



Resonances in heavy meson–heavy baryon coupled-channel interactions

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Abstract The interactions of $\bar{D}^{(*)}\Lambda_c - \bar{D}^{(*)}\Sigma_c^{(*)}$ are studied within the framework of a dynamical coupled-channel approach. A series of bound states and resonances with different spin and parity are dynamically generated in the hidden charm sector. Four S -wave bound states are found in the mass range of 4.3–4.5 GeV, close to the pentaquark states observed by LHCb. Two of the states have a spin parity of $J^P = 1/2^-$ and the other two have $J^P = 3/2^-$. In addition, several resonances with different spin and parity in higher partial waves are predicted.

1 Introduction

In 2015, two pentaquark-like resonances were observed by the LHCb Collaboration in the $J/\psi p$ invariant mass spectrum of the $\Lambda_b^0 \rightarrow J/\psi p K^-$ decay [1]. The results of these P_c states were then updated in 2019 with three clear narrow structures reported in Ref. [2]. Recently, another new pentaquark state $P_c(4337)$ has been announced with a significance between 3.1 and 3.7σ by the LHCb Collaboration [3]. The existence of these pentaquark-like states with hidden charm had been predicted before the experimental findings in Refs. [4–11], and they attracted much attention after the observation. Among the theoretical works, a popular explanation is the anticharmed meson-charmed baryon molecular

picture, where these exotic states are $\bar{D}^{(*)}\Sigma_c^{(*)}$ bound states [12–38].

In Ref. [33], the authors applied the Jülich–Bonn dynamical coupled-channel (JüBo DCC) framework [39] to the hidden-charm sector to investigate the resonance spectrum in the $\bar{D}\Lambda_c - \bar{D}\Sigma_c$ system. The JüBo approach was originally developed to extract the light baryon resonance spectrum from πN scattering data and has been successfully applied in the simultaneous analysis of pion- and photon-induced reactions [40–42]. Recently, it was extended to include also electroproduction reactions [43,44]. The scattering amplitude is obtained as the solution of a Lippmann–Schwinger equation respecting theoretical constraints on the S -matrix like unitarity and analyticity. In the exploratory study of Ref. [33], only two channels $\bar{D}\Lambda_c$ and $\bar{D}\Sigma_c$ with vector meson exchange in the t -channel were included and one pole in each partial wave up to $J^P = 5/2^+$ and $5/2^-$ was dynamically generated. For $J^P = 1/2^-$, there exist a very narrow pole around 4296 MeV, which can be considered as a $\bar{D}\Sigma_c$ bound state.

Here, we extend the work of Ref. [33] to a more complete coupled-channel calculation, including $\bar{D}\Lambda_c$, $\bar{D}\Sigma_c$, $\bar{D}^*\Lambda_c$, $\bar{D}^*\Sigma_c$, $\bar{D}\Sigma_c^*$ channels, where t -channel pseudoscalar and vector meson exchange and u -channel doubly charmed baryon exchange are taken into account. No “genuine” s -channel poles are included. Possible dynamically generated poles in different partial waves in the energy range around 4 GeV are studied. Together with their couplings to the different channels, we expect to obtain a more comprehensive picture of the hidden-charm resonance spectrum.

This work is organized as follows. In Sect. 2, we present the theoretical framework of our calculation. In Sect. 3, the numerical results for the $\bar{D}\Lambda_c$, $\bar{D}\Sigma_c$, $\bar{D}^*\Lambda_c$, $\bar{D}^*\Sigma_c$, $\bar{D}\Sigma_c^*$

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Table 1 Angular momentum structure of the coupled channels in isospin $I = \frac{1}{2}$ up to $J = \frac{9}{2}$

