation

$$P(\Theta|y, I) \propto P(y|\Theta, I) \times P(\Theta|I), \quad (11)$$

where the symbol $|$ stands for “conditional on” and Θ is the whole vector of unknown parameters, i.e., $\Theta = (k, A, \tau, \beta, \gamma, \nu)$, with $A = (A_1, \dots, A_k)$, $\tau = (\tau_1, \dots, \tau_k)$, and $\beta = (\beta_1, \dots, \beta_k)$. We stress, again, that the number k of elementary components in the model is also treated as a parameter.

The first term on the right-hand side of Eq. (11) is the likelihood function, which is the probability of the observed data y , conditional on a certain set Θ of parameter values, and on prior knowledge. In the present case, the likelihood of the data is

$$P(y|\Theta) = \prod_{i=1}^n \phi(y_i; f(t_i), \nu\sigma_i^2), \quad (12)$$

where $f(t_i) = \gamma \sum_{j=1}^k A_j \exp[-(\frac{t_i}{\tau_j})^{\beta_j}]$ and $\phi(\cdot; f(t_i), \nu\sigma_i^2)$ is the density of the $\mathcal{N}(f(t_i), \nu\sigma_i^2)$ distribution, evaluated in a specific point “ $\cdot$ ”. The conditioning on I here and in the rest of the paper has been dropped to simplify the notation.

The second term on the right-hand side of Eq. (11) is the prior distribution for the parameters of the model function, before data collection. It incorporates all prior knowledge, if any, on the problem at hand. In a Bayesian perspective, such prior distributions need to be specified for each parameter of the model, including the unknown number of mixture components k . In order to allow for complete *a priori* ignorance on k , we assume $p(k) = 1/k_{\max}$, for $k = 1, \dots, k_{\max}$, where $k_{\max}$ is reasonably chosen to be the maximum possible number of components. The vector of weights A is assumed to follow *a priori* a Dirichlet distribution with hyperparameters $\lambda = (\lambda_1, \dots, \lambda_k)$, i.e., $A \sim \mathcal{D}(\lambda)$. Complete ignorance can be accommodated by letting $\lambda_1 = \dots = \lambda_k = 1$ so that each component of the mixture has a marginal weight which is uniformly distributed in the interval $(0, 1)$. Different prior beliefs on k and A can be easily accounted for [26]. The stretching parameters β_j are assumed *a priori* to independently follow a mixed distribution made up of a continuous beta density and a probability mass in 1, that is

$$\beta_j \sim \zeta \mathcal{B}(\kappa, \psi) + (1 - \zeta) \delta_{\beta_j, 1},$$

with κ and ψ being the hyperparameters of the beta density, and ζ denoting our prior support in favor of stretched (rather than unstretched) components. Notice that letting $\zeta = 0$ implies assuming a sum of pure exponential components, as the one in Eq. (9), whereas letting $\zeta = 1$ determines a sum of stretched components. For any other $0 < \zeta < 1$, the mixture can include both stretched and unstretched components. In particular, we assumed $\zeta = 0.5$, so that the probability of having a stretched or unstretched j th component is, *a priori*, the same, for all $j = 1, \dots, k$. In addition, we set $\kappa = 1$ and $\psi = 1$ so that the stretching parameters are assumed uniformly distributed on the interval $(0, 1)$. This corresponds to an uninformative prior giving the probability of 0.5 to an

unstretched component and the remaining probability of 0.5 to a stretched component, and, for a stretched component, assigning the same density to any value of β_j in $(0; 1)$. Therefore, our choice is that of being completely general: the same prior probability is given to any model made up of a sum of just stretched components or a sum of just unstretched components or a sum of stretched and unstretched components. Obviously, models with just one stretched or unstretched component are also included when $k = 1$. Our choice of prior satisfies both the requirements of the theory of correlation functions [22–24] and the possibility to describe the results in terms of theoretical models used in polymer physics [41] without pushing the data in either one or the other direction. Obviously, different priors can be considered. For example, if one wants to give a higher prior to the Zimm model, he could choose κ and ψ such that $\kappa/(\kappa + \psi) \approx 0.85$ (so that the *a priori* expected value of β_j would be 0.85), with ψ accounting for the degree of prior belief in the Zimm model: the smaller ψ , the higher the prior probability assigned to this model. If, instead, one wants to consider both Zimm and Rouse model as *a priori* more probable than other solutions, a mixture of two beta densities could be considered for β_j , with two peaks centered on 0.85 and 0.5, respectively, and mixture weights proportional to the prior probabilities assigned to the two models.

The relaxation times τ_j , for $j = 1, \dots, k$, are assumed to be *a priori* independent and identically distributed according to uniform densities in the interval $[0; \tau_{\max}]$, with $\tau_{\max}$ reasonably chosen. For the parameter γ , we also assume a mixed distribution made up of a continuous beta density and a probability mass in 1, that is

$$\gamma \sim \xi \mathcal{B}(\varphi, \rho) + (1 - \xi) \delta_{\gamma,1},$$

with φ and ρ being the hyperparameters of the beta density, and ξ denoting our prior support in favor of a refinement of the data normalization process. As for ζ , also in this case we let $\xi = 0.5$. Finally, for conjugacy reasons, ν is assumed to have an inverse gamma prior, $\nu^{-1} \sim \mathcal{G}(\iota, \varsigma)$, parametrized so that the mean and the variance are ι/ς and ι/ς^2 .

In the case of PEG AuNP data, where the first component of the model is meant to represent explicitly the center-of-mass diffusion described as in Eq. (7), we specify slightly different priors. In fact, the stretching parameter is fixed at the value 1 by using a degenerate prior distribution $P(\beta_1 = 1) = 1$. Moreover, for the translational diffusion parameter, we assumed $D \sim \mathcal{N}(\mu_D, \sigma_D^2)$, with the hyperparameters chosen in such a way that μ_D is equal to the value D_{LS} obtained by independent DLS measurements (see Sec. III A) and $\sigma_D = 0.20\mu_D$, so that D can vary by up to 40% from D_{LS} , with a probability of 95%. This was done to account for possible concentration effects on the translational diffusion, since the DLS experiment was done in a very dilute regime, whereas the NSE was done at a concentration of 5% in weight. Notice that the relaxation time of the first exponential component, conditional on D , is not random anymore and is simply obtained as $\tau_D = \frac{1}{DQ^2}$.

Bayesian inference is based on the joint posterior distribution of all the random parameters, i.e., the left-hand side in Eq. (11). Due to the complexity of the model at hand,

such posterior distribution is known up to a normalization constant, which would require solving a high-dimensional integral to be evaluated [42,43] [i.e., the integral of the right-hand side of Eq. (11) over the parameters Θ , which is required in the denominator of the right-hand side to have the joint posterior density integrate to one and the proportionality symbol replaced by equality]. For this reason, we resort to MCMC algorithms to simulate the joint posterior distribution of interest.

MCMC algorithms allow to draw samples from a target distribution (the joint posterior in Bayesian applications), known up to a normalizing constant. In this regard, these algorithms make it feasible to simulate the joint posterior distribution for the parameters of arbitrarily complicated models. MCMC algorithms are based on the construction of an irreducible and aperiodic Markov chain that has the target distribution as invariant distribution [44]. However, in our model the dimension of the parameter space is random since it depends on the random value of k . Therefore, the MCMC algorithm needs to be incremented with a reversible jump (RJ) step [25], which allows to jump between models having different dimensions.

The RJ-MCMC algorithm we propose performs M sweeps and, at each sweep, all the parameters are updated in turn, sampling the new value of a certain parameter conditionally on the data and all the remaining parameters and is able to explore models with any number of components as well as to move between models with stretched or unstretched components or any mixing of the two, thus including also Zimm and Rouse models. At each sweep, we perform a changing dimension move to update the value of k (through RJ sampling), as well as fixed dimension moves to update all the other parameters (either through Gibbs or Metropolis-Hastings sampling [45]). All the details about the algorithm and the updating moves are exhaustively explained in a dedicated paragraph of the Supplemental Material [37].

After an initial burn-in period, during which the RJ-MCMC draws are discarded until the convergence of the chain to the target distribution is reached, the algorithm produces at each sweep m a new realization of all the parameters of the model, sampled from their joint posterior distribution. Let $(k^{(m)}, A^{(m)}, \tau^{(m)}, \beta^{(m)}, \gamma^{(m)}, \nu^{(m)})$, for $m = 1, \dots, M$, be the sample obtained after M sweeps of the algorithm. This sample provides the simulated joint posterior distribution of all parameters and can be used to estimate all the quantities of interest.

The posterior distribution of the number of exponential components can be estimated as the proportion of times each model was visited by the algorithm, i.e.,

$$\hat{p}(k = \ell | y) = \sum_{m=1}^M \delta_{k^{(m)}, \ell} / M = M_\ell / M,$$

where M_ℓ is the number of times the model with ℓ components was visited. In practice, the model chosen will be the one with the highest posterior probability.

Once a particular model has been chosen, we can estimate the parameters of each exponential component under that model. The estimates of these parameters can be computed as the mean of their simulated marginal posterior distribution,

conditional on the model with ℓ components. Point estimates for the weights and relaxation times of the ℓ exponential components are given, respectively, by

$$\begin{aligned}\hat{A}_j &= \sum_{m:k^{(m)}=\ell} A_j^{(m)} / M_\ell, \\ \hat{\tau}_j &= \sum_{m:k^{(m)}=\ell} \tau_j^{(m)} / M_\ell, \\ &\text{for } j = 1, \dots, \ell\end{aligned}$$

where the sum over $m : k^{(m)} = \ell$ means over all values of m in which the model with $k = \ell$ has been visited.

