

Update on $|V_{us}|$ and $|V_{us}/V_{ud}|$ from semileptonic kaon and pion decaysChien-Yeah Seng¹, Daniel Galviz¹, William J. Marciano², and Ulf-G. Meißner^{1,3,4}¹*Helmholtz-Institut für Strahlen- und Kernphysik and Bethe Center for Theoretical Physics, Universität Bonn, 53115 Bonn, Germany*²*Department of Physics, Brookhaven National Laboratory, Upton, New York 11973, USA*³*Institute for Advanced Simulation, Institut für Kernphysik and Jülich Center for Hadron Physics, Forschungszentrum Jülich, 52425 Jülich, Germany*⁴*Tbilisi State University, 0186 Tbilisi, Georgia*

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Implications for Cabibbo universality based on progress in the study of semileptonic kaon and pion decays are discussed. Included are recent updates of experimental input along with improved radiative corrections, form factors, and isospin-breaking effects. As a result, we obtain for the Cabibbo-Kobayashi-Maskawa (CKM) quark mixing matrix element $|V_{us}| = 0.22309(56)$ from semileptonic $K_{\ell 3}$ ($K \rightarrow \pi \ell \nu$) decays and $|V_{us}/V_{ud}| = 0.22908(87)$ from the ratio between the kaon and pion (π_{e3}) semileptonic decay rates. In both, a lattice QCD value of the form factor $|f_+^K(0)| = 0.9698(17)$ is employed. The V_{us} from $K_{\ell 3}$ decays together with $V_{ud} = 0.97373(31)$ found from superallowed nuclear beta decays implies an apparent 3.2σ violation of the first-row CKM unitarity condition. The $|V_{us}|/|V_{ud}|$ obtained from the ratio of weak vector current induced meson decays is consistent with the observed unitarity violation but found to differ by 2.2σ from its extraction using the ratio of weak axial-vector leptonic decay rates $\Gamma(K \rightarrow \mu \nu)/\Gamma(\pi \rightarrow \mu \nu)$. The situation suggests a difference between vector and axial-vector derived CKM matrix elements or a problem with the lattice QCD form factor input. Prospects for future improvements in comparative precision tests involving $|V_{ud}|$, $|V_{us}|$, and their ratios are briefly described.

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I. INTRODUCTION

Precision tests of the Standard Model (SM) have become increasingly important given the null results from high-energy colliders in the search for physics beyond the Standard Model (BSM) [1]. Deviation from expectations in the muon anomalous magnetic moment [2–6], hints of lepton flavor universality violation in B decays [7–10], and tests of unitarity in the Cabibbo-Kobayashi-Maskawa (CKM) quark mixing matrix are all exhibiting potential BSM effects. In the last case, improvements in the electroweak radiative corrections (RCs) [11–18] have revealed tension in the first-row unitarity requirement $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$ [19]. Similarly, the difference in the value of $|V_{us}|$ extracted from $K_{\ell 3}$ and $K_{\mu 2}$ decays needs to be better understood.¹ Possible explanations based on BSM calculations [20–44] have been conjectured.

¹All decay processes are understood to be radiative inclusive; for example, K_{e3} means $K \rightarrow \pi \ell \nu(\gamma)$.

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The confirmation of such ideas will require improvements in both experiment and SM theory.

In this work, we focus on the extraction of $|V_{us}|$ from $K \rightarrow \pi \ell \nu$ semileptonic decay processes known as $K_{\ell 3}$. Our starting point is the comprehensive 2010 FlaviaNet Working Group Report [45] which presented a thorough review of all relevant experimental and theoretical information available at the time. Since then, significant progress has been made on both the experimental and theoretical fronts, including new measurements of the K_S lifetime [46,47] and branching ratio (BR) [48], updates of the electroweak RCs [49–51], phase space factors [52], $K\pi$ form factor, and isospin-breaking corrections ([53] and references therein). Some of those advances have been more recently discussed by FlaviaNet Working Group members in the form of proceedings [54] and conference slides [52,55] rather than more detailed research publications, making cross-checking difficult. For that reason, we present in this paper an updated status report on the value of $|V_{us}|$ extracted from $K_{\ell 3}$ decay properties. With the recent theory progress, we find that apart from the lattice input of the $K\pi$ form factor which is a universal multiplicative constant, the experimental errors from the kaon lifetimes and BRs are by far the dominant sources of uncertainties in all six channels of $K_{\ell 3}$, which is quite different from the

situation a few years ago, in particular, before the new calculation of the K_{e3} RC [50,51].

The progress described above can also be applied to the determination of the recently proposed ratio $R_V = \Gamma(K_{e3})/\Gamma(\pi_{e3})$ as an alternative approach to study $|V_{us}/V_{ud}|$ [56], complementary to the existing method based on $R_A = \Gamma(K_{\mu 2})/\Gamma(\pi_{\mu 2})$. We show, following the recent improvements in the precision level of the K_{e3} [50,51] and π_{e3} RC [57], that R_V is now a theoretically cleaner observable than R_A . It provides strong motivation to further improve the experimental precision of the π_{e3} BR and K_{e3} decay properties as much as possible. An experimental next-generation rare pion decay program proposal, PIONEER [58,59], would aim to improve the experimental π_{e3} decay rate by a factor of 3 or better, making the K_{e3} decay rate the dominant uncertainty in R_V .

The paper is organized as follows: In Sec. II, we give a detailed update on $|V_{us}f_+^K(0)|$ and $|V_{us}|$. Similarly, Sec. III provides updates on $|V_{us}f_+^K(0)|/|V_{ud}f_+^\pi(0)|$ and $|V_{us}/V_{ud}|$. We end with conclusions in Sec. IV. Some technicalities are relegated to the Appendix.

II. UPDATES ON $|V_{us}f_+^K(0)|$ AND $|V_{us}|$

We start from the quantity $|V_{us}f_+^K(0)|$, with $f_+^K(0)$ the vector form factor in the $K^0 \rightarrow \pi^-$ transition at zero momentum transfer. It is extracted from K_{e3} decays through the following master formula [1]:

$$A_{K\ell} \equiv |V_{us}f_+^K(0)|_{K\ell} = \sqrt{\frac{192\pi^3 \text{BR}(K\ell)\Gamma_K}{G_F^2 M_K^5 C_K^2 S_{\text{EW}} I_{K\ell} (1 + \delta_{\text{EM}}^{K\ell} + \delta_{\text{SU}(2)}^{K\ell})}}, \quad (1)$$

where $G_F = 1.1663787(6) \times 10^{-5} \text{ GeV}^{-2}$ is the Fermi constant obtained from muon decay, Γ_K the total kaon decay width, $\text{BR}(K\ell)$ the K_{e3} branching ratio, and C_K a simple isospin factor which equals 1 ($1/\sqrt{2}$) for K^0 (K^+) decay. The SM theory inputs to the right-hand side are as follows: $S_{\text{EW}} = 1.0232(3)_{\text{HO}}$ is a universal short-distance EW factor [60], and the uncertainty comes from higher-order QED effects [61] which is common to all channels and will not take part in the weighted average. Meanwhile, $I_{K\ell}$, $\delta_{\text{EM}}^{K\ell}$ and $\delta_{\text{SU}(2)}^{K\ell}$ are the phase space factor, the long-distance electromagnetic (EM) correction, and the isospin-breaking correction, respectively. These are channel-specific inputs. To facilitate the discussion of correlation effects, we group the values of $A_{K\ell}$ from six independent K_{e3} channels into a vector,

$$A = (A_{K_L e}, A_{K_S e}, A_{K^+ e}, A_{K_L \mu}, A_{K_S \mu}, A_{K^+ \mu})^T. \quad (2)$$

The order of the entries is important, as is seen later.

