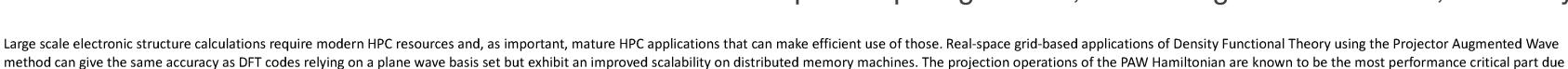


A Spherical Harmonic Oscillator Basis for Reduced Bandwidth Requirements

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to their limitation by the available memory bandwidth. We investigate on the usability of a 3D factorizable basis of Hermite functions for the atomic PAW projector functions that allows to reduce and nearly to remove the bandwidth requirements for the grid

representation of the projector functions in projection operations. This increases the fraction of exploitable floating-point operations on modern vectorized many-core architectures, like GPUs, by raising the arithmetic intensity of such operations.





1D Harmonic Oscillator

$$\hat{H}_{\text{HO}} = \hat{T}_{\text{kinetic}} + \hat{V}_{\text{potential}} = -\frac{\partial^2}{\partial x^2} + x^2$$

$$\psi_{\text{HO}}(x) \sim H_n(x) \cdot \exp(-x^2/2)$$

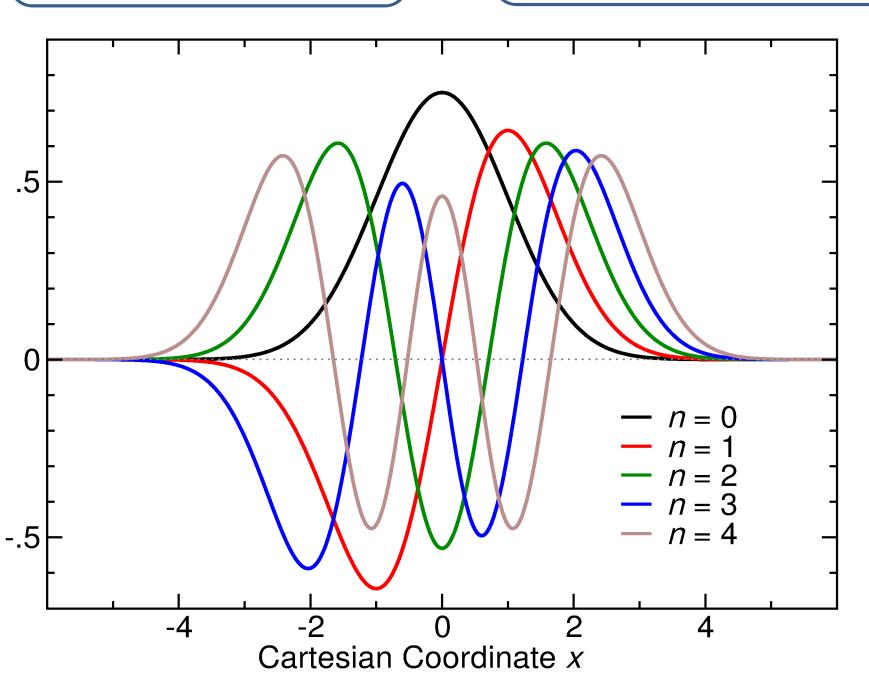
Hermite polynomial · Gaussian



Ch. Hermite



J. C. F. Gauss



Recursive definition of Hermite polynomials

$$H_{n+1}(x) = xH_n(x) - \frac{n}{2}H_{n-1}(x)$$

$$H_0(x) = 1, \ H_1(x) = x, \ H_2(x) = x^2 - \frac{1}{2}$$

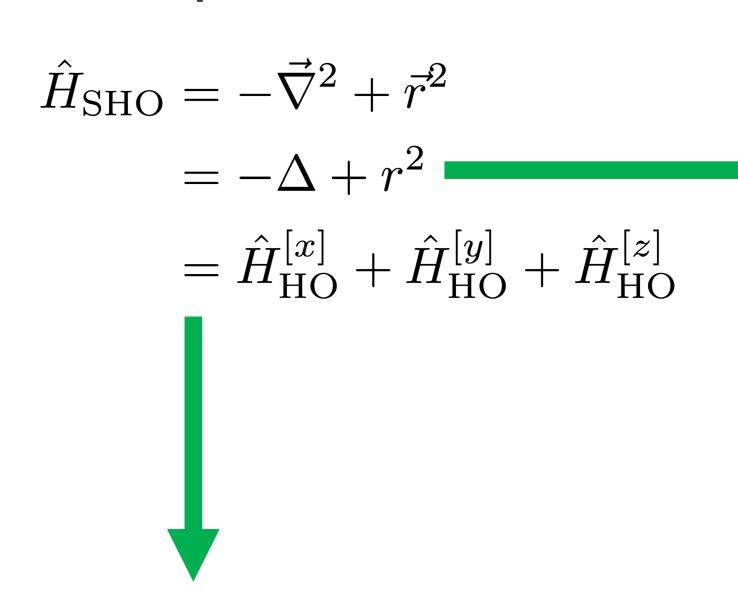
Application to Electronic Structure Calculations on Cartesian Grids