Note that in case the simulated marginal posterior distribution of a certain parameter is not symmetric, other estimators than the arithmetic mean can be considered (e.g., the simulated posterior mode). For the stretching parameters, if the simulated posterior probability of $\beta_j = 1$ is larger than its prior, i.e., $\sum_{m:k^{(m)}=\ell} \delta_{\beta_j^{(m)},1} / M_\ell \geq 1 - \zeta$, then we simply let $\hat{\beta}_j = 1$, otherwise $\hat{\beta}_j = \sum_{m \in \mathcal{S}} \beta_j^{(m)} / |\mathcal{S}|$, where $\mathcal{S}$ is the set of sweeps for which $k^{(m)} = \ell$ and $\beta_j^{(m)} < 1$, and $|\mathcal{S}|$ is the dimension of $\mathcal{S}$. All the other model parameters do not depend on the number of exponential components and can be estimated from their simulated marginal posterior distributions, averaging over all the sweeps of the algorithm.

III. EXPERIMENTAL RESULTS AND ANALYSIS

A. PEG2000 AuNP and PEG400 AuNP

A representative set of NSE data obtained for the PEG2000 AuNP is shown in Fig. 2(a). The full set is reported in the Supplemental Material (Fig. S2) [37]. As expected, the nISF decays as a function of the NSE time and the decay becomes faster as Q increases.

In the investigated Q range we probed length scales ranging from 30 Å to about 300 Å, which include distances on the same scale of the intrapolymer and interpolymer chain interactions, as well as distances on the scale typical of translational diffusion of the whole nanoparticle. The translational diffusion coefficient was estimated by DLS measurements carried out on very dilute solution of the NPs in D₂O (~ 0.02 wt.%). The values obtained for this parameter were $D_{LS} = 1.35$ and $3.5 \text{ Å}^2/\text{ns}$ at temperatures 280 and 318 K, respectively. From the D_{LS} we can estimate the hydrodynamic radius of the PEG2000 AuNP, $R_h \sim 85 \text{ Å}$ through the relation $D = k_B T / 6\pi\eta R_h$, where k_B is the Boltzmann constant, T the temperature, and η the viscosity of D₂O at the temperature T .

At low Q , where the condition $QR_h \sim 1$ holds, the dynamics should be dominated by the translational diffusion of the AuNP, whereas at higher Q where $QR_h \gg 1$ the dynamics should be dominated by the polymer relaxations. At intermediate Q values, a Q -dependent combination of these two contributions is expected. With this in mind and in analogy with Arriaga *et al.* [46] we initially described the decay of $I(Q, t)/I(Q, 0)$ as

$$I(Q, t)/I(Q, 0) = A(Q)F_{NP}(Q, t) + [1 - A(Q)]F_{pol}(Q, t), \quad (13)$$

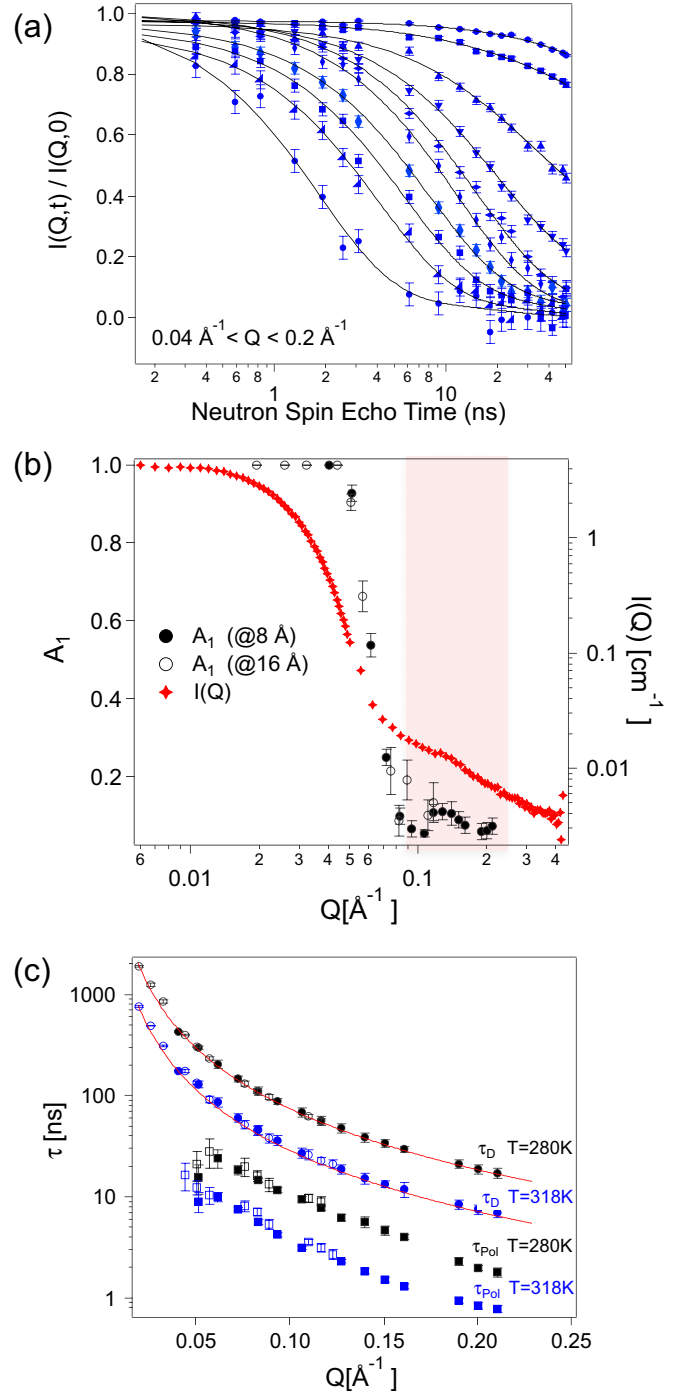


FIG. 2. (a) A representative series of NSE curves representing $I(Q, t)/I(Q, 0)$ of the PEG2000 AuNPs dispersed in D₂O at neutron wavelength $\lambda = 10 \text{ Å}$ and a temperature of 280 K (symbols). The Q values increase from above to bottom. The lines represent the best fits to the data. (b) Q dependence of the A_1 parameter obtained by the RJ-MCMC analysis and of the $I(Q)$ measured for a solution of PEG2000 AuNPs in D₂O at the 1 wt.% concentration; the red diamond symbols correspond to a SANS measurement performed on the sample. (c) The two relaxation times τ_D and τ_{pol} obtained by the fitting analysis at different temperatures. The full and empty symbols correspond to the values obtained at 10 and 16 Å. The red solid lines indicate the relaxation time obtained from the translational diffusion coefficient as derived by DLS at the respective temperatures.

where

$$F_{\text{NP}}(Q, t) = \exp(-t/\tau_D) \quad (14)$$

is the translational diffusion of the NP, and

$$F_{\text{pol}}(Q, t) = \exp[(-t/\tau_{\text{pol}})^\beta] \quad (15)$$

is the term approximately associated to the polymer dynamics (see above) often described with a Kohlrausch-Williams-Watts function. In the previous equations, A is the Q -dependent relative portion of translational diffusion with respect to the polymer dynamics contribution, $\tau_D = 1/DQ^2$ is the characteristic time of the translational diffusion of the nanoparticle in D_2O , τ_{pol} is a characteristic time associated to the polymer dynamics [although it contains also the translational diffusive dynamics, cf. Eq. (7)], and β the stretching exponent, which often accounts for the possibility of having a distribution of relaxation processes. However, we wanted to leave open the possibility to describe the polymer relaxations with a more general model. Consequently, we chose to describe the polymer relaxations with sums of exponential (stretched or simple) functions in order to find the combination that described the data with the greatest statistical significance.

In Eq. (8) we imposed the first component to be a simple exponential [which corresponds, as discussed in Sec. II C, to set $P(\beta_1 = 1) = 1$ as degenerate prior distribution for β_1] since it is associated to the translational diffusion of the NP [see, for example, Eq. (7)]. As already discussed in Sec. II C, we assumed $D \sim \mathcal{N}(\mu_D, \sigma_D^2)$. We allowed the remaining components to either be simple or stretched exponentials.

The results of the fitting procedures are listed in Tables SI, SII, SIII, and SIV in the Supplemental Material [37]. The column labeled $P(k|y)$ lists the probability associated to the most statistically supported model conditional to our data. In Fig. 2 we show a summary of the main results concerning the fitting analysis for the PEG2000 AuNP. For $Q < 0.04 \text{ \AA}^{-1}$, the RJ-MCMC algorithm privileged a solution with only the translational diffusive component at both measured temperatures [Figs. 2(b) and 2(c)].

Above this value of Q , the solution with two relaxations was always chosen with a probability $P(k = 2|y)$ usually between 70% and 90% especially at high Q and at the higher temperature. The mean β_2 stretching coefficient was estimated by the algorithm to be close to 1, and in any case the probability to have $\beta_2 = 1$ was most of the times greater than 90%.