In what follows, we summarize all the data input needed in this work. All the experimental data of decay widths and

BRs are obtained from the 2021 online update [62] of the Particle Data Group (PDG) review [1]. Knowing that different choices of inputs of statistical analyses of the same problem may lead to different quantitative conclusions, throughout the discussion we clearly explain the similarities and differences in the data inputs between our work and existing global analysis (in particular, the 2010 FlaviaNet review [45] and its updates [52,54,55]) and present all the essential steps in some detail despite that most of them are familiar to experts; the basic mathematical tools needed in this work are summarized in the Appendix. With such, all the intermediate and final results in this paper are fully transparent and can be easily crossed-checked by interested readers or compared to similar analyses with possibly different statistical approaches.

A. K_L experimental inputs

The K_L decay width quoted in PDG is obtained through a combined fit from Refs. [63–65], whereas the $K_L e$, $K_L \mu$ BRs are fitted from Refs. [63,66], all which have been included in the FlaviaNet 2010 review. The results read

$$\begin{aligned} K_L^{\text{exp}} &= (\text{BR}(K_L e) \text{BR}(K_L \mu) \Gamma_{K_L})^T \\ &= (0.4055(11) \ 0.2704(7) \ 1.2866(53) \times 10^{-14} \text{ MeV})^T, \end{aligned} \quad (3)$$

with the correlation matrix (it is symmetric, so we only show the upper right components for simplicity),

$$\text{Corr}(K_L^{\text{exp}}) = \begin{pmatrix} 1 & -0.2149932 & -0.2676291 \\ & 1 & -0.08759115 \\ & & 1 \end{pmatrix}, \quad (4)$$

from which we can obtain the covariance matrix $\text{Cov}(K_L^{\text{exp}})$ using Eq. (A4). The contribution of K_L^{exp} to the covariance matrix of A is given by

$$\text{Cov}(A)_{K_L^{\text{exp}}} = \left(\frac{\partial A}{\partial K_L^{\text{exp}}} \right) \cdot \text{Cov}(K_L^{\text{exp}}) \cdot \left(\frac{\partial A}{\partial K_L^{\text{exp}}} \right)^T, \quad (5)$$

where

$$\left(\frac{\partial A}{\partial K_L^{\text{exp}}} \right) = \begin{pmatrix} \frac{\partial A_{K_L e}}{\partial \text{BR}(K_L e)} & 0 & \frac{\partial A_{K_L e}}{\partial \Gamma_{K_L}} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & \frac{\partial A_{K_L \mu}}{\partial \text{BR}(K_L \mu)} & \frac{\partial A_{K_L \mu}}{\partial \Gamma_{K_L}} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}. \quad (6)$$

B. K_S experimental inputs

The K_S decay width quoted in PDG is fitted from Refs. [46,47,67–70] and the $K_S e$ BR from Refs. [71–73]. Notice that not all of them were utilized in the FlaviaNet review and updates but only the more recent results from the KTeV [46], KLOE [47,72], and NA48 [67,71]

$$\begin{aligned} K_S^{\text{exp}} &= (\text{BR}(K_S e) \quad \text{BR}(K_S \mu) \quad \Gamma_{K_S})^T \\ &= (7.04(8) \times 10^{-4} \quad 4.56(20) \times 10^{-4} \quad 7.3510(33) \times 10^{-12} \text{ MeV})^T. \end{aligned} \quad (7)$$

The PDG provides the correlation coefficient between the two BRs but not between the BRs and the total decay width. The latter is expected to be very small from Ref. [45], so here we simply set it to zero. With that, we obtain the following correlation matrix:

$$\text{Corr}(K_S^{\text{exp}}) = \begin{pmatrix} 1 & -0.00144257 & 0 \\ & 1 & 0 \\ & & 1 \end{pmatrix}. \quad (8)$$

The contribution of K_S^{exp} to the covariance matrix of A is given by

$$\text{Cov}(A)_{K_S^{\text{exp}}} = \left(\frac{\partial A}{\partial K_S^{\text{exp}}} \right) \cdot \text{Cov}(K_S^{\text{exp}}) \cdot \left(\frac{\partial A}{\partial K_S^{\text{exp}}} \right)^T, \quad (9)$$

collaborations. Also, only five channels in $K_{\ell 3}$ were analyzed in all the past reviews because the $K_S \mu$ BR was not independently measured. The first direct measurement of this BR appeared in year 2020 [48], which allows us to finally include all six channels in the combined analysis of V_{us} for the first time. The experimental results read

where

$$\left(\frac{\partial A}{\partial K_S^{\text{exp}}} \right) = \begin{pmatrix} 0 & 0 & 0 \\ \frac{\partial A_{K_S e}}{\partial \text{BR}(K_S e)} & 0 & \frac{\partial A_{K_S e}}{\partial \Gamma_{K_S}} \\ 0 & 0 & 0 \\ 0 & \frac{\partial A_{K_S \mu}}{\partial \text{BR}(K_S \mu)} & \frac{\partial A_{K_S \mu}}{\partial \Gamma_{K_S}} \\ 0 & 0 & 0 \end{pmatrix}. \quad (10)$$

C. K^+ experimental inputs

The K^+ decay width quoted in PDG is fitted from Refs. [74–78] and the $K^+ e$, $K^+ \mu$ BRs from Refs. [79,80]. All of them, except two earlier experiments [77,80], were also used in the FlaviaNet analysis. The results read

$$\begin{aligned} K_+^{\text{exp}} &= (\text{BR}(K^+ e) \quad \text{BR}(K^+ \mu) \quad \Gamma_{K^+})^T \\ &= (5.07(4) \times 10^{-2} \quad 3.352(33) \times 10^{-2} \quad 5.3167(86) \times 10^{-14} \text{ MeV})^T, \end{aligned} \quad (11)$$

with the correlation matrix

$$\text{Corr}(K_+^{\text{exp}}) = \begin{pmatrix} 1 & 0.8959847 & 0.01425396 \\ & 1 & 0.01376368 \\ & & 1 \end{pmatrix}. \quad (12)$$