For reasons of parallel scalability, the quantum mechanical wave functions in Density Functional Theory (DFT) are represented on 3D Cartesian grids. Using the Projector Augmented Wave (PAW) method [Blöchl95], a non-local operator accurately describes the scattering of valence electrons at the ions.

$$\hat{V}_{\text{ion}} = |\tilde{p}_i\rangle D_{ij} \langle \tilde{p}_j|$$

The non-locality can be implemented as projection and expansion operations with a set of localized functions refered to as projectors. If we manage to represent the projectors in a small SHO basis, we can omit storing and loading of precomputed projector function values in both, projection and expansion operation. This releases memory capacity constraints and, even more important, reduces the requirements for memory bandwidth which has become the most costly resource in HPC. The 1D HO basis functions are cheap enough to be generated on-the-fly.

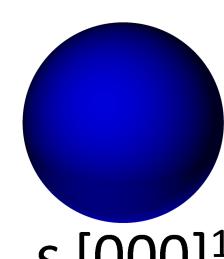
3D Spherical Harmonic Oscillator



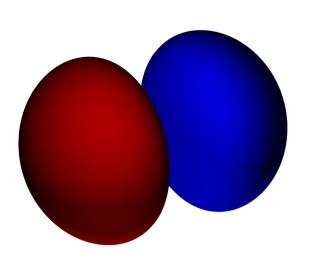
Cartesian Representation

$$\Psi_{\text{SHO}}(x, y, z) = \psi_{n_x}(x)\psi_{n_y}(y)\psi_{n_z}(z)$$

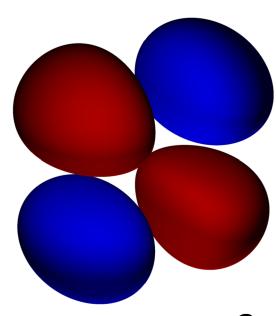
as 3D Cartesian product of 1D-HO eigenfunctions → Tensor compression