In Fig. 2(b), the trend in Q of the parameter $A_1(Q)$ is displayed at the temperature of 280 K. At $Q < 0.04 \text{ \AA}^{-1}$ the

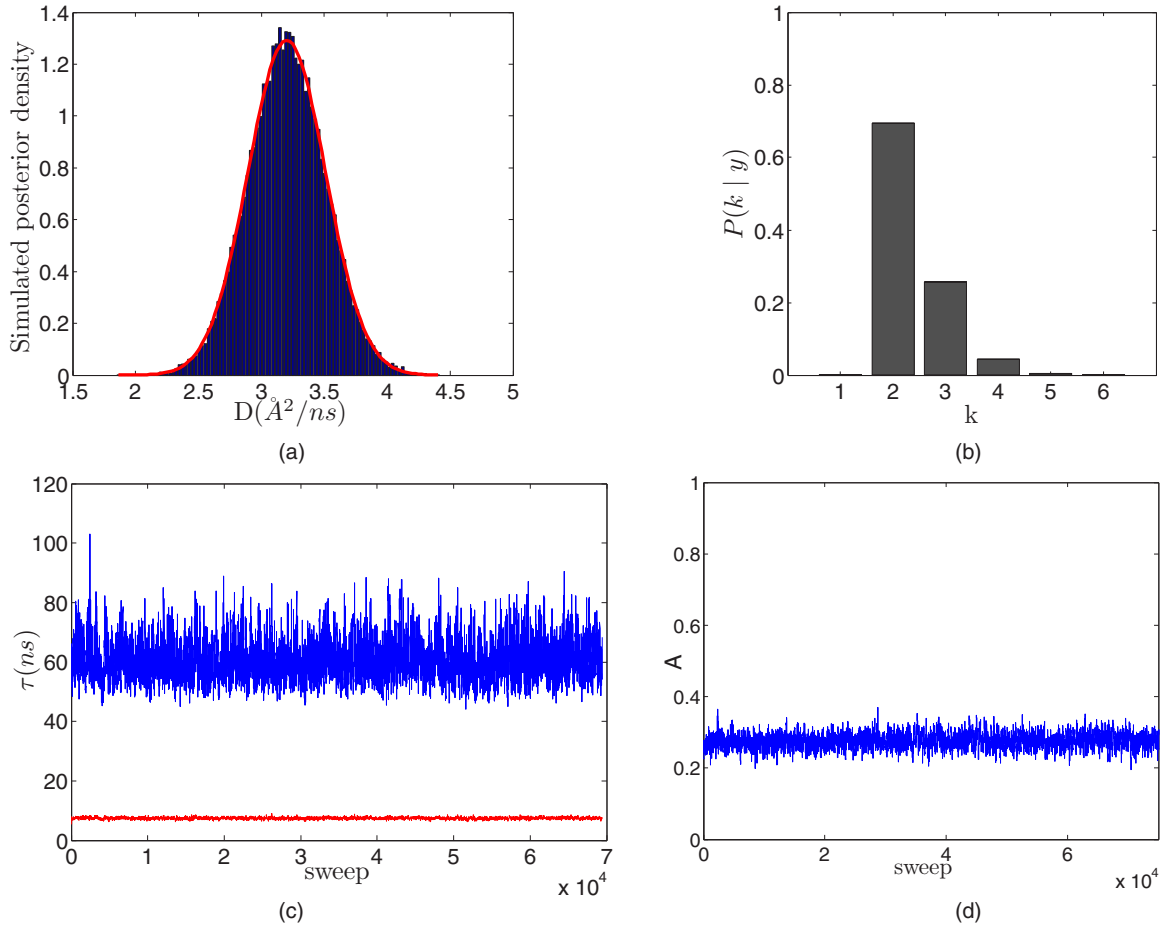


FIG. 3. (a) An example of the posterior distribution for the parameter D for the PEG2000 AuNPs at 318 K, $\lambda = 10 \text{ \AA}$, $Q = 0.72 \text{ \AA}^{-1}$. The solid red curve is a Gaussian fit to the simulated posterior distribution. (b) Posterior distribution for the number of relaxation components for the same example in (a). (c) Trace plot for the relaxation time τ_D (bottom curve) and τ_{pol} (upper curve). (d) Trace plot for the A_1 parameter.

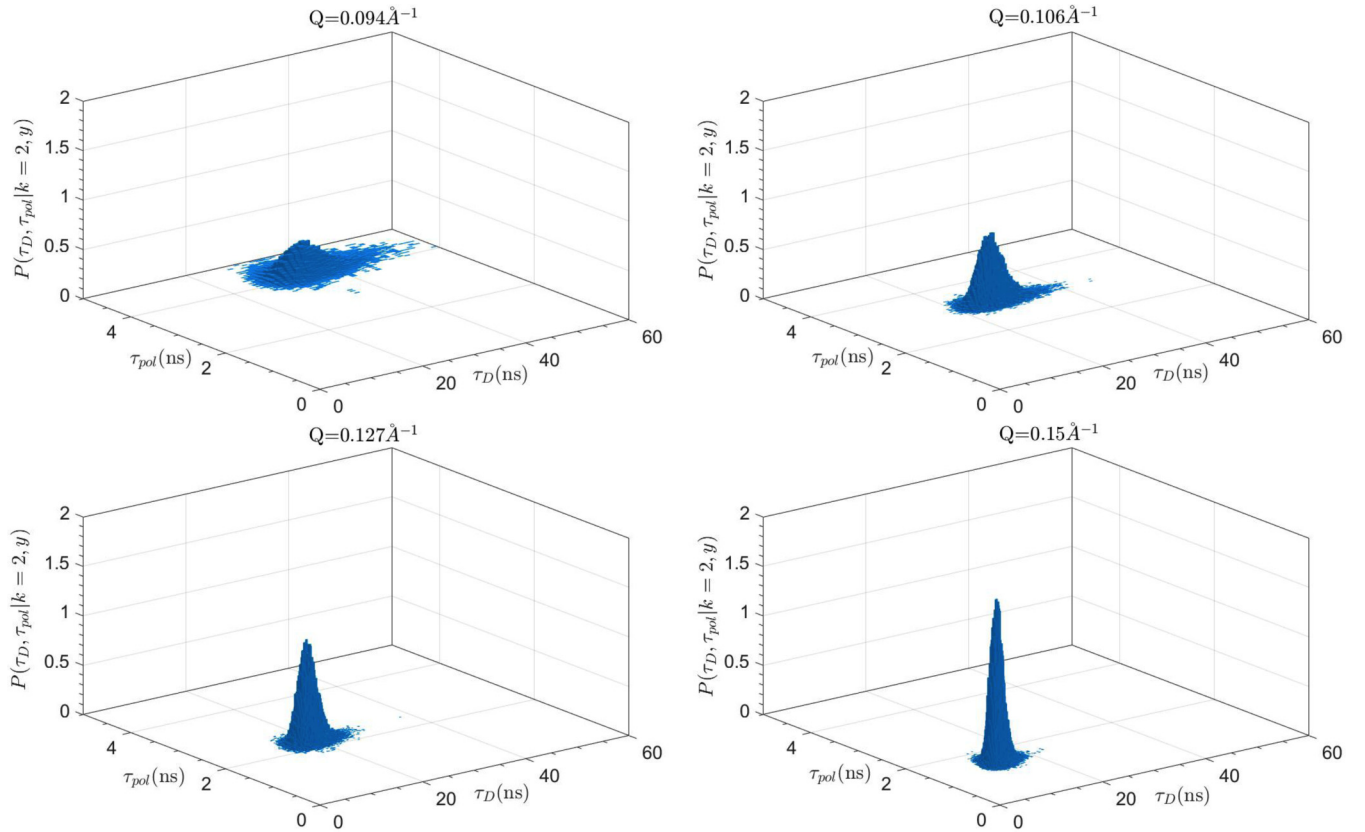


FIG. 4. Joint posterior distribution at four selected Q values, of the two relaxation times, conditional on $k = 2$ and marginalized with respect to the other model parameters for the PEG2000 AuNP sample at $T = 318$ K and $\lambda = 10$ Å.

main contribution to the dynamics came from the translational diffusion of the NP. At $Q > 0.1$ Å⁻¹ the dominant contribution was the polymer dynamics. In the intermediate Q regime 0.04 Å⁻¹ $< Q < 0.1$ Å⁻¹, we observed a transition between these two regimes.

The value of the translational diffusion coefficient estimated by the algorithm was in perfect agreement with the one measured by DLS [see Fig. 2(c), solid red line] which confirms that the contribution of the rotational term can actually be neglected. From Fig. 2(c) it is also clear that the relaxation times estimated at the two wavelengths are consistent with each other. As expected, the relaxation times decreased both with increasing the temperature and Q .

In Fig. 3 we report as example a typical set of results coming from the fitting algorithm obtained at $\lambda = 10$ Å and at $Q = 0.72$ Å⁻¹. Figure 3(a) shows an example of the posterior distribution of the translational diffusion coefficient D at $Q = 0.72$ Å⁻¹. In this specific case, the simulated posterior distribution was a Gaussian, as shown by the fit (red curve). The spread of this distribution is related to the polydispersity of the NP and to the errors bars of the experimental $I(Q, t)$. Figure 3(b) displays the posterior distribution of the number of components in Eq. (10) probed by the algorithm. In Figs. 3(c) and 3(d), the trace plots for the two relaxation times τ_D and τ_{pol} and for the A parameter are shown, always for the same λ and Q . In Fig. 4 we report as an example the joint posterior distribution $P(\tau_D, \tau_{pol} | k = 2, y)$ for the two relaxation times marginalized with respect to all the other model parameters as

a function of Q for the PEG2000 AuNP sample at 318 K. It can be noticed that at small Q values, the posterior distribution of τ_D is more dispersed than that of τ_{pol} , implying a minor accuracy in the estimation of the first parameter, compared to the second one. As Q increases, the difference between the dispersions of the two parameters vanishes and, in addition, the joint distribution becomes more concentrated, implying that more precise estimates of both parameters can be obtained at large Q values. Also, notice the absence of correlation between the two parameters.