The contribution of K_+^{exp} to the covariance matrix of A is given by

$$\text{Cov}(A)_{K_+^{\text{exp}}} = \left(\frac{\partial A}{\partial K_+^{\text{exp}}} \right) \cdot \text{Cov}(K_+^{\text{exp}}) \cdot \left(\frac{\partial A}{\partial K_+^{\text{exp}}} \right)^T, \quad (13)$$

where

$$\left(\frac{\partial A}{\partial K_+^{\text{exp}}} \right) = \begin{pmatrix} 0 & 0 & 0 \\ \frac{\partial A_{K^+ e}}{\partial \text{BR}(K^+ e)} & 0 & \frac{\partial A_{K^+ e}}{\partial \Gamma_{K^+}} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & \frac{\partial A_{K^+ \mu}}{\partial \text{BR}(K^+ \mu)} & \frac{\partial A_{K^+ \mu}}{\partial \Gamma_{K^+}} \end{pmatrix}. \quad (14)$$

D. Phase space factor

The phase space factor is defined as²

²Notice that there is a typo in the $I_{K\ell}$ formula in many important references, e.g., Refs. [45,52,55,81].

$$I_{K\ell} = \int_{m_\ell^2}^{(M_K - M_\pi)^2} \frac{dt}{M_K^8} \bar{\lambda}^{3/2} \left(1 + \frac{m_\ell^2}{2t}\right) \left(1 - \frac{m_\ell^2}{t}\right)^2 \left[\bar{f}_+^2(t) + \frac{3m_\ell^2 \Delta_{K\pi}^2}{(2t + m_\ell^2) \bar{\lambda}} \bar{f}_0^2(t) \right], \quad (15)$$

where $\bar{\lambda} = [t - (M_K + M_\pi)^2][t - (M_K - M_\pi)^2]$ and $\Delta_{K\pi} = M_K^2 - M_\pi^2$. It probes the t dependence of the (rescaled) $K\pi$ vector and scalar form factors $\bar{f}_{+,0}(t)$, which are obtained by fitting to the $K_{\ell 3}$ Dalitz plot. There are different ways to parametrize the form factors, including the Taylor expansion [45], the z parametrization [82], the pole parametrization [83], and the dispersive parametrization [84–86]. In this work, we take the results of the dispersive parametrization from the latest FlaviaNet updates which claim the smallest uncertainty [52,55,87],

$$I_K = (I_{K^0 e} \quad I_{K^+ e} \quad I_{K^0 \mu} \quad I_{K^+ \mu})^T \\ = (0.15470(15) \quad 0.15915(15) \quad 0.10247(15) \quad 0.10553(16))^T, \quad (16)$$

with the correlation matrix [87],

$$\text{Corr}(I_K) = \begin{pmatrix} 1 & 1 & 0.530 & 0.521 \\ & 1 & 0.530 & 0.521 \\ & & 1 & 1 \\ & & & 1 \end{pmatrix}. \quad (17)$$

The contribution of I_K to the covariance matrix of A is given by

$$\text{Cov}(A)_{I_K} = \left(\frac{\partial A}{\partial I_K} \right) \cdot \text{Cov}(I_K) \cdot \left(\frac{\partial A}{\partial I_K} \right)^T, \quad (18)$$

where

$$\left(\frac{\partial A}{\partial I_K} \right) = \begin{pmatrix} \frac{\partial A_{K^0 e}}{\partial I_{K^0 e}} & 0 & 0 & 0 \\ \frac{\partial A_{K^+ e}}{\partial I_{K^0 e}} & 0 & 0 & 0 \\ 0 & \frac{\partial A_{K^+ e}}{\partial I_{K^+ e}} & 0 & 0 \\ 0 & 0 & \frac{\partial A_{K^0 \mu}}{\partial I_{K^0 \mu}} & 0 \\ 0 & 0 & \frac{\partial A_{K^+ \mu}}{\partial I_{K^0 \mu}} & 0 \\ 0 & 0 & 0 & \frac{\partial A_{K^+ \mu}}{\partial I_{K^+ \mu}} \end{pmatrix}. \quad (19)$$

E. Isospin-breaking corrections

The isospin-breaking correction $\delta_{\text{SU}(2)}^{K\ell}$ is defined through the deviation of $f_+^{K\pi}(0)$ from $f_+^{K^0\pi^-}(0)$ (after scaling out the isospin factor C_K),

$$\delta_{\text{SU}(2)}^{K\ell} = \left(\frac{f_+^{K\pi}(0)}{f_+^{K^0\pi^-}(0)} \right)^2 - 1, \quad (20)$$

so it resides in the K^+ channel only by construction. Upon neglecting small EM contributions, it is given by [45]

$$\delta_{\text{SU}(2)}^{K^+\ell} = \frac{3}{2} \frac{1}{Q^2} \left[\frac{\hat{M}_K^2}{\hat{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left(1 + \frac{m_s}{\hat{m}} \right) \right], \quad (21)$$

where $\hat{M}_{K,\pi}$ are the meson masses in the isospin limit, $Q^2 = (m_s^2 - \hat{m}^2)/(m_d^2 - m_u^2)$, and $\chi_{p^4} \simeq 0.219$ is calculable in chiral perturbation theory (ChPT) [88]. The pure quantum chromodynamics (QCD) mass parameters can only be obtained through lattice simulations. Here, we quote the most precise results of Q and $m_s/\hat{m}$ from the latest web update [89] of the Flavor Lattice Averaging Group (FLAG) review [53],

$$Q = 23.3(5), \quad m_s/\hat{m} = 27.42(12), \\ N_f = 2 + 1, \quad \text{Refs. [90–94]}, \\ Q = 24.0(8), \quad m_s/\hat{m} = 27.23(10), \\ N_f = 2 + 1 + 1, \quad \text{Refs. [95–98]}. \quad (22)$$