μ	$J^P =$	$\frac{1}{2}^-$	$\frac{1}{2}^+$	$\frac{3}{2}^+$	$\frac{3}{2}^-$	$\frac{5}{2}^-$	$\frac{5}{2}^+$	$\frac{7}{2}^+$	$\frac{7}{2}^-$	$\frac{9}{2}^-$	$\frac{9}{2}^+$
1	$\bar{D}\Lambda_c$	S_{11}	P_{11}	P_{13}	D_{13}	D_{15}	F_{15}	F_{17}	G_{17}	G_{19}	H_{19}
2	$\bar{D}\Sigma_c$	S_{11}	P_{11}	P_{13}	D_{13}	D_{15}	F_{15}	F_{17}	G_{17}	G_{19}	H_{19}
3	$\bar{D}^*\Lambda_c(S = 1/2)$	S_{11}	P_{11}	P_{13}	D_{13}	D_{15}	F_{15}	F_{17}	G_{17}	G_{19}	H_{19}
4	$\bar{D}^*\Lambda_c(S = 3/2, J - L = 1/2)$	–	P_{11}	P_{13}	D_{13}	D_{15}	F_{15}	F_{17}	G_{17}	G_{19}	H_{19}
5	$\bar{D}^*\Lambda_c(S = 3/2, J - L = 3/2)$	D_{11}	–	F_{13}	S_{13}	G_{15}	P_{15}	H_{17}	D_{17}	I_{19}	F_{19}
6	$\bar{D}^*\Sigma_c(S = 1/2)$	S_{11}	P_{11}	P_{13}	D_{13}	D_{15}	F_{15}	F_{17}	G_{17}	G_{19}	H_{19}
7	$\bar{D}^*\Sigma_c(S = 3/2, J - L = 1/2)$	–	P_{11}	P_{13}	D_{13}	D_{15}	F_{15}	F_{17}	G_{17}	G_{19}	H_{19}
8	$\bar{D}^*\Sigma_c(S = 3/2, J - L = 3/2)$	D_{11}	–	F_{13}	S_{13}	G_{15}	P_{15}	H_{17}	D_{17}	I_{19}	F_{19}
9	$\bar{D}\Sigma_c^*(J - L = 1/2)$	–	P_{11}	P_{13}	D_{13}	D_{15}	F_{15}	F_{17}	G_{17}	G_{19}	H_{19}
10	$\bar{D}\Sigma_c^*(J - L = 3/2)$	D_{11}	–	F_{13}	S_{13}	G_{15}	P_{15}	H_{17}	D_{17}	I_{19}	F_{19}

interactions and relevant discussions are presented. We end with a brief summary in Sect. 4. Many technicalities are relegated to the Appendix.

2 Formalism

In the JüBo model, the scattering equation describing the coupled-channel interactions for a given partial wave takes the form

$$T_{\mu\nu}(p'', p', z) = V_{\mu\nu}(p'', p', z) + \sum_{\kappa} \int_0^{\infty} dp p^2 V_{\mu\kappa}(p'', p, z) \times G_{\kappa}(p, z) T_{\kappa\nu}(p, p', z), \quad (1)$$

where z is the scattering energy in the center-of-mass system, $p'' = |\mathbf{p}''|$, $p' = |\mathbf{p}'|$ and $p = |\mathbf{p}|$ represent the outgoing, in-coming, and intermediate three-momenta, respectively, which may be on- or off-shell. The subscripts μ, ν, κ are channel indices, and the integral term contains a sum over all intermediate possible quantum numbers and channels involved in the model. The propagator $G_{\kappa}(p, z)$ is given by

$$G_{\kappa}(p, z) = \frac{1}{z - E_a(p) - E_b(p) + i\varepsilon}, \quad (2)$$

where $E_a(p) = \sqrt{m_a^2 + p^2}$ and $E_b(p) = \sqrt{m_b^2 + p^2}$ are the on-shell energies of the intermediate particles a and b in the channel κ . In Eq. (1) left-hand cuts and the correct structure of complex branch points are implemented, which allows for a well-defined determination of resonances as poles in the complex energy plane on the second Riemann sheet, see Ref. [45] for a detailed discussion.

It should be pointed out that there are different possibilities for the $\bar{D}^*\Lambda_c$, $\bar{D}^*\Sigma_c$ and $\bar{D}\Sigma_c^*$ channels to couple to a given spin-parity J^P . Different coupling possibilities for these channels are denoted as different channels μ . In Table 1, we present the complete coupling scheme utilized in this work, where S is the total spin and L is the orbital angular momentum. The notation of the quantum numbers is the common spectroscopic notation, $L_{2I,2J}$, with I denoting the isospin.

The potential V iterated in Eq.(1) is constructed from the effective interactions based on the Lagrangians of Wess and Zumino [46,47], using time-ordered perturbation theory (TOPT). The involved effective Lagrangians are