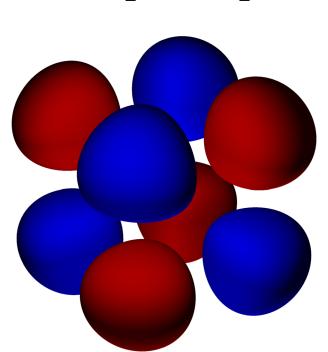
 $s [000]^1$



 $p [100]^3$

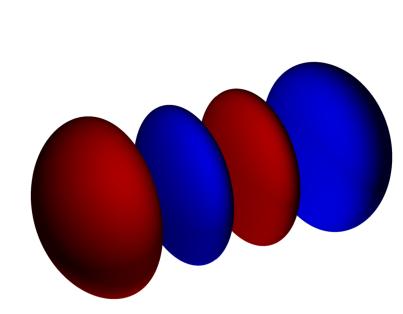


 $d [110]^3$



 $f[111]^1$

 $[200]^3$



 $[300]^3$

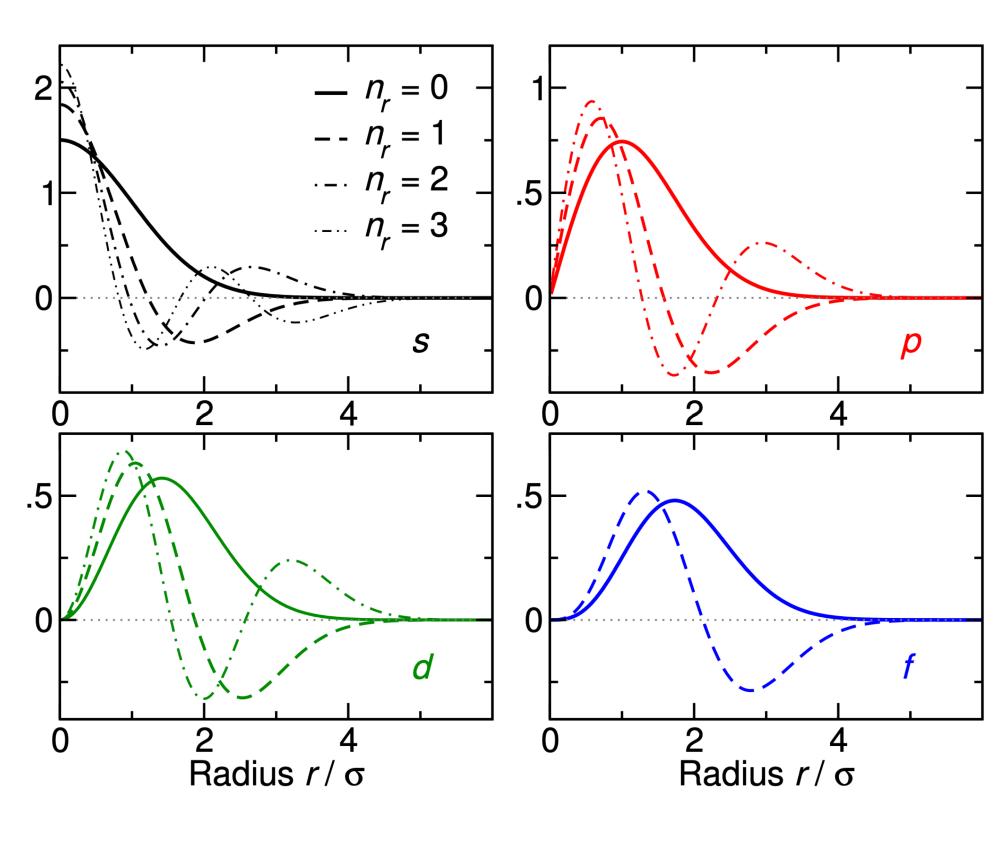
Radial Representation

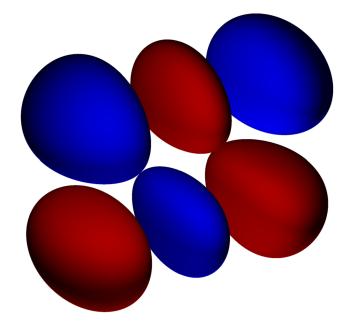
$$\Psi_{\rm SHO}(r, \vartheta, \varphi) = R_{n_r \ell}(r) \cdot Y_{\ell m}(\vartheta, \varphi)$$