Since Q is by far the main contributor of the uncertainty in $\delta_{\text{SU}(2)}^{K^+\ell}$, we choose the more precise data set from $N_f = 2 + 1$ for numerical applications. Meanwhile, Ref. [53] did not provide the explicit values of $\hat{M}_{K,\pi}$, so we quote them from the 2017 FLAG review [99]: $\hat{M}_\pi = 134.8(3)$ MeV, $\hat{M}_K = 494.2(3)$ MeV. Putting pieces together, we have

$$\delta_{\text{SU}(2)}^{K^+\ell} = 0.0457(20). \quad (23)$$

The contribution of $\delta_{\text{SU}(2)}$ to the covariance matrix of A is given by

$$\text{Cov}(A)_{\delta_{\text{SU}(2)}} = \left(\frac{\partial A}{\partial \delta_{\text{SU}(2)}} \right) \cdot \text{Cov}(\delta_{\text{SU}(2)}) \cdot \left(\frac{\partial A}{\partial \delta_{\text{SU}(2)}} \right)^T, \quad (24)$$

where

$$\left(\frac{\partial A}{\partial \delta_{\text{SU}(2)}}\right) = \begin{pmatrix} 0 & 0 & \frac{\partial A_{K^+e}}{\partial \delta_{\text{SU}(2)}^{K^+e}} & 0 & 0 & \frac{\partial A_{K^+\mu}}{\partial \delta_{\text{SU}(2)}^{K^+\mu}} \end{pmatrix}^T. \quad (25)$$

It should be pointed out that the parameter Q can also be obtained phenomenologically. For instance, Ref. [100] obtained $Q = 22.1(7)$ from $\eta \rightarrow 3\pi$ decay, which is marginally discrepant with the FLAG average based on lattice calculations. We notice that different versions of FlaviaNet updates in the past few years had adopted different choices for their quark mass parameters. Reference [54] took the parameter Q from lattice, while Refs. [52,55] adopted the phenomenological value $Q = 22.1(7)$ (together with a slightly different ChPT parameter $\chi_{p^4} \simeq 0.252$, of which the origin was not clearly explained), which returned a somewhat larger isospin-breaking correction $\delta_{\text{SU}(2)}^{K^+e} = 0.0522(34)$.

F. Long-distance EM corrections

The last theory input is the long-distance EM correction. It was taken in the FlaviaNet review and its updates from the ChPT calculation at $\mathcal{O}(e^2 p^2)$ [101], with a theory uncertainty of the order of 10^{-3} . However, a novel framework based on Sirlin's representation of RC [102] was recently pioneered [103,104]. With this framework and new lattice QCD inputs of the meson γW -box diagrams [49,57], δ_{EM}^{Ke} were reevaluated with a significant increase in precision level reaching 10^{-4} [50,51]. A similar update on $\delta_{\text{EM}}^{K\mu}$ is not yet available due to the more complicated error analysis but will be carried out in the near future. Meanwhile, an important cross-check would be to compute the full $K_{\ell 3}$ RC (both the virtual and real corrections). This can be based on the existing technique that was proven successful in the study of the $K_{\mu 2}/\pi_{\mu 2}$ RCs [105], although its generalization to $K_{\ell 3}$ is expected to be more challenging and could take up to a decade to reach 10^{-3} precision [106].

In this paper, we choose to take the results of $\delta_{\text{EM}}^{K\mu}$ from ChPT, and that of δ_{EM}^{Ke} from the new calculation. Since these two evaluations are based on very different starting points, it is only natural to assume that they are uncorrelated. In the $K\mu$ channels we have [101]

$$\begin{aligned} \delta_{\text{EM}}^{K\mu} &= (\delta_{\text{EM}}^{K^0\mu} \quad \delta_{\text{EM}}^{K^+\mu})^T \\ &= (1.40(22) \times 10^{-2} \quad 0.016(250) \times 10^{-2})^T, \end{aligned} \quad (26)$$

with the correlation matrix

$$\text{Corr}(\delta_{\text{EM}}^{K\mu}) = \begin{pmatrix} 1 & 0.081 \\ & 1 \end{pmatrix}. \quad (27)$$

Meanwhile, in the Ke channels we have

$$\delta_{\text{EM}}^{Ke} = (\delta_{\text{EM}}^{K^0e} \quad \delta_{\text{EM}}^{K^+e})^T, \quad (28)$$

where [50,51]

$$\begin{aligned} \delta_{\text{EM}}^{K^0e} &= 1.16(2)_{\text{sg}}(1)_{\text{lat}}(1)_{\text{NF}}(2)_{e^2 p^4} \times 10^{-2} \\ &= 1.16(3) \times 10^{-2}, \\ \delta_{\text{EM}}^{K^+e} &= 0.21(2)_{\text{sg}}(1)_{\langle t_K^2 \rangle}(1)_{\text{lat}}(4)_{\text{NF}}(1)_{e^2 p^4} \times 10^{-2} \\ &= 0.21(5) \times 10^{-2}. \end{aligned} \quad (29)$$

The correlation matrix of δ_{EM}^{Ke} was not given in Refs. [50,51] and is derived here for the first time.

First of all, we realize that most of the uncertainties in Eq. (29) were estimated through simple power-counting arguments on top of the central values in each respective channel, so the most natural choice is to take them as independent since assuming any correlation would be equally arbitrary. However, there is a piece that has well-defined correlations, namely, the lattice calculation of the meson axial γW -box diagrams, which enters δ_{EM}^{Ke} effectively through the low-energy constants (LECs) in ChPT [104],

$$\begin{aligned} \delta_{\text{EM}}^{K^0e} &= 2e^2 \left[\frac{4}{3}X_1 - \frac{1}{2}\bar{X}_6^{\text{phys}} \right] + \dots, \\ \delta_{\text{EM}}^{K^+e} &= 2e^2 \left[2 \left(-\frac{2}{3}X_1 - \frac{1}{2}\bar{X}_6^{\text{phys}} \right) - \left(\frac{4}{3}X_1 - \frac{1}{2}\bar{X}_6^{\text{phys}} \right) \right] + \dots. \end{aligned} \quad (30)$$

These two combinations of LECs were pinned down by two independent lattice calculations of the axial γW -box diagrams [49,57],

$$\begin{aligned} \square_{\gamma W}^{\pi} &: -\frac{2}{3}X_1 - \frac{1}{2}\bar{X}_6^{\text{phys}} = -7.0(3) \times 10^{-3}, \\ \square_{\gamma W}^K &: \frac{4}{3}X_1 - \frac{1}{2}\bar{X}_6^{\text{phys}} = -11.3(5) \times 10^{-3}. \end{aligned} \quad (31)$$

Therefore, the variations of δ_{EM}^{Ke} due to the lattice uncertainties are given by

$$\delta(\delta_{\text{EM}}^{K^0e}) = 2e^2 \sigma_{K,\text{lat}}, \quad \delta(\delta_{\text{EM}}^{K^+e}) = 2e^2 [2\sigma_{\pi,\text{lat}} - \sigma_{K,\text{lat}}], \quad (32)$$

where $\sigma_{K,\text{lat}} = 0.5 \times 10^{-3}$ and $\sigma_{\pi,\text{lat}} = 0.3 \times 10^{-3}$. The two expressions above depend on a common quantity $\sigma_{K,\text{lat}}$, which gives a nonzero correlation,

$$\langle \delta_{\text{EM}}^{K^0e} \delta_{\text{EM}}^{K^+e} \rangle - \langle \delta_{\text{EM}}^{K^0e} \rangle \langle \delta_{\text{EM}}^{K^+e} \rangle = -4e^4 \sigma_{K,\text{lat}}^2. \quad (33)$$

As a consequence, the correlation matrix reads

$$\text{Corr}(\delta_{\text{EM}}^{Ke}) = \begin{pmatrix} 1 & \Delta \\ & 1 \end{pmatrix}, \quad (34)$$

where

TABLE I. Individual values and weighted average of $|V_{us}f_+^K(0)|$, with independent uncertainties displayed separately. Notice that the uncertainties from the BR and Γ_K are generally correlated, so we display only their combined uncertainty as “exp”.