$$\begin{aligned} \mathcal{L}_{PPV} &= i\sqrt{2}g_{PPV}(P\partial^\mu P - \partial^\mu PP)V_\mu, \\ \mathcal{L}_{VVP} &= \frac{g_{VVP}}{m_V}\epsilon_{\mu\nu\alpha\beta}\partial^\mu V^\nu\partial^\alpha V^\beta P, \\ \mathcal{L}_{VVV} &= ig_{VVV}\langle V^\mu[V^\nu, \partial_\mu V_\nu] \rangle, \\ \mathcal{L}_{BBP} &= \frac{g_{BBP}}{m_P}\bar{B}\gamma^\mu\gamma^5\partial_\mu PB, \\ \mathcal{L}_{BBV} &= -g_{BBV}\bar{B}(\gamma^\mu - \frac{\kappa}{2m_B}\sigma^{\mu\nu}\partial_\nu)V_\mu B, \\ \mathcal{L}_{BDP} &= -\frac{g_{BDP}}{m_P}(\bar{B}\partial^\mu PD_\mu + \bar{D}_\mu\partial^\mu PB), \\ \mathcal{L}_{BDV} &= i\frac{g_{BDV}}{m_V}[\bar{B}\gamma_\mu\gamma^5D_\nu(\partial^\mu V^\nu - \partial^\nu V^\mu) \\ &\quad + \bar{D}_\mu\gamma_\nu\gamma^5B(\partial^\mu V^\nu - \partial^\nu V^\mu)], \\ \mathcal{L}_{DDV} &= g_{DDV}\bar{D}^\tau(\gamma^\mu - \frac{\kappa}{2m_D}\sigma^{\mu\nu}\partial_\nu)V_\mu D_\tau, \end{aligned} \quad (3)$$

where P, V, B and D denote the pseudoscalar, vector meson, and baryon octet, in order, and the decuplet baryons. We have extended the original formulation [39] for three flavors to four, but of course this flavor SU(4) symmetry is strongly broken by the meson and baryon masses. We apply SU(3) and SU(4) flavor symmetry [48,49] to relate the specific coupling constants. The expressions of the amplitudes, detailed

relations of the couplings and other theoretical formulae are given in Appendix A.

3 Results and discussions

In this section, we first discuss some generalities pertinent to the considered processes. As explained in Appendix A, each exchange process is multiplied by a form factor which contains a cut-off value. In general in the JüBo model, those cut-offs are fitted to experimental data and the extracted pole positions depend on the values of those internal parameters. In the present study, however, we use the cut-off to postdict the position of the LHCb pentaquarks without a direct fit to the published line shapes, predicting additional states at the same time. The exact values of the cut-offs used can be found in Appendix A. It should be noted that we do not attribute any physical meaning to those model-dependent parameters. A fit to the experimental data of the P_c states as well as an uncertainty analysis is beyond the scope of this work and postponed to the future.

We find that the potentials of $\bar{D}\Sigma_c \rightarrow \bar{D}\Sigma_c$, $\bar{D}^*\Sigma_c \rightarrow \bar{D}^*\Sigma_c$ and $\bar{D}\Sigma_c^* \rightarrow \bar{D}\Sigma_c^*$ are attractive, and S -wave bound states can be formed in these channels. To confirm this, we first do the single channel calculation for all the considered channels as a test. The absolute value of the T -matrix for different channels in the single channel case, such as $\bar{D}\Sigma_c \rightarrow \bar{D}\Sigma_c$, in S -wave is shown in Fig. 1. A peak structure around 4310 MeV is located just below the $\bar{D}\Sigma_c$ threshold. Similarly, when performing a single channel calculation, we find a peak for the $\bar{D}\Sigma_c^*$ channel in the S_{13} partial wave. For the $\bar{D}^*\Sigma_c$ channel, there exist two peaks in the S_{11} and S_{13} partial waves, respectively. Contrary to this, the interactions of $\bar{D}\Lambda_c \rightarrow \bar{D}\Lambda_c$ and $\bar{D}\Lambda_c^* \rightarrow \bar{D}\Lambda_c^*$ are repulsive indicating that no bound states can be formed.

We observe that the effect of the u -channel diagrams are negligible due to the large mass of the exchanged doubly charmed baryon. Moreover, for the t -channel diagrams, vector meson exchange is more important than pseudoscalar meson exchange in generating bound states for $\bar{D}^{(*)}\Sigma_c^{(*)} \rightarrow \bar{D}^{(*)}\Sigma_c^{(*)}$. Except for the two partial waves S_{11} and S_{13} mentioned above, we have not observed any peak structures in other partial waves in the single channel case.