Spherical Harmonics

$$R_{n_r\ell}(r) \sim r^{\ell} L_{n_r}^{(\ell+\frac{1}{2})}(r^2) \exp(-r^2/2)$$





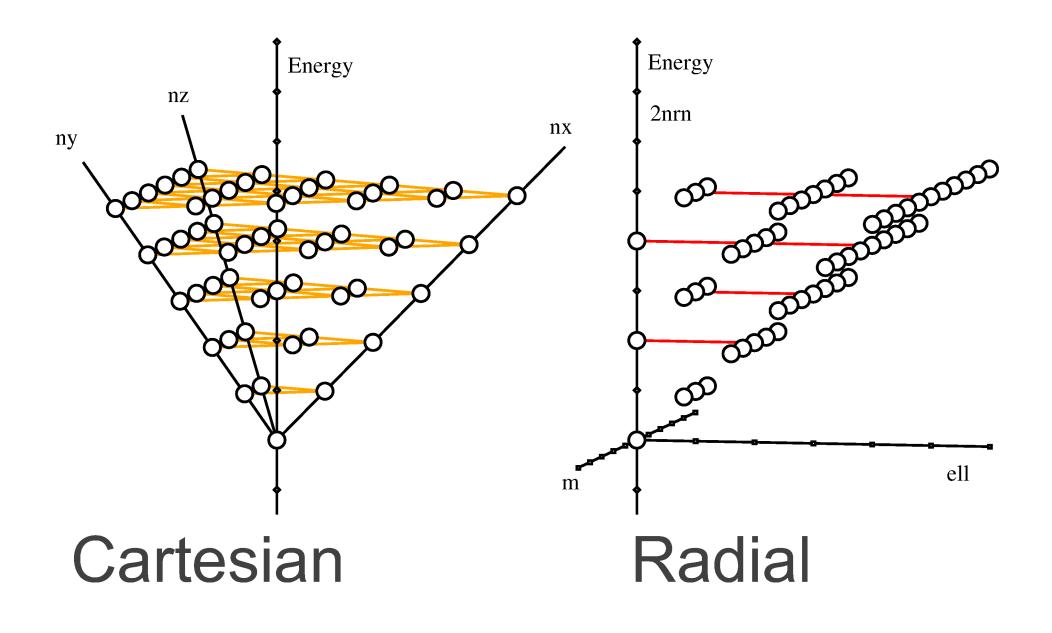
 $[210]^6$

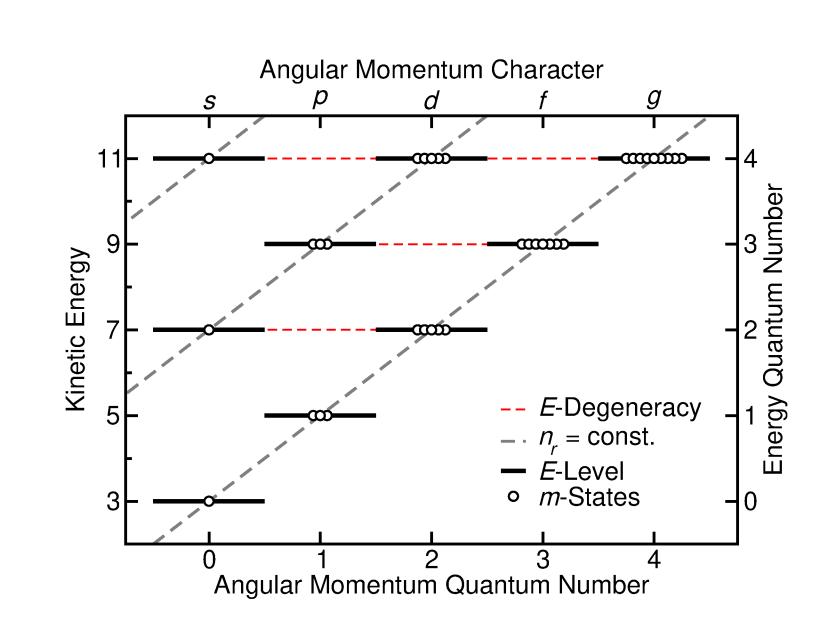


Paper@arXiv 1905.00480

$$E = 2(n_x + n_y + n_z) + 3$$
 Energy Degeneracy

$$E = 2(\ell + 2n_r) + 3$$





Unitary Transformation in energy-degenerate subspaces

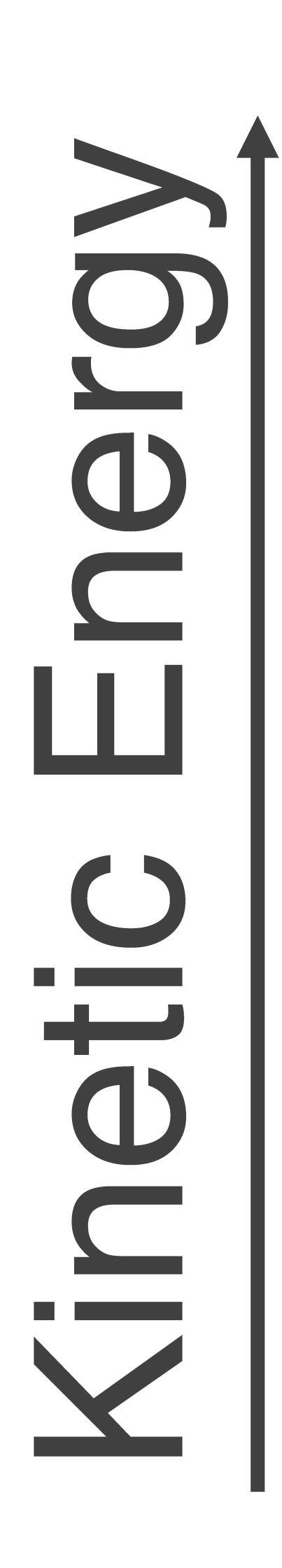
$$\Psi_{n_x n_y n_z}(x, y, z) \equiv \sum_{n_r \ell m} U_{n_x n_y n_z}^{n_r \ell m} \ \Psi_{n_r \ell m}(r, \vartheta, \varphi)$$