	$ V_{us}f_+^K(0) $	Correlation matrix						
$K_L e$	$0.21617(46)_{\text{exp}}(10)_{I_K(3)}\delta_{\text{EM}}$	1	0.021	0.025	0.519	0.004	0.017	
$K_S e$	$0.21530(122)_{\text{exp}}(10)_{I_K(3)}\delta_{\text{EM}}$		1	0.009	0.012	0.000	0.006	
$K^+ e$	$0.21714(88)_{\text{exp}}(10)_{I_K(21)}\delta_{\text{SU}(2)}(5)\delta_{\text{EM}}$			1	0.016	0.002	0.871	
$K_L \mu$	$0.21664(50)_{\text{exp}}(16)_{I_K(24)}\delta_{\text{EM}}$				1	0.029	0.047	
$K_S \mu$	$0.21265(466)_{\text{exp}}(16)_{I_K(23)}\delta_{\text{EM}}$					1	0.006	
$K^+ \mu$	$0.21703(108)_{\text{exp}}(16)_{I_K(21)}\delta_{\text{SU}(2)}(26)\delta_{\text{EM}}$						1	
Average: Ke	$0.21626(40)_{K(3)}\text{HO}$							
Average: $K\mu$	$0.21667(52)_{K(3)}\text{HO}$							
Average: tot	$0.21635(39)_{K(3)}\text{HO}$							

$$\Delta = -\frac{4e^4\sigma_{K,\text{lat}}^2}{\sigma_{\delta_{\text{EM}}^{K^0e}}\sigma_{\delta_{\text{EM}}^{K^+e}}}, \quad (35)$$

with $\sigma_{\delta_{\text{EM}}^{K^0e}} \simeq 0.03 \times 10^{-2}$, $\sigma_{\delta_{\text{EM}}^{K^+e}} \simeq 0.05 \times 10^{-2}$ as given above.

We may now combine all the independent long-distance EM corrections as

$$\delta_{\text{EM}} = (\delta_{\text{EM}}^{K^0e} \quad \delta_{\text{EM}}^{K^+e} \quad \delta_{\text{EM}}^{K^0\mu} \quad \delta_{\text{EM}}^{K^+\mu})^T. \quad (36)$$

Its correlation matrix is given by the following block-diagonal matrix:

$$\text{Corr}(\delta_{\text{EM}}) = \begin{pmatrix} \text{Corr}(\delta_{\text{EM}}^{K^0e}) & 0 \\ 0 & \text{Corr}(\delta_{\text{EM}}^{K^+\mu}) \end{pmatrix}. \quad (37)$$

The contribution of δ_{EM} to the covariance matrix of A is given by

$$\text{Cov}(A)_{\delta_{\text{EM}}} = \left(\frac{\partial A}{\partial \delta_{\text{EM}}} \right) \cdot \text{Cov}(\delta_{\text{EM}}) \cdot \left(\frac{\partial A}{\partial \delta_{\text{EM}}} \right)^T, \quad (38)$$

where

$$\left(\frac{\partial A}{\partial \delta_{\text{EM}}} \right) = \begin{pmatrix} \frac{\partial A_{K_L e}}{\partial \delta_{\text{EM}}^{K^0e}} & 0 & 0 & 0 \\ \frac{\partial A_{K_S e}}{\partial \delta_{\text{EM}}^{K^0e}} & 0 & 0 & 0 \\ 0 & \frac{\partial A_{K^+ e}}{\partial \delta_{\text{EM}}^{K^+e}} & 0 & 0 \\ 0 & 0 & \frac{\partial A_{K_L \mu}}{\partial \delta_{\text{EM}}^{K^0\mu}} & 0 \\ 0 & 0 & \frac{\partial A_{K_S \mu}}{\partial \delta_{\text{EM}}^{K^0\mu}} & 0 \\ 0 & 0 & 0 & \frac{\partial A_{K^+ \mu}}{\partial \delta_{\text{EM}}^{K^+\mu}} \end{pmatrix}. \quad (39)$$

G. Final result

With the above, we may obtain the values of $|V_{us}f_+^K(0)|$ from each channel, which are summarized in the left panel of Table I. Making use of the total covariance matrix of A ,

$$\text{Cov}(A) = \text{Cov}(A)_{K_L^{\text{exp}}} + \text{Cov}(A)_{K_S^{\text{exp}}} + \text{Cov}(A)_{K_+^{\text{exp}}} + \text{Cov}(A)_{I_K} + \text{Cov}(A)_{\delta_{\text{SU}(2)}} + \text{Cov}(A)_{\delta_{\text{EM}}}, \quad (40)$$

from which the total correlation matrix in the right panel of Table I can be calculated; the weighted average of $|V_{us}f_+^K(0)|$ can then be obtained using Eq. (A9). Given the different theory statuses of K_{e3} and $K_{\mu 3}$, we present simultaneously the average values of $|V_{us}f_+^K(0)|$ by weighting over the Ke channels, the $K\mu$ channels, and both. Notice that the uncertainty from S_{EW} is common to all channels and does not enter the weighting process. Therefore, we choose to display it only in the weighted averages, i.e., the last three rows in Table I. We find that the Ke and $K\mu$ averages agree well with each other within uncertainties. Finally, the 2020 website update of the FLAG review quoted [89]

$$\begin{aligned} |f_+^K(0)| &= 0.9698(17), \quad N_f = 2 + 1 + 1, \quad \text{Refs. [107, 108]} \\ |f_+^K(0)| &= 0.9677(27), \quad N_f = 2 + 1, \quad \text{Refs. [109, 110]}, \\ |f_+^K(0)| &= 0.9560(57)(62), \quad N_f = 2, \quad \text{Ref. [111]}. \end{aligned} \quad (41)$$

We choose the most precise value from $N_f = 2 + 1 + 1$ for numerical applications. With that we obtain

$$|V_{us}|_{K_{e3}} = 0.22309(40)_{K(3)}\text{lat}(3)_{\text{HO}}. \quad (42)$$

Let us discuss the results above. First, both the central value and the total uncertainty in the weighted average of $|V_{us}f_+^K(0)|$ experience no significant change compared to those in previous reviews (e.g., 0.21654(41) in Ref. [54]), but not the composition of uncertainties in each channel.