We now turn to the results of the complete calculation with all five channels included. The amplitudes squared of partial waves up to $J = 7/2$ for the $\bar{D}\Lambda_c \rightarrow \bar{D}\Lambda_c$ process are shown in Fig. 2. The relation of the dimensionless partial-wave amplitude $\tau_{\mu\nu}$ to the scattering amplitude $T_{\mu\nu}$ of Eq.(1) is:

$$\tau_{\mu\nu} = -\pi\sqrt{\rho_\mu\rho_\nu}T_{\mu\nu}, \quad (4)$$

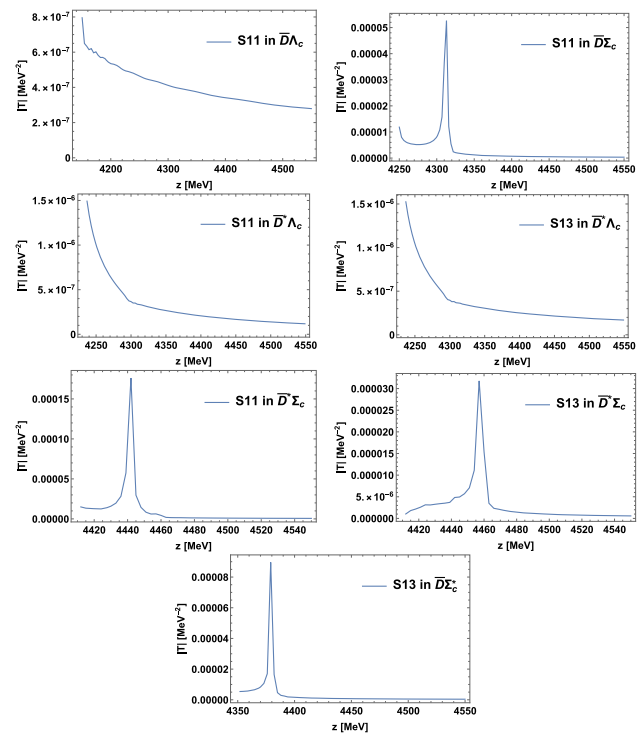


Fig. 1 The absolute value of the T -matrix for the single channel calculation of different channels in S -wave

where the phase factor is $\rho_\mu = k_\mu E_\mu \omega_\mu / z$, with k_μ , E_μ and ω_μ the on-shell three-momentum, baryon energy and meson energy of channel μ , respectively.

We need to point out that although some amplitudes are still increasing within the range we plot, all of them will eventually fall with increasing energy.

In Fig. 2, two peaks can be observed in the S_{11} partial wave. One is located just below the threshold of $\bar{D}\Sigma_c$ and the other one is below the threshold of $\bar{D}^*\Sigma_c$. Based on their positions, a reasonable assumption is that they are S -wave $\bar{D}\Sigma_c$ and $\bar{D}^*\Sigma_c$ bound states. Likewise, there are two peaks in the D_{13} wave located below the thresholds of $\bar{D}\Sigma_c^*$ and $\bar{D}^*\Sigma_c$, respectively. Thus, a similar conjecture is that they are the bound states of these two channels. It should be pointed out that here D_{13} is the name of the partial wave for $\bar{D}\Lambda_c$ channel meaning $J^P = 3/2^-$. This name can be different when talking about another channel for the same state, while J^P is always the same, c.f. Table 1. If not specifically pointed out, we use the name of the $\bar{D}\Lambda_c$ partial wave. For example, if the lower observed $J^P = 3/2^-$ state in the D_{13} partial wave is a $\bar{D}\Sigma_c^*$ bound state, it could be an S - or D -wave bound state.

In order to substantiate the above inferences and clarify the situation in the other partial waves, we extend the scattering matrix T to the complex energy plane and perform a search for dynamically generated poles on the second Riemann sheet. Note that in the present study no “genuine” s -