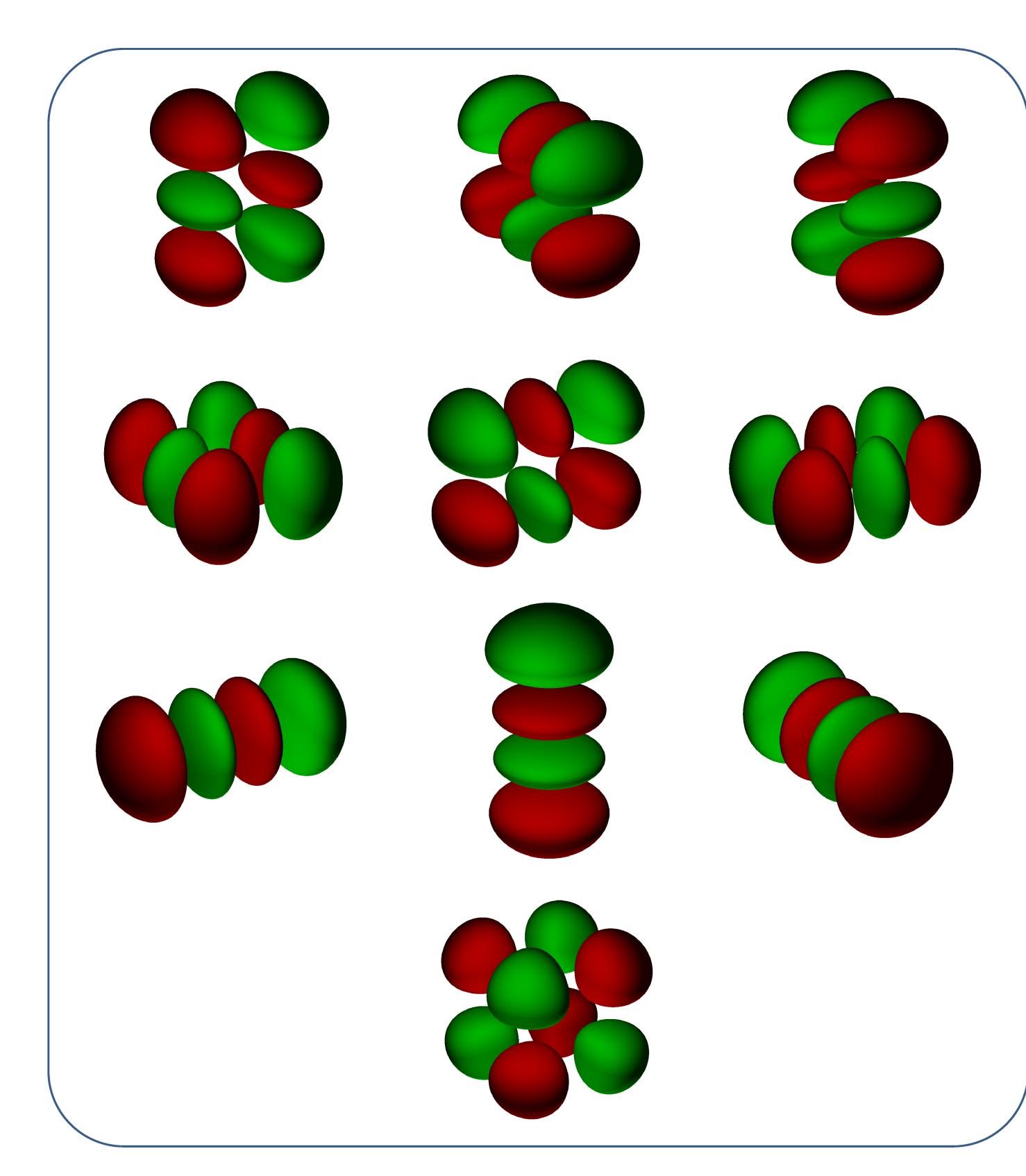
Angular momentum characters

 $s^* \oplus d$, $p^* \oplus f$, $s^{**} \oplus d^* \oplus g$, ...