A representative set of the NSE data measured on the PEG400 AuNPs dispersed in D₂O are shown in Fig. 5(a). At both the investigated temperatures, the most probable model predicted by the algorithm to describe the decay of $I(Q, t)/I(Q, 0)$ was a single exponential relaxation. At 280 K, even if the quality of data at $Q \geq 0.8$ Å⁻¹ degraded (see Fig. S4 in the Supplemental Material [37]), the probability assigned by the algorithm to a one-component model was, except for a few data sets, significantly greater than 60%. The estimated relaxation time at each Q value is very close to the one expected for the NP translational diffusion as measured by DLS [$D_{LS} = 2.2$ Å²/ns; cf. Fig. 5(b), red solid lines]. At 318 K the situation was clearer in view of an improved data quality. This allowed an unambiguous fitting of the relaxation curves given that the algorithm assigned to a one-component relaxation model a posterior probability $P(k = 1|y)$ almost always above 90%. The translational diffusion coefficient from the fits is in agreement within the errors with that

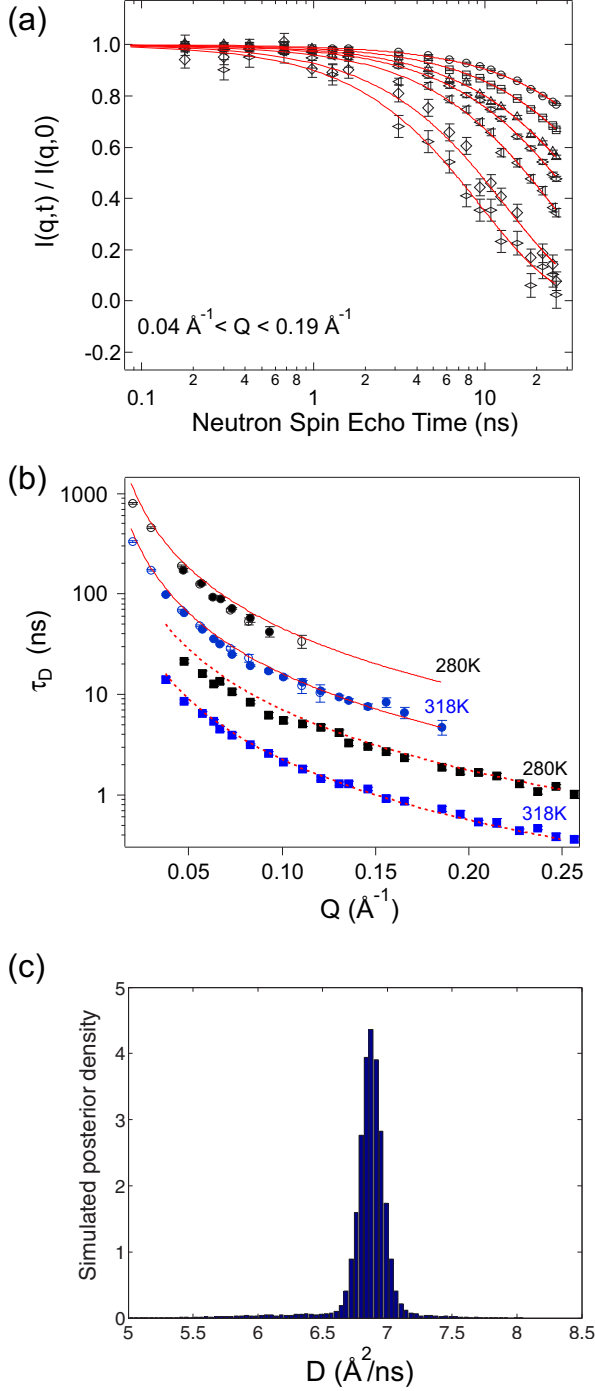


FIG. 5. (a) Representative set of NSE curves measured on the PEG400 AuNPs at $\lambda = 8$ Å and $T = 318$ K; the lines represent the best fits to the data. (b) Diffusion relaxation time as derived by the RJ-MCMC algorithm for the PEG400 AuNPs (circles) and for the PEG400 polymers in 5% solution in D₂O (squares). The empty and solid symbols refer to $\lambda = 8$ and 16 Å, respectively. The continuous red lines correspond to the relaxation times obtained from the translational diffusion coefficients as measured by DLS measurements for the PEG400 AuNPs. The dashed lines indicate the relaxation times obtained from the translational diffusion coefficient of PEG400 solution as obtained from literature [47]. (c) An example of the simulated posterior distribution for the diffusion coefficient obtained at $T = 318$ K and $Q = 0.048$ Å⁻¹.

obtained with the DLS ($D_{LS} = 6.2$ Å²/ns) and corresponds to a hydrodynamic radius of ~ 48 Å. The complete set of parameters obtained by the fits are listed in Tables SV, SVI, SVII, and SVIII in the Supplemental Material [37].

B. Polymer solutions of PEG2000 and PEG400

1. PEG2000

In order to better characterize the dynamics of the polymer grafted to the NP, we compared the results obtained with the PEG-coated NPs and those obtained with a reference system identified as a solution of PEG homopolymers. We used PEG2000 and PEG400 that possess a molecular weight very similar to those tethered to the NPs. As previously shown [13], the concentration of polymer as a function of the distance from the gold surface is not uniform, but varies being rather high in the vicinity of the surface of the NP and lower away from it. Ideally, we should compare the results obtained on the PEG-coated NP with several PEG solutions at different concentrations. The PEG2000 was studied in D₂O

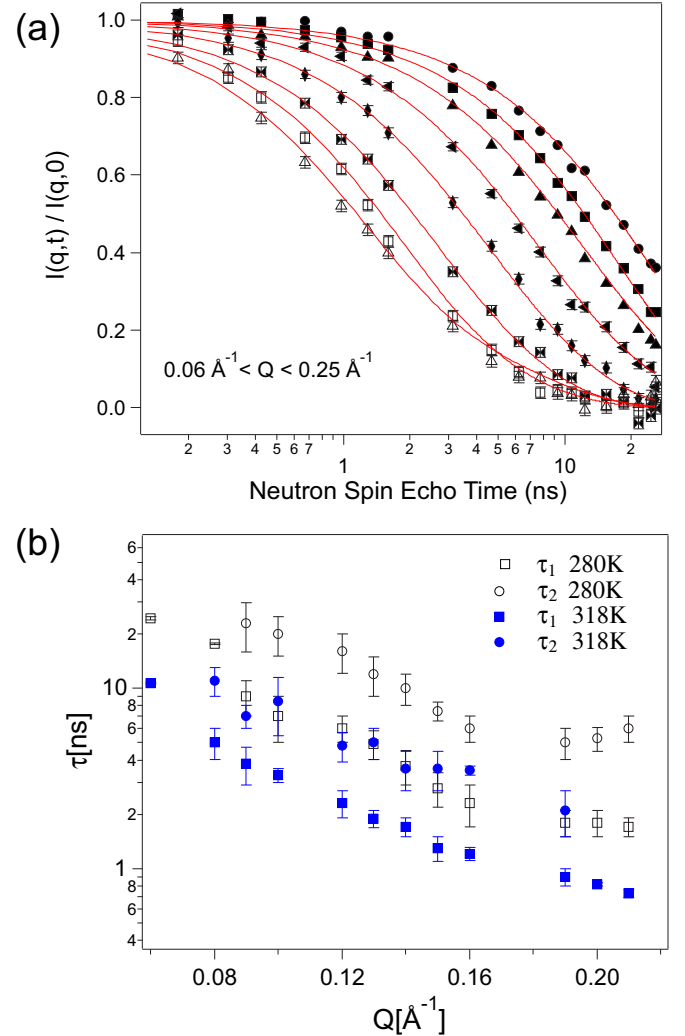


FIG. 6. (a) Representative set of NSE curves measured on the PEG2000 at $\lambda = 8$ Å and $T = 280$ K. The lines represent the best fits to the data. (b) The two relaxation times τ_1 and τ_2 obtained by the fitting analysis at different temperatures.

at weight concentrations of 10%, 20%, and 45% (the latter is the solubility limit for PEG2000 in D₂O) as a reference. However, due to an unfavorable combination of coherent and incoherent components that contribute to the neutron polarization with opposite sign, the measurements at the higher concentrations (20% and 45% in weight) suffered of a too weak spin echo signal and were not taken into account. The choice of the reference system was based on a compromise between the counting statistics and the available measuring time. Ultimately, the PEG2000 and PEG400 were studied at concentrations of 10% and 5% in weight, respectively. A representative set of data obtained for the PEG2000 at 10% is displayed in Fig. 6(a).

The protocol used to analyze the NSE data of the polymers was slightly different from the one used for the NPs, as we did not use any informative prior for the first relaxation. In fact, we left the algorithm free to search for the most statistically probable solution. The prior distribution of the relaxation times was chosen as to be uniform in an interval $[0; \tau_{\max}]$.

In a first attempt, we used a combination of simple exponential functions to describe $I(Q, t)/I(Q, 0)$. The results for PEG2000 homopolymers at a concentration of 10% and at $\lambda = 8 \text{ \AA}$ are shown in Fig. 6(b). The parameters obtained from the fits are listed in Tables SIX and SX in the Supplemental Material [37]. At $T = 280 \text{ K}$, the algorithm favors a single exponential for $Q \leq 0.08 \text{ \AA}^{-1}$ and two exponentials for $Q > 0.08 \text{ \AA}^{-1}$ [see Figs. 6(b) and 7, top panels]. At $T = 318 \text{ K}$ we have a similar behavior up to $Q = 0.19 \text{ \AA}^{-1}$.

Above this value, the algorithm would still favor the solution with two components. However, while the parameters of the first component are nicely estimated to continue the trend of their respective values at lower Q 's, the parameter τ_2 results scarcely identified. In fact, the simulated posterior distribution of τ_2 turns out to be substantially uniform in a large range of time [Fig. S5(b)]. As a consequence, any value of τ_2 sufficiently large would define a second exponential function able to fit the decay of the time correlation function at large spin echo time. We, therefore, decided to retain both components in the model [a single component would simply misconside the last few points; see Fig. S5(a) in the Supplemental Material [37]] but we do not trust estimates for the slower dynamics, since meaningless. The algorithm still provides the best probabilistically supported solution given the experimental data, but at the highest Q values the number of data points available for a precise determination of τ_2 is insufficient. The application of our method particularly for data sets as the ones at $Q > 0.19$ illustrates the efficacy of the Bayesian analysis since we can infer the reliability of the parameter estimates from the simulated marginal posterior distributions of the different parameters. This represents effectively a diagnosis tool to exploit at maximum the collected experimental data.