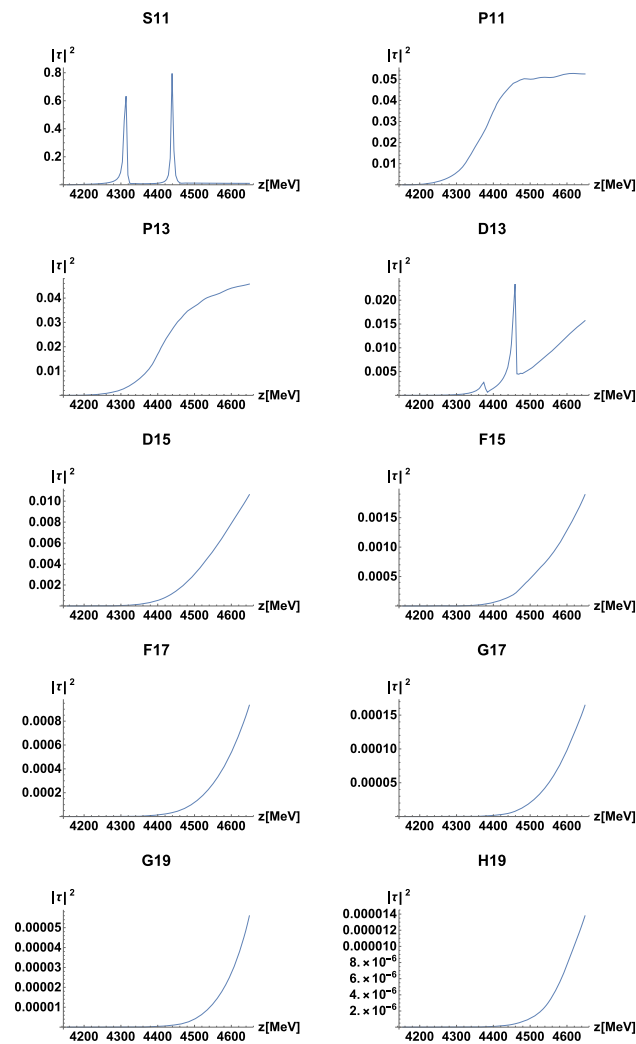


Fig. 2 Partial-wave amplitudes squared for $\bar{D}\Lambda_c \rightarrow \bar{D}\Lambda_c$ of the full coupled-channel model

channel pole diagrams are included. For details of the pole searching and the determination of the coupling strength, i.e. the residue, of a pole with respect to different channels, we refer to Refs. [33,39,45]. In Fig. 3, we present the absolute value $|T|$ of $\bar{D}\Lambda_c \rightarrow \bar{D}\Lambda_c$ for the S_{11} and D_{13} partial waves in the complex energy plane. Two peak structures are clearly visible in both cases, which are in accordance with the observed peaks in Fig. 2.

The pole positions and their couplings to different channels are listed in Table 2. It can be seen that all these four poles are very narrow with the widths between 5 to 15 MeV. The pole at $z_R = 4312.43 - i2.9$ MeV is about 8 MeV below the threshold of $\bar{D}\Sigma_c$ and it couples strongly to this channel. Thus, it can be considered as a $\bar{D}\Sigma_c$ bound state with $J^P = 1/2^-$, and it could be assigned to the $P_c(4312)$ state. Among all the couplings of the pole at $z_R = 4375.89 - i7.6$ MeV, its coupling to the $\bar{D}\Sigma_c^*$ channel is the strongest. We regard it as a $\bar{D}\Sigma_c^*$ bound state with $J^P = 3/2^-$, where the bind-

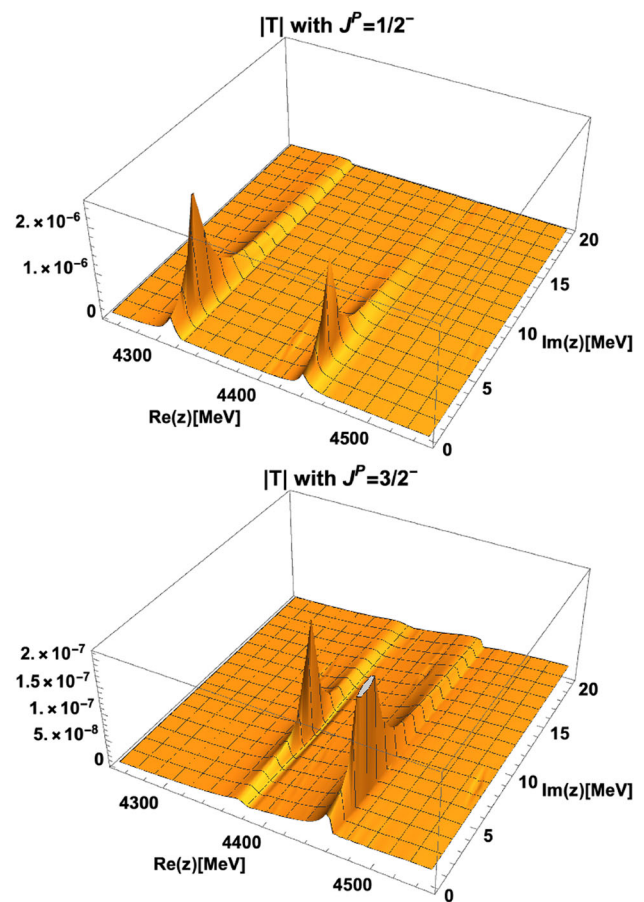


Fig. 3 Absolute value $|T|$ of $\bar{D}\Lambda_c \rightarrow \bar{D}\Lambda_c$. The upper one is the S_{11} wave and the lower one is the D_{13} partial wave

ing energy is about 10 MeV. The positions of the remaining two poles are below the threshold of $\bar{D}^*\Sigma_c$, and their couplings with this channel is the strongest. These two poles are both viewed as the S -wave $\bar{D}^*\Sigma_c$ bound states with spin-parity $J^P = 1/2^-$ and $J^P = 3/2^-$, respectively. We assign them to the observed $P_c(4440)$ and $P_c(4457)$ states. Note, however, the exact positions of these states do depend on the cut-offs employed in our calculations. Furthermore we would like to point out that conjectures about the nature of the observed states are beyond the scope of the present work.