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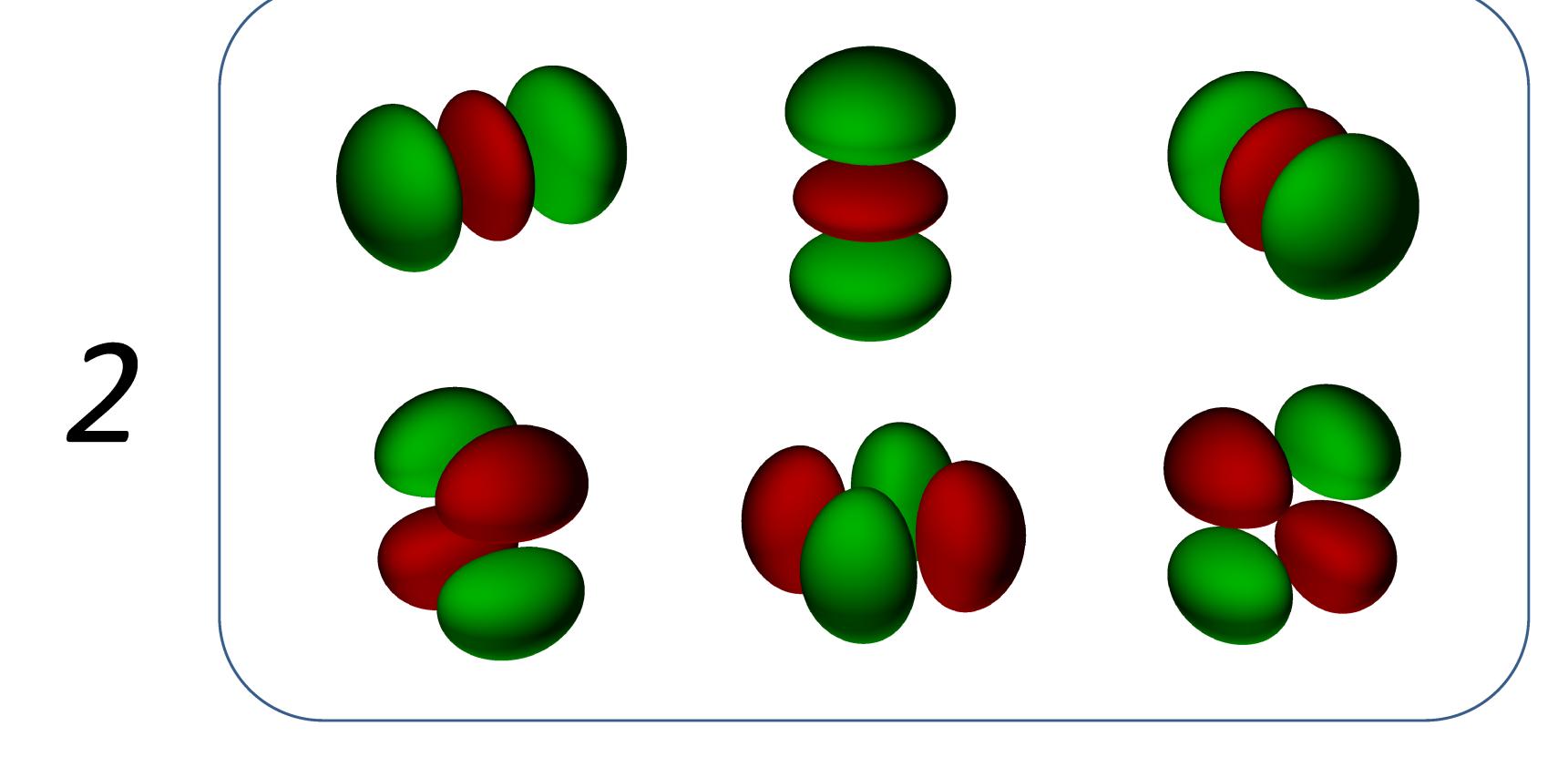