In our latest analysis, the combined experimental uncertainty from the kaon lifetime and BRs are by far the dominant source of uncertainty in all channels. This is quite different from a few years ago, wherein some channels (e.g., K^+e) the theory and experimental uncertainties are comparable. Such changes are mainly due to the improved theory precision in $\delta_{\text{SU}(2)}^{K\ell}$ and δ_{EM}^{Ke} .

We may also review the status of the top-row CKM unitarity. The best extraction of $|V_{ud}|$ comes from superallowed $0^+ \rightarrow 0^+$ beta decays, but its precise value depends on the theory inputs of the single-nucleon RC and nuclear structure corrections. In particular, it was recently pointed out that several potentially large new nuclear corrections (NNCs) that reside in the nuclear γW -box diagrams were missed in the existing nuclear structure calculations [12,112]; their true sizes are poorly understood and are at present only roughly estimated based on a simple Fermi gas model. After alerting the readers about the possible (small) quantitative difference due to different choices of theory inputs, let us quote, just for this work, the result from the latest review by Hardy and Towner [19],

$$|V_{ud}|_{0^+} = 0.97373(11)_{\text{exp}}(9)_{\text{RC}}(27)_{\text{NS}}, \quad (43)$$

where “exp” is the combined uncertainty from experiment and the so-called “outer” correction, “RC” the theory uncertainty from the single-nucleon (inner) RC, and “NS” the nuclear structure uncertainty that originates primarily from the NNCs. Combining Eqs. (42) and (43) gives

$$\begin{aligned} |V_{ud}|_{0^+}^2 + |V_{us}|_{K_{e3}}^2 - 1 \\ = -0.0021(2)_{V_{ud},\text{exp}}(2)_{V_{ud},\text{RC}}(5)_{V_{ud},\text{NS}}(2)_{V_{us},K}(2)_{V_{us},\text{lat}}, \end{aligned} \quad (44)$$

while the SM prediction (after neglecting the small $|V_{ub}|^2$) is 0. The above indicates an apparent anomaly in the top-row CKM unitarity with the significance level of 3.2σ , which could increase to as much as 5.6σ if we imagine that the NS uncertainty was significantly reduced while the central value of $|V_{ud}|$ remained unchanged. This provides a strong motivation for nuclear theorists to perform *ab initio* calculations of the NS correction in superallowed beta decays to reduce its theory uncertainty.

III. $|V_{us}f_+^K(0)|/|V_{ud}f_+^\pi(0)|$ AND $|V_{us}/V_{ud}|$ FROM K/π SEMILEPTONIC DECAYS

In addition to its contribution to an apparent violation of the top-row CKM unitarity, Eq. (42) also shows a direct disagreement at the level $\sim 2.8\sigma$ with the same quantity extracted from $K_{\mu 2}$ decay: $|V_{us}|_{K_{\mu 2}} = 0.2252(5)$ [1]. The latter is obtained from the ratio $R_A = \Gamma(K_{\mu 2})/\Gamma(\pi_{\mu 2})$,

which gives the value of $|V_{us}f_{K^+}|/|V_{ud}f_{\pi^+}|$, with f_{K^+} and f_{π^+} the decay constants of the charged kaon and pion [113],

$$\left| \frac{V_{us}f_{K^+}}{V_{ud}f_{\pi^+}} \right| = 0.23871(20)_{\text{RC}} \times R_A^{1/2}, \quad (45)$$

where the theory uncertainty of 0.084% at the right-hand side originates from residual long-distance RCs that do not cancel in the ratio. This residual RC has been calculated using both ChPT [114] and lattice QCD [105] with comparable sizes of theory uncertainties. The two calculations agree well with each other, showing that the theory error in this input is under good control. Following PDG, we utilize the ChPT input in Eq. (45) for illustration (throughout this work, we add nothing new in the R_A analysis; everything is the same as in the 2021 online version of the PDG review).

The above discrepancy may indicate the presence of BSM effects or possible unidentified SM corrections that are not reflected in the existing error estimation, such as a smaller value for $|f_+^K(0)|$ outside the range of the quoted $N_f = 2 + 1 + 1$ lattice QCD result. To further explore these possibilities, in particular, the latter, Ref. [56] suggested to study a new ratio $R_V = \Gamma(K_{e3})/\Gamma(\pi_{e3})$ (which takes different values in different $K\ell$ channels), where the V denotes the fact that such decays are due to weak vector current interactions. Like R_A , it results from a ratio of weak interaction meson decays (induced by vector rather than axial-vector interactions) for which theoretical uncertainties partially cancel. A comparison of R_V and R_A can, in principle, unveil the influence of BSM physics.

We first recall the SM prediction of the π_{e3} decay width,

$$\Gamma_{\pi_{e3}} = \frac{G_F^2 |V_{ud}|^2 M_{\pi^+}^5 |f_+^\pi(0)|^2}{64\pi^3} (1 + \text{RC}_\pi) I_\pi. \quad (46)$$

The left-hand side is calculated from the experimental measurement of the charged pion lifetime and the semi-leptonic decay BR [1],

$$\tau_{\pi^+} = 2.6033(5) \times 10^{-8} \text{ s}, \quad \text{BR}(\pi_{e3}) = 1.038(6) \times 10^{-8}. \quad (47)$$

Notice the slight modification of the BR from the PDG value that took into account the effect of the updated $\text{BR}(\pi_{e2})$ normalization [56]. On the right-hand side, $|f_+^\pi(0)|$ is the $\pi^+ \rightarrow \pi^0$ form factor at zero momentum transfer which equals 1 in the isospin limit (isospin-breaking correction is negligible due to the Behrends-Sirlin-Ademollo-Gatto theorem [115,116]). In this limit, the phase space integral I_π is calculable,

$$I_\pi = \int_{2\sqrt{r_\pi}}^{1+r_\pi-r_e} dz \int_{c(z)-d(z)}^{c(z)+d(z)} dy [4(1-y)(y+z-1) + r_e(4y+3z-3) - 4r_\pi + r_e(r_\pi - r_e)]$$

$$= 7.3764 \times 10^{-8}, \quad (48)$$

where

$$r_\pi = \frac{M_{\pi^0}^2}{M_{\pi^+}^2}, \quad r_e = \frac{m_e^2}{M_{\pi^+}^2},$$