For other partial waves, no structures are observed directly close to the real axis (c.f. Fig. 2), while we do find several poles in those partial waves located farther away from the real axis. All the poles obtained in this study have isospin $I = 1/2$, and no states with isospin $3/2$ are found. In Table 3, we present the pole positions and spin-parities for some of the observed states, which can be regarded as resonances in P - or higher partial waves with relatively large widths. The state at 4339.3 MeV having $J^P = 1/2^+$ is close to the mass of the newly observed $P_c(4337)$ state. Although the width of this state here is larger than that of the $P_c(4337)$ state, which is about 29 MeV, the situation may change after

Table 2 Pole positions z_R , spin-parities J^P and couplings g_μ for the states in S_{11} and D_{13} with $I = \frac{1}{2}$

$g_\mu [10^{-3}\text{MeV}^{-\frac{1}{2}}]$	$z_R(J^P)$			
	$4312.43 - i2.9(\frac{1}{2}^-)$	$4375.89 - i7.6(\frac{3}{2}^-)$	$4439.36 - i2.8(\frac{1}{2}^-)$	$4460.00 - i3.9(\frac{3}{2}^-)$
g_1	$1.2 + i0.15$	$0.23 - i0.014$	$0.96 + i0.026$	$0.25 - i0.30$
g_2	$7.7 + i1.0$	$0.21 + i0.25$	$-0.84 + i0.29$	$-0.15 + i0.82$
g_3	$0.023 - i0.43$	$0.053 + i0.82$	$1.2 \cdot 10^{-3} - i0.33$	$-0.13 + i0.16$
g_4	0	$0.065 + i0.77$	0	$-0.13 + i0.054$
g_5	$0.022 + i0.24$	$-0.30 + i2.5$	$0.032 + i0.067$	$0.13 - i7.3 \cdot 10^{-3}$
g_6	$-1.3 + i6.8$	$-0.95 + i1.3$	$-0.62 + i11.4$	$0.055 + i0.070$
g_7	0	$1.7 + i8.0$	0	$-0.027 - i0.13$
g_8	$-1.3 - i12$	$-1.9 + i3.9$	$-0.10 + i0.48$	$0.061 - i2.1$
g_9	0	$0.20 + i0.023$	0	$-1.1 - i1.5$
g_{10}	$0.17 + i0.49$	$10.1 + i3.6$	$-0.081 - i0.021$	$0.38 - i0.85$

Table 3 Pole positions z_R and spin-parities J^P for the states in other partial waves. Some poles with even larger widths are not listed

J^P	$z_R[\text{MeV}]$
$\frac{1}{2}^+$	$4339.3 - i106.3$
$\frac{3}{2}^+$	$4401.4 - i128.8$
$\frac{3}{2}^+$	$4463.1 - i90.1$
$\frac{5}{2}^+$	$4386.0 - i95.2$
$\frac{7}{2}^-$	$4430.8 - i214.1$

adjusting the cut-offs to the experimental data. Similarly, the state at 4386.0 MeV with $J^P = 5/2^+$, which also has a large width, could be related to the broad $P_c(4380)$ state proposed in 2015, where the preferred spin is $3/2$ or $5/2$ [1]. In Ref. [2], the existence of $P_c(4380)$ is weakened but not ruled out. Furthermore, we see inconclusive indications for two poles with even larger imaginary parts in the D_{13} and D_{15} partial waves, which are not listed in Table 3.

4 Summary

Based on the previous exploratory study, we extend the Jülich-Bonn dynamical coupled-channel model to the hidden charm sector with more channels taken into account. The channels included here are $\bar{D}\Lambda_c$, $\bar{D}\Sigma_c$, $\bar{D}^*\Lambda_c$, $\bar{D}^*\Sigma_c$, $\bar{D}\Sigma_c^*$. Predictions of partial-wave amplitudes are provided with possible dynamically generated poles examined. A series of bound states and resonances that can be assigned to the hidden-charm pentaquarks found by LHCb with different spin and parity are generated. Four S -wave bound states with isospin $I = 1/2$, whose widths are between 5 to 15 MeV, are found in the mass range of the experimentally observed P_c states. The $\bar{D}\Sigma_c$ bound state at $z_R = 4312.43 - i2.9$ MeV with $J^P = 1/2^-$ can be assigned to the $P_c(4312)$ state. The one at $z_R = 4375.89 - i7.6$ MeV is regarded as a $\bar{D}\Sigma_c^*$ bound

state with $J^P = 3/2^-$, since its coupling with this channel is the strongest, and its binding energy is about 10 MeV. The other two S -wave states are treated as the $\bar{D}^*\Sigma_c$ bound states with one having $J^P = 1/2^-$ and the other having $J^P = 3/2^-$. They are considered as the experimentally observed $P_c(4440)$ and $P_c(4457)$ states. In higher partial waves, several additional resonances with different spin and parity are also dynamically generated. However, their decay widths are very large and need further confirmation by fitting the free parameters of the model to experimental data. Such work is underway.