We stress that the posterior probability for the chosen values of k is always larger or much larger than 60% (see an example in Fig. 8). Also, the estimate of the relaxation times is accurate for the majority of data sets and providing

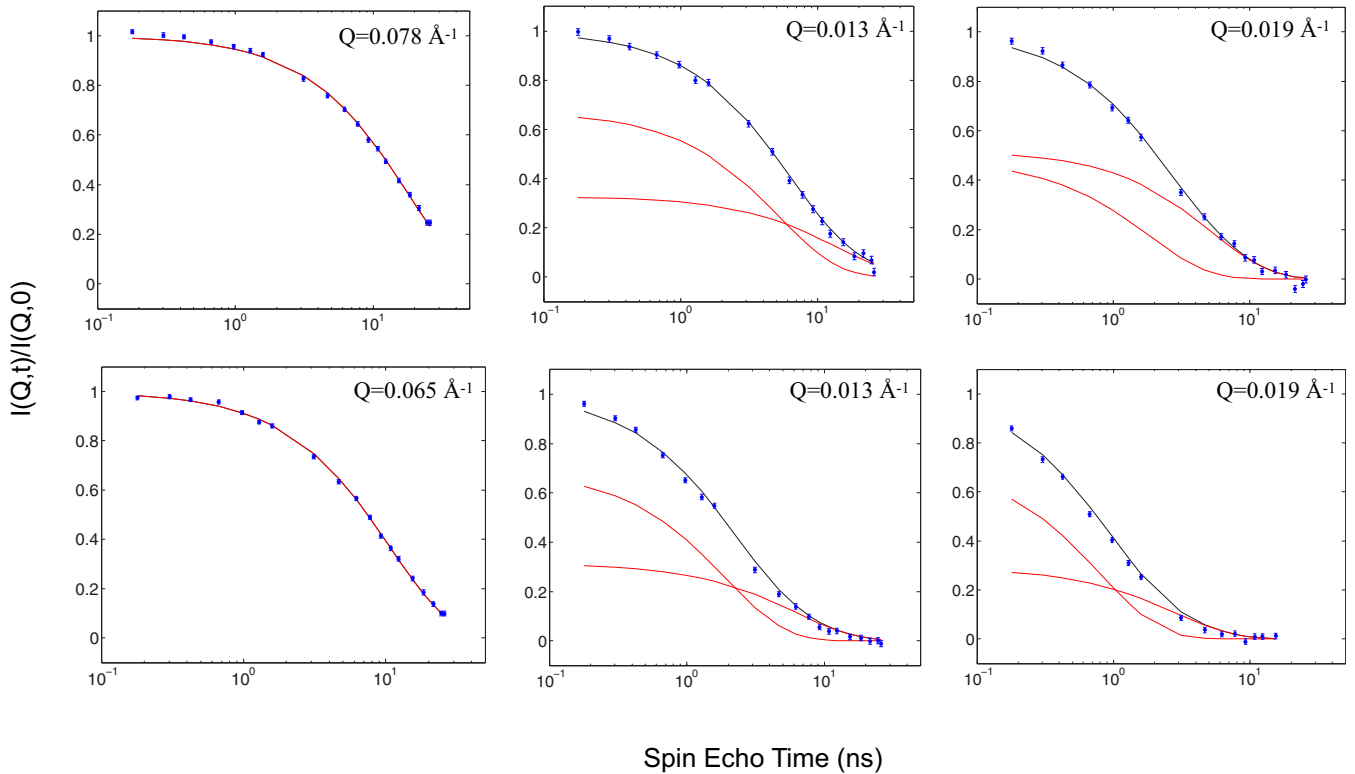


FIG. 7. NSE curves representing $I(Q, t)/I(Q, 0)$ measured on the PEG2000 homopolymer 10% concentration, at 280 K (top panels) and 318 K (bottom panels) at three selected Q values. For single relaxation ($Q = 0.078 \text{ \AA}^{-1}$, $Q = 0.065 \text{ \AA}^{-1}$) the lines are the best fit to the data. For multiple relaxations ($Q = 0.013 \text{ \AA}^{-1}$, $Q = 0.019 \text{ \AA}^{-1}$) the black lines fully overlapping to the data points are the best fit. We show also the relaxation components (red curves).

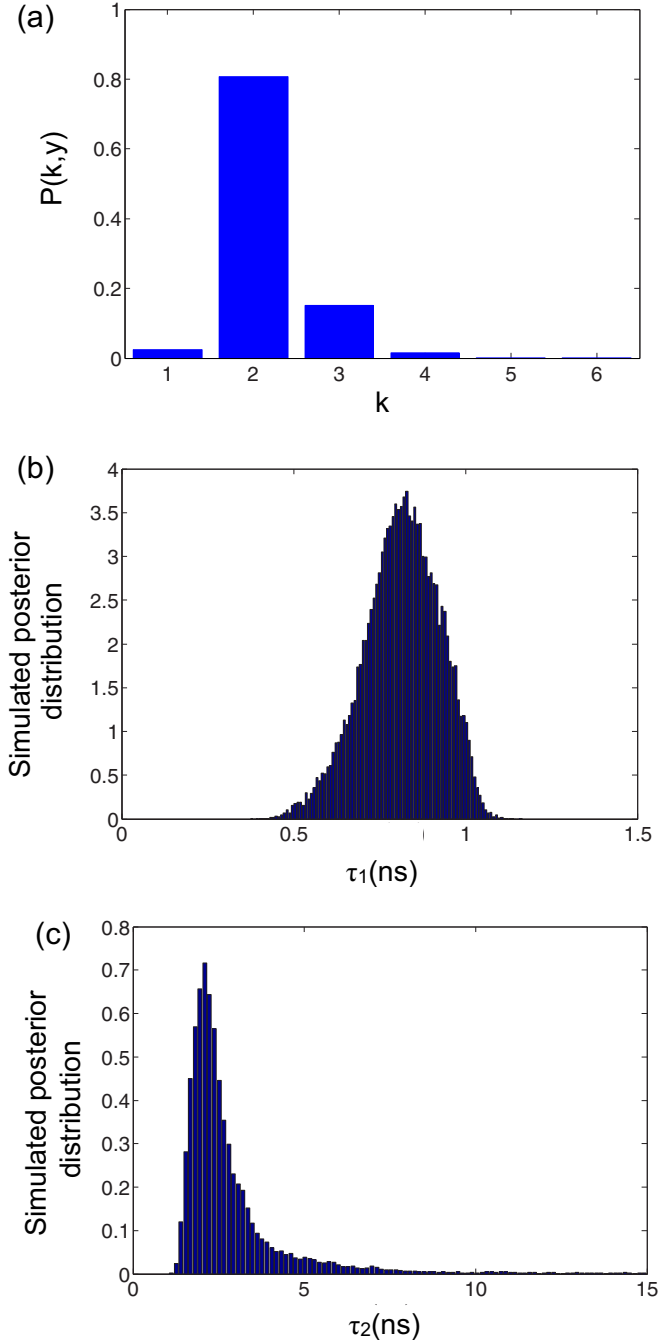


FIG. 8. We provide as example the posterior distribution obtained for PEG2000 homopolymer at 280 K and at $Q = 0.19 \text{ \AA}^{-1}$ for (a) the number k of relaxation components (b) τ_1 and (c) τ_2 .

for the posterior parameter distribution function a Gaussian distribution or at least a well-shaped unimodal distribution although not necessarily symmetric [see Figs. 8(b) and 8(c)].

Further measurements on the PEG2000 homopolymers dissolved in D_2O at 10% in weight were carried out on the spin echo spectrometers IN11 at ILL with an incident wavelength of $8.5 \text{ \AA}$. These measurements expanded the accessible time and momentum transfer, covering a time window between 0.02 and 5 ns and a Q range between 0.09 and $0.42 \text{ \AA}^{-1}$. In order to further validate the scenario so far achieved, we

joined the data collected with the two spin echo instruments IN11 and IN15 extending the Bayesian MCMC data analysis to this enlarged time window. The available data from IN11 were only at 280 K. In Fig. 9 we show the relaxation curves at four selected momentum transfer wave vectors ($Q = 0.12, 0.14, 0.20$, and $0.23 \text{ \AA}^{-1}$) at which we collected the data from both spectrometers.

At $Q = 0.12, 0.20$, and $0.23 \text{ \AA}^{-1}$ we reobtain substantially the same results we found from the analysis of the IN15 data only, both for the characteristic relaxation times and the weights A_1 and $(1 - A_1)$. At $Q = 0.14 \text{ \AA}^{-1}$ there is a slight difference for the faster relaxation, but the results are fully consistent within the uncertainties (Table II).

2. PEG400

The study of the PEG400 homopolymer solution has been undertaken at a concentration of 5 wt.%. The description once again is rather simple. As in the case of the PEG400 AuNP's, one relaxation is sufficient to give account of the time correlation decays at both the investigated temperatures. The posterior distribution function for k predicts a probability $P(k = 1 | y)$ for a single relaxation model always in the range 80%–98% for all the spectra analyzed, with few exceptions where anyway such probability was larger than 60%. The relaxation times obtained at the two temperatures are summarized in Fig. 5(b). These relaxation times are coherent with the values of translational diffusion of PEG400 at the same temperature and concentration [cf. Fig. 5(b), dotted red lines] [47]. A selection of the experimental NSE curves at the two investigated temperatures and at three selected Q values and fits are displayed in Fig. S6 of the Supplemental Material [37]. In Fig. S7 we report two examples of the simulated posterior distribution function for the relaxation time τ_1 for two different values of Q at $T = 318 \text{ K}$. In these cases, the posterior distribution functions look unimodal and almost symmetric. The relaxation times obtained for the PEG400 indicate a faster dynamics compared to that found in the PEG400 AuNP [Fig. 5(b)].