$$c(z) = \frac{(2-z)(1+r_e+r_\pi-z)}{2(1+r_\pi-z)},$$

$$d(z) = \frac{\sqrt{z^2 - 4r_\pi(1+r_\pi-r_e-z)}}{2(1+r_\pi-z)}, \quad (49)$$

in analogy to the well-known $K_{\ell 3}$ phase space formula (see, e.g., Ref. [51]). RC_π is the electroweak RC in the pion semileptonic decay which was recently determined to high precision with lattice QCD: $\text{RC}_\pi = 0.0332(1)_{\gamma W}(3)_{\text{HO}}$ [57]. Combining Eqs. (1) and (46) gives

$$\left| \frac{V_{us}f_+^K(0)}{V_{ud}f_+^\pi(0)} \right|_{K\ell}$$

$$= \sqrt{\frac{3}{C_K^2} \left(\frac{M_{\pi^+}}{M_K} \right)^5 \frac{I_\pi}{I_{K\ell}} \frac{1 + \text{RC}_\pi}{S_{\text{EW}}(1 + \delta_{\text{EM}}^{K\ell} + \delta_{\text{SU}(2)}^{K\ell})}} \times (R_V^{K\ell})^{1/2}, \quad (50)$$

which provides a measure of $|V_{us}f_+^K(0)|/|V_{ud}f_+^\pi(0)|$.

There are several benefits in studying R_V . First, uncertainties from short-distance electroweak RCs (contained in RC_π and S_{EW} , although they are numerically smaller than the other SM theory uncertainties, e.g., those coming from I_K and $\delta_{\text{SU}(2)}$) as well as BSM effects that are common to the numerator and denominator (e.g., those correcting G_F through the muon lifetime) cancel each other in the R_V ratio. This means as follows: should one observe a significant discrepancy between the values of $|V_{us}/V_{ud}|$ obtained from R_V and R_A , then its possible BSM explanations would be more limited than those which could be used to explain the discrepancy between the values of $|V_{us}|$ from $K_{\mu 3}$ and R_A . This makes R_V a useful gauge to search for possibly large nonuniversal systematic effects, especially those from the SM. Second, the recent improvements in SM theory precision, in particular, the electroweak RCs, makes R_V an extremely clean observable from the theory aspect. Consider, for instance, the $K_L e$ channel. Substituting all the SM theory inputs we discussed above, one obtains

$$\left| \frac{V_{us}f_+^K(0)}{V_{ud}f_+^\pi(0)} \right|_{K_L e} = 4.9786(24)_{I_K}(7)_{\text{RC}_K}(2)_{\text{RC}_\pi} \times 10^{-5}$$

$$\times (R_V^{K_L e})^{1/2}. \quad (51)$$

We see that the total theory uncertainty on the right-hand side is only 0.051%, which is already better than R_A . Moreover, the above is only for one channel in $K_{\ell 3}$. A further reduction of uncertainty is achieved once all six channels are weighted over.

Upon substituting the experimental inputs into Eqs. (45) and (50), we find

$$\left| \frac{V_{us}f_{K^+}}{V_{ud}f_{\pi^+}} \right| = 0.27600(29)_{\text{exp}}(23)_{\text{RC}},$$

0.13% precision,

$$\left| \frac{V_{us}f_+^K(0)}{V_{ud}f_+^\pi(0)} \right| = 0.22216(64)_{\text{BR}(\pi_{e3})}(40)_K(2)_{\tau_{\pi^+}}(1)_{\text{RC}_\pi},$$

0.34% precision, (52)

where the first line is from R_A , and the second line is from R_V weighted over all six channels. The precision of the latter is limited primarily by the uncertainty in $\text{BR}(\pi_{e3})$ and secondarily by the $K_{\ell 3}$ experiments. Future improvements of the experimental precision in these areas are therefore urgently needed.

At this point, it is interesting to discuss the relevance with respect to PIONEER, the proposed next-generation experiment for rare pion decays which may take place in PSI or TRIUMF [58,59]. It is originally designed for an improved measurement of the ratio $R_{e/\mu} = \Gamma(\pi^+ \rightarrow e^+ \nu(\gamma))/\Gamma(\pi^+ \rightarrow \mu^+ \nu(\gamma))$ to test lepton universality, but the optimized detector is also ideal for a high-precision measurement of $\text{BR}(\pi_{e3})$. The current best measurement of the latter is from the PIBETA experiment [117], which leads to the following extracted value of $|V_{ud}|$ [57]:

$$|V_{ud}|_{\pi_{e3}} = 0.9740(28)_{\text{exp}}(1)_{\text{th}}. \quad (53)$$

Despite being theoretically clean, it is 10 times less precise than the superallowed beta decay extraction [see Eq. (43)]. To make $|V_{ud}|_{\pi_{e3}}$ competitive requires a 10 times reduction of the experimental uncertainty, which may require 100 times the statistics of the existing measurement and comparable reduction of systematics and backgrounds, a very ambitious long-term goal. However, the introduction of R_V provides a new physical significance to π_{e3} not just in terms of $|V_{ud}|$ but also $|V_{us}/V_{ud}|$. With this, it is most beneficial to plan the next-generation π_{e3} experiment for two stages:

- (i) The first stage would primarily aim to improve the precision of the ratio, $R_{e/\mu}$, (its primary goal) by an order of magnitude. That same phase could be used

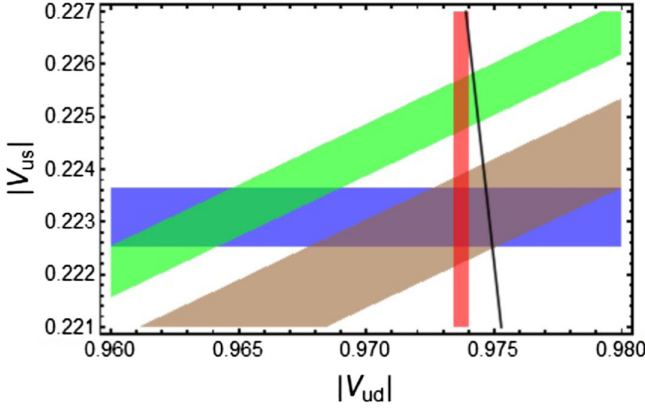


FIG. 1. A combined plot of $|V_{ud}|$ from superallowed beta decays (red band), $|V_{us}|$ from $K_{\ell 3}$ (blue band), $|V_{us}/V_{ud}|$ from R_A (green band), and R_V (brown band), together with the first-row CKM unitarity requirement (black line). Notice that the blue and brown bands are highly correlated as they rely on the same set of inputs from $K_{\ell 3}$, so they should not be taken as independent constraints.

to improve the precision of $\text{BR}(\pi_{e3})$ by a factor of 3 or better compared to the existing PIBETA result.

That would reduce the uncertainty in R_V to a level comparable to R_A , making for an interesting confrontation. If accompanied by future improvement in $K_{\ell 3}$ experiments, R_V could eventually surpass R_A as the primary means to constrain $|V_{us}/V_{ud}|$.