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Appendix A: Exchange potentials

We consider t -channel meson and u -channel baryon exchanges in this work. The corresponding Feynman diagrams are shown in Fig. 4, where the labeling of the particles and the momenta are specified.

For each exchange diagram, there is a kinematic normalization factor, an isospin factor (IF) and form factors F_1 , F_2 in addition to the pseudo-potential $\mathcal{V}$. The kinematic normalization factor is $N = \frac{1}{(2\pi)^3} \frac{1}{2\sqrt{\omega_2\omega_4}}$. The isospin factors (IF) is calculated using

$$IF(I) = \sum_{m_1, m_2, m_3, m_4} \langle I_1 m_1 I_2 m_2 | II_z \rangle \langle I_3 m_3 I_4 m_4 | II_z \rangle \times \langle I_3 m_3 I_4 m_4 | M^{iso} | I_1 m_1 I_2 m_2 \rangle, \quad (\text{A.1})$$

where the notations and details can be found in Ref. [50]. All the isospin factors for involved processes are given in Table 4. Note that the baryon-first convention is always applied in the isospin calculation. At each vertex, a form factor is applied, its specific form is given by

$$F(q) = \left(\frac{\Lambda^2 - m_{ex}^2}{\Lambda^2 + \mathbf{q}^2} \right)^n, \quad (\text{A.2})$$

where m_{ex} , $\mathbf{q}$ and Λ are the mass, the three-momentum of the exchanged particle and the cut-off, respectively. The powers $n = 1, 2$ correspond to monopole or dipole form factors. Different powers for the form factors are used to ensure that the integral over the off-shell momenta in the scattering equation is convergent. For the t -channel pseudoscalar and vector meson exchange, we use a dipole form factor, i.e. $n = 2$, while a monopole form factor is applied for u -channel Ξ_{cc} exchange, i.e. $n = 1$. The transition amplitude is then expressed by

$$V_{\lambda_1, \lambda_2, \lambda_3, \lambda_4}^{\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3, \mathbf{p}_4} = N F_1 F_2 (\text{IF}) \mathcal{V}. \quad (\text{A.3})$$

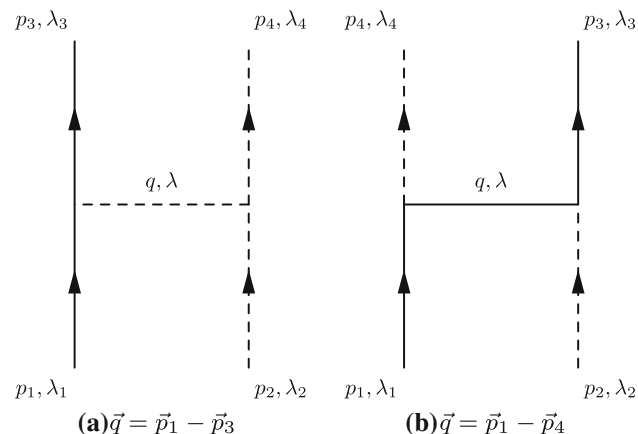


Fig. 4 t - and u -channel exchange processes

Table 4 The isospin factors for all the involved processes in this work

Process	Exchanged particle	IF($\frac{1}{2}$)	IF($\frac{3}{2}$)
$\Lambda_c \bar{D}^{(*)} \rightarrow \Lambda_c \bar{D}^{(*)}$	ω	1	0
	Ξ_{cc}	1	0
$\Lambda_c \bar{D}^{(*)} \rightarrow \Sigma_c^{(*)} \bar{D}^{(*)}$	π/ρ	$-\sqrt{3}$	0
	Ξ_{cc}	$\sqrt{3}$	0
$\Sigma_c^{(*)} \bar{D}^{(*)} \rightarrow \Sigma_c^{(*)} \bar{D}^{(*)}$	$\eta/\eta'/\omega$	1	1
	π/ρ	-2	1
	Ξ_{cc}	-1	2

The detailed expressions of all the needed potentials in this work are presented below.