IV. DISCUSSION

The fitting analysis of the PEG2000 AuNP data details two relaxation processes that occur in the time window probed by the NSE experiment. The slower relaxation process with characteristic time τ_D corresponds to the translational diffusion of the PEG2000 AuNP. As it appears from Fig. 2(c), it perfectly overlaps with the relaxation times obtained from the expression $\tau_D = 1/DQ^2$ where D is the translational diffusion as derived by DLS measurement of PEG2000 AuNP at infinite dilution and at the respective temperatures. To further check if the interaction of the NP due to the finite concentration used in the experiments had a significant effect on the dynamics, we obtained the structure factor $S(Q)$ by performing SANS experiments at increasing concentrations (see Fig. S3 in the Supplemental Material [37]). Following the de Gennes narrowing, a significant interaction would produce a modulation of the translational diffusion according to the inverse of the structure factor $1/S(Q)$. As it is shown in Fig. S3, this modulation is barely visible within the

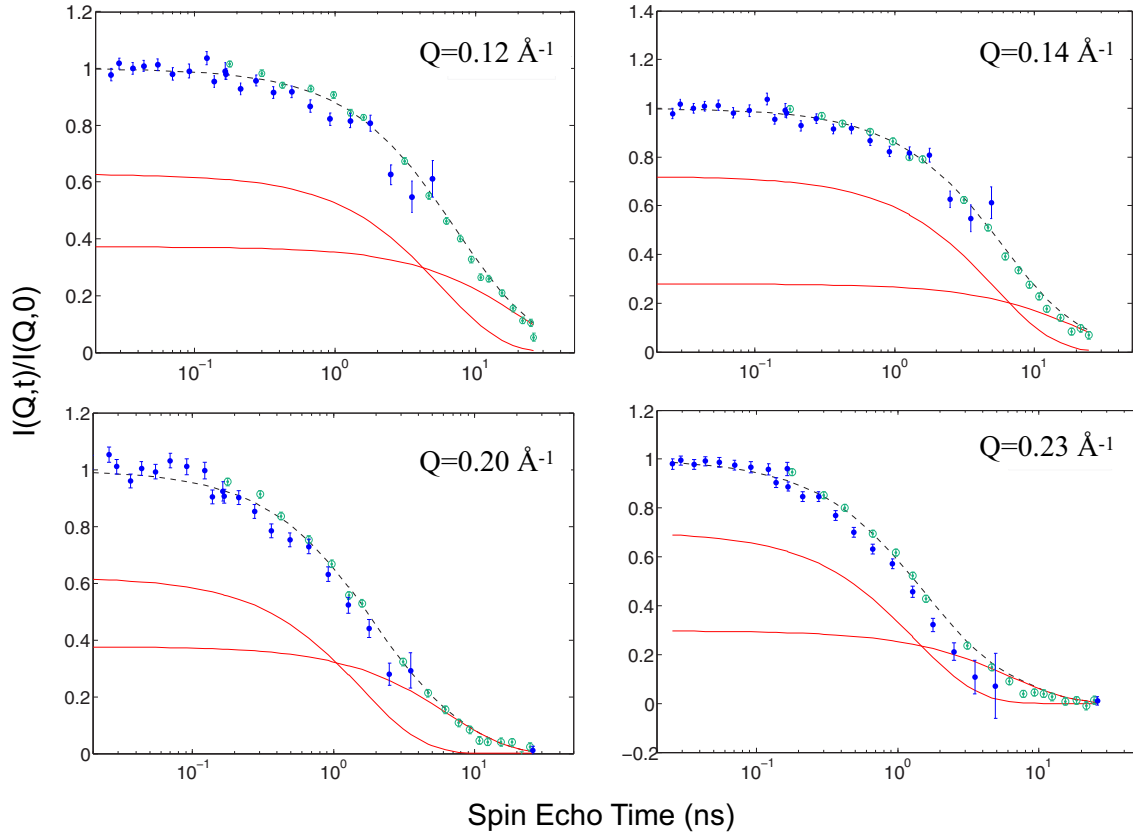


FIG. 9. NSE curves for the data collected with IN11 (blue filled symbols) and IN15 (green empty symbols) joined together on a PEG2000 homopolymer solution in D_2O at 10% concentration, at four selected Q values and $T = 280$ K.

error bars. This makes the assumption of negligible interactions between the nanoparticles reasonably applicable in this case.

While at small Q the signal is dominated by the coherent contribution due to the translational diffusion of the PEG AuNP, at high Q the signal is dominated by the chain relaxation ($A_1 \leq 0.2$) and it is composed by similar proportion of coherent and incoherent (Fig. S1). The observation of only one single exponential decay at those high Q values implies that the coherent and incoherent scatterings have a similar relaxation behavior within the limit of the sensitivity of the

technique. A possible explanation for this feature is that we are looking to a system of polymers closely tethered to a small surface from the perspective of a relatively large length scales (>25 Å). In these conditions it is reasonable to think that the polymers cannot move independently from each other and that the self-motion and self-coordinated motion of the polymers relax on similar timescales.

The second relaxation process is related to the dynamics of the polymer chains tethered to the NP, as shown by the fact the τ_{pol} resembled the relaxation times found in solution of PEG2000 in D_2O (Fig. 10). The relative contribution of the two relaxation processes is expressed by the Q -dependent parameter A_1 , that represents the portion of translational diffusion, with the sum rule $A_1 + A_2 = 1$. As shown in Fig. 2(b), the translational diffusion of the NP dominates at low Q and the polymer dynamics dominates at high Q . Interestingly, the trend of A_1 as a function of Q suggests a maximum at the same Q where the $I(Q)$ obtained by SANS measurements presents a shoulder [see the shaded region in Fig. 2(b)]. In fact, the form of the function $A_1(Q)$ reflects a modulation of the structure factor. At the Q values where the maximum occurs, the incoherent part of the scattering is around 30%. This incoherent part also contributes to the $A_1(Q)$ with a term given by $I_{inc}(Q, \infty)/I_{inc}(Q, 0)$. This term is the usually called EISF (elastic incoherent structure factor), which relates to the volume where the hydrogen atoms are confined in, i.e., the nanoparticle corona.

TABLE II. Comparison between the values of the relaxation times and their weights obtained by fitting only the IN15 data and by fitting the nISF obtained by gathering together the data from IN11 and IN15.

Q (Å ⁻¹)	Instrument	τ_1 (ns)	τ_2 (ns)	A_1
0.12	IN15	6 ± 1	16 ± 4	0.59
	IN15 and IN11	6 ± 1	12 ± 4	0.60
0.14	IN15	3.7 ± 0.8	10 ± 2	0.46
	IN15 and IN11	4.7 ± 0.8	10 ± 4	0.60
0.20	IN15	1.8 ± 0.3	5.3 ± 0.8	0.49
	IN15 and IN11	1.6 ± 0.2	5.2 ± 0.8	0.62
0.23	IN15	1.6 ± 0.2	5 ± 1	0.69
	IN15 and IN11	1.3 ± 0.2	4.0 ± 0.8	0.69

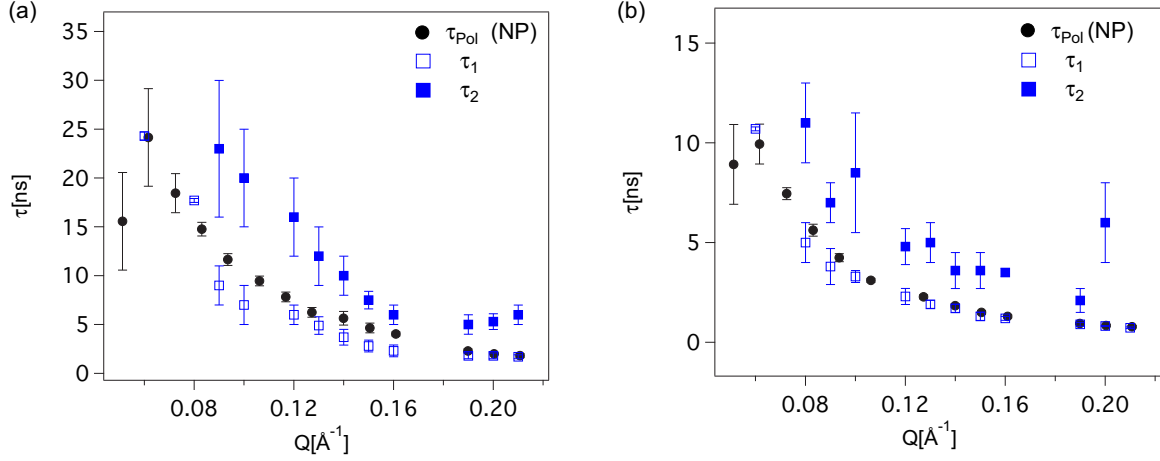


FIG. 10. Characteristic relaxation times τ_1 and τ_2 for the homopolymer PEG2000 solution at 10% concentration and relaxation time of the polymer in the AuNP, τ_{pol} , at (a) $T = 280$ K and (b) $T = 318$ K.

It is useful to plot the inverse of the relaxation times in terms of $\Gamma = 1/\tau$. As detailed in Eq. (7), the $\Gamma_{\text{int}} = (DQ^2 + \frac{1}{\tau_{\text{int},j}})$ contains also the term DQ^2 related to the translational diffusion of the whole NP. In Figs. 11(a) and 11(b) we report as a function of Q^2 the $\Gamma_D = 1/\tau_D$ and the $\Gamma_{\text{pol,cor}}$, where the latter corresponds to the internal dynamics corrected from the translational diffusion of the whole NP. Both Γ_D and $\Gamma_{\text{pol,cor}}$ have a clear linear trend that intercepts the origin. A dynamic whose relaxation can be described by a simple exponential decay in time, with a relaxation rate $\Gamma \sim Q^2$, has a diffusionlike behavior. The coefficient diffusion associated to the relaxation time of the tethered polymer chains is around $9 \pm 2 \text{ \AA}^2 \text{ ns}^{-1}$ and $25 \pm 3 \text{ \AA}^2 \text{ ns}^{-1}$ at temperatures of 280 and 318 K, respectively. Such a behavior is also associated to a linear trend with t of the mean square displacement. This is in qualitative agreement with that observed in coarse grained molecular dynamics simulations [48] that described the dynamics of brush polymers on spherical nanoparticles. This is an important result since the fact that the dynamics of a corona is represented by a single exponential in this timescale and

length scale is not due to intrinsic limitation of the technique or of the analysis that is not able to distinguish a higher number of relaxations but it corroborates computational studies. The linear trend for the mean square displacement found in [48] was qualitatively related to a Rouse description of a polymer brush under the boundary condition that constrains the tethered part of the polymer to the surface [49]. Under this assumption, all the modes but one are quickly damped. The only remaining mode has a relaxation behavior that can effectively be described with a simple exponential. The diffusion constant associated to this dynamics was shown to be twice the diffusion constant of a free chain in solution [49]. This is in fair agreement, for example, with the center-of-mass diffusion found for the PEG2000 solution of our study (see below). The fact that the Q dependence of $\Gamma_{\text{pol,cor}}$ has a linear diffusivelike behavior and that the associated diffusion coefficient is half of the diffusion of the free polymer suggests that the physical origin of this dynamics could be related to the $p = 1$ mode of the Rouse model for tethered polymer as described in [49]. This mode involves chain motions of the length scale of the end-to-end distance of the polymer [50].