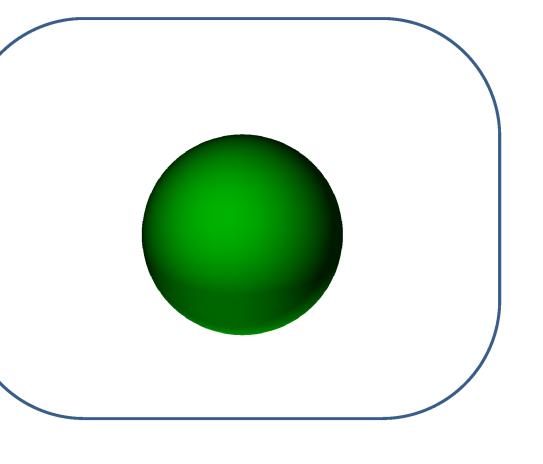
fand p*

fxyz



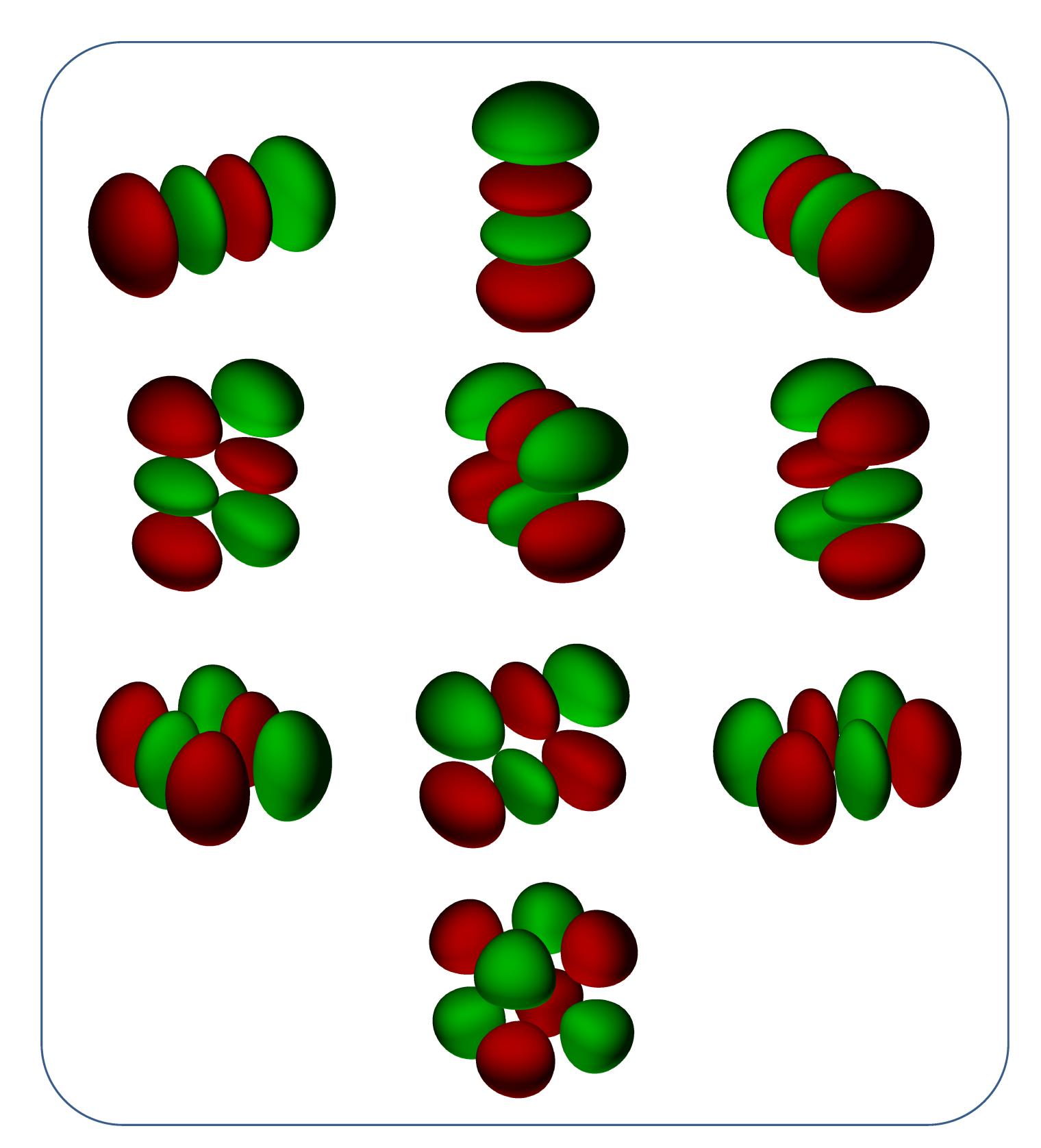
eg and s*