- (ii) In the second stage, an overall improvement of a factor of 10 improvement in the $\text{BR}(\pi_{e3})$ precision is required to compete with superallowed beta decays for precision in extracting V_{ud} . It is, however, much more challenging and is not yet at the achievable level in the present technical design [118].

We close this section by reporting the current extracted values of $|V_{us}/V_{ud}|$ from R_A and R_V respectively, by supplementing Eq. (52) with relevant lattice QCD inputs [89],

$$|f_{K^+}/f_{\pi^+}| = 1.1932(21), \quad N_f = 2 + 1 + 1,$$

$$\text{Refs. [95, 119 – 121]},$$

$$|f_+^K(0)/f_+^\pi(0)| \approx |f_+^K(0)| = 0.9698(17),$$

$$N_f = 2 + 1 + 1, \quad \text{Refs. [107, 108]}. \quad (54)$$

They give

$$|V_{us}/V_{ud}| = \begin{cases} 0.23131(41)_{\text{lat}}(24)_{\text{exp}}(19)_{\text{RC}} & \text{from } R_A \\ 0.22908(66)_{\text{BR}(\pi_{e3})}(41)_K(40)_{\text{lat}}(2)_{\tau_{\pi^+}}(1)_{\text{RC}_\pi} & \text{from } R_V \end{cases}. \quad (55)$$

The difference between the two determinations is at the level of 2.2σ . All the determinations of $|V_{us}|$, $|V_{ud}|$, and their ratio quoted in this paper are summarized in Fig. 1, from which the mutual disagreements between different determinations and the deviations from the first-row CKM unitarity requirement are clearly shown.

IV. CONCLUSIONS

This work updates the values of $|V_{us}|$ and $|V_{us}/V_{ud}|$ determined from kaon and pion semileptonic decays using the most recent inputs from theory and experiment. The uncertainties in these quantities have been experiment and lattice dominated, which is even more the case in the recent years due to the more precise SM electroweak theory inputs. Their values along with V_{ud} from superallowed beta decays correspond to $2\text{--}3\sigma$ deviations from CKM unitarity and related axial current induced weak decays. Those differences may provide hints of BSM physics or deficiencies in SM theory or experiment. Such anomalies provide a strong motivation for future improvements of the experimental precision of the π_{e3} BR as well as the kaon lifetimes and $K_{\ell 3}$ BRs.

The experiment-dominated uncertainties do not imply that future improvements from the theory side are not important. It is quite the opposite; the aforementioned

anomalies require us to carefully reexamine all the SM theory inputs in order to ensure that no unexpected large systematic errors exist. This was recently done for the long-distance EM corrections to the K_{e3} decay rate, and no large corrections were found. Other inputs, such as the lattice calculations of $f_+^K(0)$ and f_{K^+}/f_{π^+} , should be cross-checked with the same level of rigor. There are several other theory works that remain to be done for completeness and internal consistency. For instance, the reevaluation of the EM corrections should be generalized to the $K\mu$ channels, and the fitting of the $K\pi$ form factors should, in principle, also be updated to account for the modified EM corrections to the $K_{\ell 3}$ Dalitz plot.

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APPENDIX: MATHEMATICAL TOOLS

In this Appendix, we review all the mathematical tools needed in this work.

1. Covariance matrix and correlation matrix

For a set of variables $X = (x_1, x_2, \dots, x_n)^T$, we define the symmetric covariance matrix $\text{Cov}(X)$ as

$$\text{Cov}(X)_{ij} = \langle x_i x_j \rangle - \langle x_i \rangle \langle x_j \rangle. \quad (\text{A1})$$

In particular, the diagonal terms give the variance of x_i ,

$$\sigma_i^2 = \text{Cov}(X)_{ii}. \quad (\text{A2})$$

We also define the symmetric correlation matrix $\text{Corr}(X)$ as

$$\text{Corr}(X)_{ij} = \frac{\langle x_i x_j \rangle - \langle x_i \rangle \langle x_j \rangle}{\sigma_i \sigma_j}. \quad (\text{A3})$$

Its diagonal elements are always 1, while the off-diagonal elements range between -1 and 1 . Its relation to the covariance matrix is given by

$$\text{Cov}(X) = \sigma_X \cdot \text{Corr}(X) \cdot \sigma_X, \quad (\text{A4})$$

where $\sigma_X \equiv \text{diag}(\sigma_1, \sigma_2, \dots, \sigma_n)$.

2. Propagation of the covariance matrix

For a set of variables $F = (f_1, f_2, \dots, f_m)^T$ that are functions of X [i.e., $f_i = f_i(x_1, \dots, x_n)$], given the covariance matrix of X , we can immediately obtain the covariance matrix of F as

$$\text{Cov}(F) = \left(\frac{\partial F}{\partial X} \right) \cdot \text{Cov}(X) \cdot \left(\frac{\partial F}{\partial X} \right)^T, \quad (\text{A5})$$

where $\partial F / \partial X$ is a $m \times n$ matrix, with matrix elements

$$\left(\frac{\partial F}{\partial X} \right)_{ij} = \frac{\partial f_i}{\partial x_j}. \quad (\text{A6})$$

If $F = (f_1, f_2, \dots, f_m)^T$ depends on two independent sets of variables $X = (x_1, x_2, \dots, x_n)^T$ and $Y = (y_1, y_2, \dots, y_{n'})^T$ with their respective covariance matrices given, then the covariance matrix of F is simply the sum of the two contributions,

$$\begin{aligned} \text{Cov}(F) &= \left(\frac{\partial F}{\partial X} \right) \cdot \text{Cov}(X) \cdot \left(\frac{\partial F}{\partial X} \right)^T \\ &+ \left(\frac{\partial F}{\partial Y} \right) \cdot \text{Cov}(Y) \cdot \left(\frac{\partial F}{\partial Y} \right)^T. \end{aligned} \quad (\text{A7})$$

This is also generalizable if F is a function of more than two independent sets of variables.

For definiteness, throughout this work, we always calculate partial derivatives numerically as

$$\frac{\partial f}{\partial a} \approx \frac{f(a + \delta a) - f(a - \delta a)}{2\delta a}. \quad (\text{A8})$$

3. Weighted average

If $X = (x_1, x_2, \dots, x_n)^T$ has a covariance matrix $\text{Cov}(X)$, then the weighted average between $x_1, \dots, x_n$ is given by

$$\bar{x} = \sigma_X^2 (J^T W X), \quad (\text{A9})$$

where the variance of $\bar{x}$ is given by

$$\sigma_{\bar{x}}^2 = (J^T W J)^{-1}. \quad (\text{A10})$$

Here, we have defined $W = [\text{Cov}(X)]^{-1}$ and $J = (1, \dots, 1)^T$ (length = n).

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