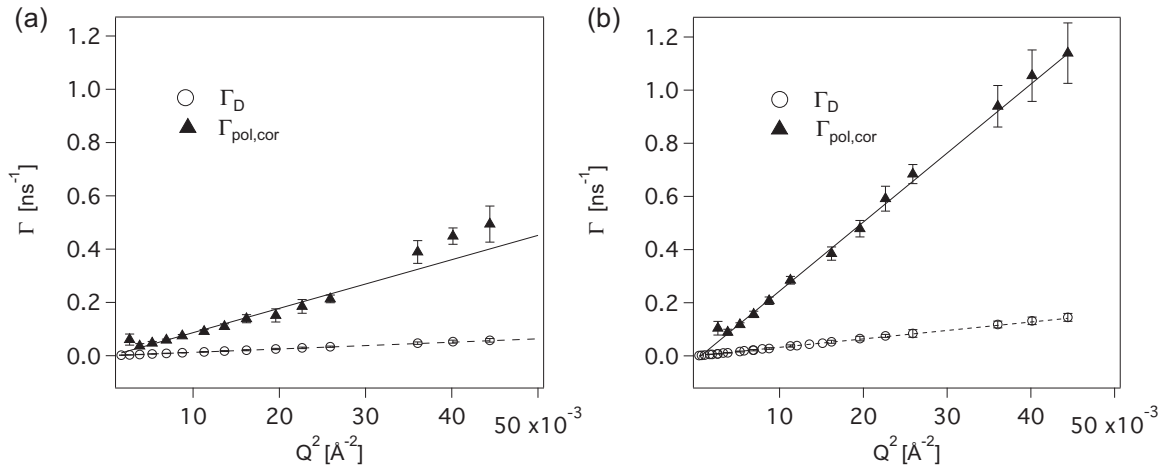


FIG. 11. The values of $\Gamma_D = 1/\tau_D$ and $\Gamma_{\text{pol,cor}}$ as a function of Q^2 obtained for PEG2000 AuNP dispersed in D_2O at the temperatures $T = 280$ K (a), $T = 318$ K (b).

If we tried to frame these dynamical results in a Zimm scenario, the relaxation times of the polymers would be related to the viscosity η_s of the liquid medium by the following relation: $\eta_{s,280\text{ K}}/\eta_{s,318\text{ K}} = [280\tau_1(Q)^{280\text{ K}}/318\tau_1(Q)^{318\text{ K}}] = R_z$ [50]. The average ratio between the viscosity of the D₂O between the two temperatures 280 and 318 K is ~ 2.6 whereas the $R_z \sim 2.3$. Even if the trend of the relaxation times with temperature is roughly compatible with what the Zimm model predicts, the absolute value of the viscosity that it is associated with these relaxation times ($\eta_{\text{solvent}} = K_B T \tau_{\text{pol}} Q^3 / 6\pi$) [15] as predicted by the Zimm model is around 0.004 Pa s at 280 K and 0.0016 Pa s at 318 K, i.e., roughly two times bigger than expected. A scaling of the relaxation time based only on the solvent viscosity neglects the important contribution to the dynamics related to the fact that the polymers are tethered to the NP surface. This has a significant impact on the resulting dynamics. This fact suggests that our results cannot be fully represented in the context of this commonly used model.

This $\Gamma \sim Q^2$ behavior was already reported for amphiphilic diblock copolymer micelles in aqueous solution like in the work of Matsuoka *et al.* [14] and Castelletto *et al.* [16]. In the latter, however, the Γ corresponding to the internal dynamics of the micelle did not intercept the origin. Our results are similar also to those found for polystyrene-butadiene diblock copolymer micelles in that the dynamical processes are ascribable to two relaxation times, one for the translational diffusion of the whole micelle and one accounting for the dynamics of the polymer corona. However, it differs in the fact that a Zimm behavior was used to describe the latter. In agreement with all the above cited work, the dynamics found for the polymer corona in this case is not compatible with the so called *breathing mode*, a collective mode of the polymer chains driven by a competition between the osmotic compressibility and the entropic force of the polymer chains predicted by de Gennes [51]. Such a dynamics is characterized by a distribution of relaxation times, at variance with our experimental observation that gives evidence of only one relaxation time associated with the polymer chains.

The fitting analysis of the NSE data for the PEG2000 solution dissolved in D₂O at 10% in weight illustrates also the presence of two dynamical processes with relaxation times τ_1 and τ_2 . By a close inspection of Figs. 10(a) and 10(b), we see that the faster relaxation times τ_1 measured for the homopolymer were very close to the relaxation times τ_{pol} obtained from the analysis of the NSE data of the PEG2000 AuNP at each Q . This is particularly evident at the highest temperature [Fig. 10(b)]. The slower dynamical process τ_2 is compatible to the relaxation time associated to the translational diffusion of a PEG2000 polymer in D₂O as measured by NMR experiments at similar concentration [47]. A direct comparison is, however, not obvious since the data in literature are performed up to 8% concentration and the extrapolation in this range of concentrations and temperatures is not trivial for the PEG2000. Therefore, we tentatively ascribed this relaxation time to the translational dynamics of the polymer in the solution. We will come back later to the consistency of this assumption. Previous attempts to fit the data corresponding to the PEG2000 homopolymers with a standard fitting procedure showed that these data could be fitted also using a single stretched exponential in the Q range 0.06–0.19 Å^{−1}. By let-

ting the β parameter free, the fit returned a value close to 0.9. A β parameter smaller than one is a sign of a more complex dynamic scenario involving a distribution of relaxation times [52]. Indeed, a distribution of relaxation times is very likely to occur in a macromolecular system such as a polymer in solution. As we have seen previously in this section, instead of considering a continuous distribution of relaxation times, one can represent this complexity in a simplified way by describing the dynamical properties of the system with a finite number of components. But, which is then the best way to describe our data?

With the aim of finding the model that supports our data with the highest statistical significance, we then run the fitting analysis by letting the RJ-MCMC algorithm use a model consisting of a linear combination of β -stretching exponentials [Eq. (10)] and leaving the algorithm free to estimate all the parameters. In looking for the best fit between the model and the data, the algorithm explored all the different possibilities such as sums of simple and/or stretched exponential functions, and clearly also the single stretched exponential function that, as we have seen, can still describe the data.

The prior distribution function for the β_j (with $j = 1 \dots n$) described in Sec. II C was used. Quite interestingly, the algorithm converged to the same combination of simple exponentials described in Sec. III B 1. The results found by the Bayesian algorithm showed quite clearly that the combination of one (only at very small Q) or two simple exponentials was indeed the most probable description of $I(Q, t)/I(Q, 0)$. At each Q , the algorithm reproduced within the errors the results obtained for a combination of simple exponentials. The estimated average values for β_1 and β_2 (when present) were both very close to one (typically $0.98/0.99 \pm 0.02$) and the posterior probability $P(\beta_1 = 1)$ and $P(\beta_2 = 1)$ were close to 100% and between 85% and 96%, respectively.

More interestingly, the solution with two single exponential functions that picks out a finite number of relaxation processes was preferred over the one with one single stretched exponential that is often used to describe the presence of multiple relaxation processes. We can conclude that the proposed method has allowed a sort of “*fine-structure*” analysis of the dynamic scenario for the polymer solutions providing also a probabilistic statistical support to the model adopted for the description of the time correlation function relaxations. We stress that by choosing arbitrarily to fit the data with just one Kohlrausch-Williams-Watts (KWW) function in aprioristic way, part of the information contained in the data might be lost. Moreover, in this particular case the choice of an apparently simpler model (one single stretched exponential rather than two exponentials) would not be even justified by some parameters’ economy criterion since in both cases the number of fitting parameters would be the same.

This analytical compact form is often used to represent NSE data in the context of the Zimm and the Rouse models [41], where the Langevin equation is solved through an expansion in normal coordinates. A relaxation time is associated to each of these normal modes. The KWW results from an average over the normal modes used to solve the Langevin equation. Although a stretched exponential function is ruled out by the algorithm as highly improbable, this by no means

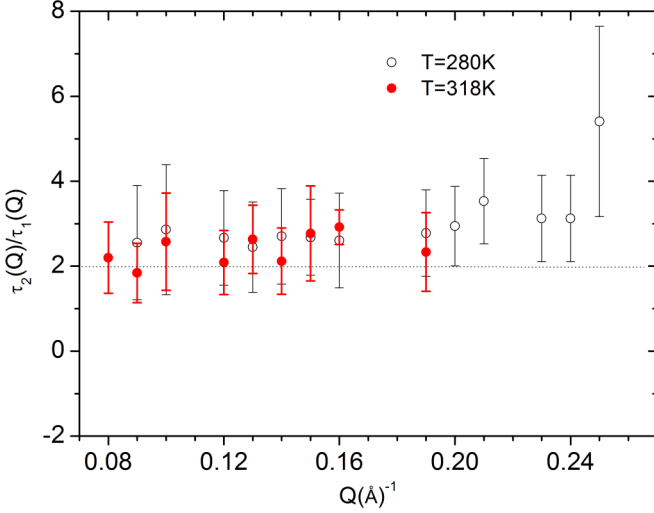


FIG. 12. Ratio between the two relaxation times estimated by the RJ-MCMC algorithm for the PEG2000 solution, at the two investigated temperatures.

implies that dynamic processes found in our analysis are at variance with what it is predicted by polymer theories.