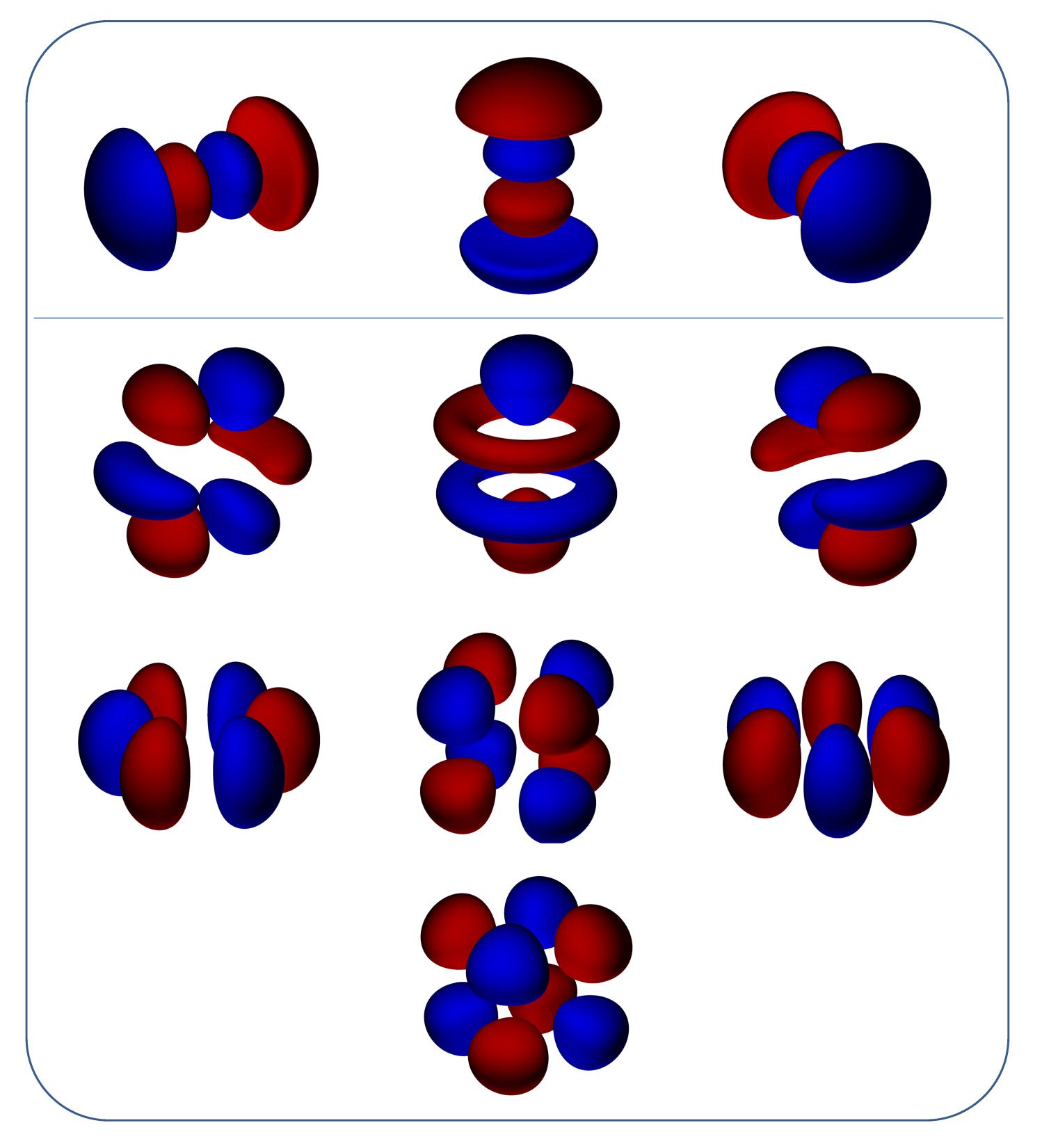
 t_{2g}



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Cartesian factorizable