In fact, an interesting experimental result of our analysis of homopolymers solutions is that the decay time ratio of the two exponential modes τ_1 and τ_2 found by the application of Eq. (9) has a constant value of about 2 (to within experimental uncertainties), as displayed in Fig. 12.

In the Rouse model [41,50], the $p=0$ normal mode corresponds to the center-of-mass diffusion with a diffusion coefficient D_R . All other modes ($p \geq 1$) account for the chain internal motions. Since the Bayesian analysis detected only two modes, including the overall polymer translational diffusion that we tentatively ascribed to the τ_2 relaxation, we worked out the $I(Q, t)$ [41] in the context of the Rouse model, taking into account only the first internal mode ($p=1$), that has the longest relaxation time. In particular, it can be shown that $I'_{\text{int}}(Q, t)$, i.e., the time-dependent part of the internal dynamics in Eq. (7), which now reduces to a single exponential, contributes to $I(Q, t)/I(Q, 0)$ with a term

$$\exp(-D_R Q^2 t) \exp\left(-\frac{t}{\tau_{p=1}^{\text{Rouse}}}\right) \quad (16)$$

with

$$\tau_{p=1}^{\text{Rouse}}(Q) = \frac{3\pi^2}{Nb^2} \frac{1}{Q^2} = \frac{1}{D_R Q^2}, \quad (17)$$

where $\tau_R = N^2 b^2 \zeta / 3\pi^2 k_B T$ is the Rouse time [50], $N=45$ is the number of monomer units of PEG2000, $b=3.5 \text{ \AA}$ [53] is the segment length in the Rouse chain, and ζ is the microscopic friction coefficient. Therefore, we finally find that $I(Q, t)/I(Q, 0)$ is given by the sum of two exponential terms with decay times expressed as $1/D_R Q^2$ and $1/2D_R Q^2$, respectively. Hence, in the context of Rouse model and under the approximation used, the ratio between the relaxation times τ_2 and τ_1 as expressed in Eq. (9) is 2, in fair agreement with the experimental finding. Furthermore, this result provides us an internal consistency proof that τ_2 can be confidently attributed

to the polymer center-of-mass diffusion and allow us to calculate from their values average values of $D_R \sim 12 \text{ \AA}^2/\text{ns}$ and $5 \text{ \AA}^2/\text{ns}$ at 318 and 280 K, respectively, and from the relation $\zeta = k_B T / N D_R$ a ratio $\zeta_{280 \text{ K}} / \zeta_{318 \text{ K}} \sim 2.1$. If again a Rouse context is assumed to model the internal dynamics of the polymer chains in solution, the same ratio between the friction coefficients at the two investigated temperatures can be calculated using the estimated relaxation times $\tau_{\text{int},1}$ [see Eq. (7)] and they follow the temperature scaling relationship $280 \tau_{\text{int},1}(Q)^{280 \text{ K}} / 318 \tau_{\text{int},1}(Q)^{318 \text{ K}} = \zeta_{280 \text{ K}} / \zeta_{318 \text{ K}} \sim 2.1$. This shows a further internal consistency of the rescaling in temperatures of relaxation times according to the Rouse model.

Unlike for the PEG2000 AuNP, for the PEG 400 AuNP, $I(Q, t)/I(Q, 0)$ could be always described only with one term corresponding to the translational dynamics of the NP, without the need of including a term accounting for the polymer dynamics. This could be due to different reasons. The net amount of the polymer, as well as the volume fraction of polymer with respect to that of the whole NP, is significantly lower for PEG400 AuNP compared to PEG2000 AuNP. As a matter of fact, if we approximate the volume of the NP to $4/3\pi R_h^3$ and if we consider that the radius of the gold core is $R_{\text{Au}} = 25 \text{ \AA}$ we have that in the case of the PEG2000 AuNP the hydrated polymer corona occupies ~ 0.90 of the NP volume, whereas only 0.70 in the case of PEG400 AuNP. Moreover, the concentration of the NP in the solution was 5% in weight for the PEG2000 AuNP and only 2% in weight for the PEG400 AuNP. Clearly, the contribution of the polymer corona in the two experiments was significantly different, in particular the scattering intensity due to the regions where polymer was present was much higher in the case of the PEG2000 AuNP than in the case of the PEG400 AuNP. Furthermore, the condition $Q R_h \gg 1$ that determines the Q range where the polymer dynamics dominates with respect to the translational diffusion is less effectively fulfilled in the case of the PEG400 AuNP due to the significantly smaller value of R_h . The combination of these two conditions might have obscured the scattering of the polymers in the case of the PEG400 AuNP. The lack of evidence of a polymer dynamics in our NSE experiment could also be due to an effect of confinement. The short PEG400 chains are anchored to the surface of the NP with an area per molecule around $14 \text{ \AA}^2$ [13]. The effect of the anchor on a solid surface combined with the presence of neighboring polymers highly packed might slow down the fluctuation of the PEG400 to values outside the experimental window of the instrument. Unlike the PEG400 that extends only up to $25 \text{ \AA}$ away from the surface, the PEG2000 extends up to $60 \text{ \AA}$ away from the surface, allowing a significant part of the molecule to fluctuate in the larger volume available.

The NSE data measured on the 5% solution of PEG400 in D_2O evidenced only one relaxation process in the time window of the experiment. This relaxation corresponds very well with the NMR results concerning the translational diffusion of the polymer at the same concentration and temperatures [47]. A possible explanation for the absence of internal modes might be related to the fact that the relaxation time corresponding to the faster translational diffusion of

PEG400 (compared to PEG2000) might be comparable with the relaxation time of the internal dynamics of PEG400. If the two processes occur at a very similar timescale, it would be very difficult to distinguish between them.

V. CONCLUSIONS

We have exploited a Bayesian approach together with the very important property of decomposition of any time correlation function in terms of exponential functions to explore the nanosecond dynamics of polymer-coated nanoparticles dissolved in D₂O in comparison with the dynamics of equivalent polymers free in solution measured through NSE spectroscopy. After decomposing the intermediate scattering function in elementary components, we associated each component to a particular kind of dynamics (center-of-mass translational diffusion or internal polymeric dynamics) both for the grafted nanoparticles and for the homopolymer solutions, by comparing the relaxation times obtained from the NSE analysis to those obtained with DLS and NMR data from literature and checking the agreement with polymer science theory. For the PEG AuNP we disentangled the translational dynamics of the whole nanoparticle from that of the grafted polymer chains. The dynamic phenomena observed depended strongly on the molecular weight of the tethered polymer. In the case of longer PEG2000 grafted chains we could clearly see the translational dynamics of the NPs and the internal dynamics of the polymer chains. The polymers did not follow a Zimm dynamics in the range of times and momentum transfer explored, nor the often invoked breathing modes. Instead, our analysis indicates that the polymer corona relaxation follows a pure exponential decay, which appears to be an important confirmation of the behavior predicted by coarse grained molecular dynamics simulations [48] and consistent with a Rouse description of a tethered polymer brush where all the Rouse modes are damped but the longest one related to a movement over the entire polymer chain length [49]. In the case of shorter PEG400 grafted chains, only the translational dynamics of the NP was visible. This could be due to the significantly lower amount of polymeric material (and the NSE signal associated with it) in this case, or due to a high mechanical constraint of the PEG400 in proximity of the gold surface.

As for polymer solutions, the analysis of the NSE data evidenced two relaxation times in the case of longer PEG2000 polymers associated to the translational dynamics of the polymer and to the $p = 1$ Rouse mode, respectively. The proposed analysis establishes that aside from the center-of-mass molecule diffusion, only one relaxation time is statistically supported and this is then necessarily described as a simple exponential decay. As an important outcome, the more common formal description of the time correlation function decay in terms of stretched exponential, as in the averaged form of Zimm and Rouse theory, is excluded by the joint posterior

distribution of the adopted model conditional on the observed data. Nevertheless, the internal relaxation time is absolutely consistent with the dominant Rouse time for a polymer free chain and effectively both estimated relaxation times scale with temperature according the Rouse theory. In the case of the shorter PEG400, only one relaxation time comparable with the translational diffusion of the polymer was evidenced.

Our methodology addresses one of the most important issues in the analysis of inelastic and quasielastic neutron scattering data: the arbitrariness in the choice of the number of components needed to describe the studied dynamical processes. In conventional fitting analysis, a particular model is imposed, based on a reasonable physical hypothesis and without any knowledge of a posterior probability that this choice of the model is statistically significant. Hence, the physical hypothesis, although well motivated, affects substantially the result of the fitting analysis. In the specific case of NSE experiments, nISF is often described in terms of stretched exponential that can account for the superposition of more relaxation mechanisms, in terms of an effective relaxation process modulated by a β parameter. In view of our lack of knowledge about the most reliable model describing the data, no such invasive hypothesis was done to run the RJ-MCMC algorithm used in this work to explain the experimental data. We have shown, for example, how it is possible to make explicit the richness of the dynamic properties of a system such as the PEG2000 solution that would be instead hidden if nISF's were modeled as a stretching exponential function. At the same time, we provided a probabilistic evidence about the model used to describe the experimental data and without sacrificing the simplicity of the model description. In doing so, we obtained that the output of the Bayesian analysis complies fully with the theoretically demonstrated general result on the exponential expansion of any time correlation function [28–30], which has been here applied to the description of the dynamics of a macromolecular system at the nanosecond timescale typically investigated by means of NSE. We believe that the approach presented here will serve as an effective tool that can be used in general to provide an unbiased interpretation of neutron spin echo data.

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