1. Amplitudes

There are in total 15 different kinds of exchange pseudo-potentials $\mathcal{V}$ considered here. As clarified in Eq. (A.3), other factors including the normalization factor, isospin factors and form factors have been isolated. Here, we present the explicit expressions of all the needed pseudo-potentials using the notations in Fig. 4.

Type 1: $BP \rightarrow BP$ via vector meson exchange,

$$\mathcal{V} = -g_a g_b \bar{u}_3 \left(\frac{\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} q_\nu}{z - \omega_t - \omega_2 - E_3} + \frac{\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} \tilde{q}_\nu}{z - \omega_t - E_1 - \omega_4} \right) \times \frac{(p_2 + p_4)_\mu}{2\omega_t} u_1. \quad (\text{A.4})$$

Type 2: $BP \rightarrow BP$ via baryon exchange,

$$\mathcal{V} = \frac{g_a g_b}{m_p^2} \bar{u}_3 \gamma^5 p_2 \left(\frac{\not{q} + m}{z - E_u - \omega_2 - \omega_4} + \frac{\not{\tilde{q}} + m}{z - E_u - E_1 - E_3} \right) \times \frac{1}{2E_u} \gamma^5 p_4 u_1. \quad (\text{A.5})$$

Type 3: $BP \rightarrow BV$ via pseudoscalar meson exchange,

$$\mathcal{V} = i \frac{g_a g_b}{m_p} \bar{u}_3 \gamma^5 \left(\frac{\not{q} (p_2 - q)_\mu}{z - \omega_t - \omega_2 - E_3} + \frac{\not{\tilde{q}} (p_2 - \tilde{q})_\mu}{z - \omega_t - E_1 - \omega_4} - 2\omega_t \gamma^0 \delta_\mu^0 \right) \times \frac{1}{2\omega_t} u_1 \epsilon_4^{*\mu}. \quad (\text{A.6})$$

Type 4: $BP \rightarrow BV$ via vector meson exchange,

$$\mathcal{V} = \frac{g_a g_b}{m_V} \bar{u}_3 \left(\frac{\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} q_\nu}{z - \omega_t - \omega_2 - E_3} + \frac{\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} \tilde{q}_\nu}{z - \omega_t - E_1 - \omega_4} \right) \times \frac{1}{2\omega_t} u_1 \epsilon_{\alpha\beta\lambda\mu} p_4^\alpha \epsilon_4^{*\beta} p_2^\lambda. \quad (\text{A.7})$$

Type 5: $BP \rightarrow BV$ via baryon exchange,

$$\mathcal{V} = -i \frac{g_a g_b}{m_P} \bar{u}_3 \gamma^5 p_2 \left(\frac{\not{q} + m}{z - E_u - \omega_2 - \omega_4} + \frac{\not{\tilde{q}} + m}{z - E_u - E_1 - E_3} \right) \frac{1}{2E_u} \times \left(\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} p_{4\nu} \right) \epsilon_{4\mu}^* u_1. \quad (\text{A.8})$$

Type 6: $BP \rightarrow DP$ via vector meson exchange,

$$\mathcal{V} = \frac{g_a g_b}{m_V} \bar{u}_3^\mu \gamma^5 \left(\frac{\gamma_\nu q_\mu - \not{q} g_{\mu\nu}}{z - \omega_t - \omega_2 - E_3} + \frac{\gamma_\nu \tilde{q}_\mu - \not{\tilde{q}} g_{\mu\nu}}{z - \omega_t - E_1 - \omega_4} \right) \frac{(p_2 + p_4)^\nu}{2\omega_t} u_1. \quad (\text{A.9})$$

Type 7: $BP \rightarrow DP$ via baryon exchange,

$$\mathcal{V} = \frac{g_a g_b}{m_P^2} \bar{u}_3^\mu \left(\frac{\not{q} + m}{z - E_u - \omega_2 - \omega_4} + \frac{\not{\tilde{q}} + m}{z - E_u - E_1 - E_3} \right) \times \frac{p_{2\mu}}{2E_u} \gamma^5 p_4 u_1. \quad (\text{A.10})$$

Type 8: $BV \rightarrow BV$ via vector meson exchange,

$$\mathcal{V} = -g_a g_b \bar{u}_3 \left[\frac{(\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} q_\nu) [\epsilon_2^\nu \epsilon_{4\mu}^* (p_4 + q)_\nu - \epsilon_2^\nu \epsilon_{4\nu}^* (p_2 + p_4)_\mu + \epsilon_{2\mu} \epsilon_4^{*\nu} (p_2 - q)_\nu]}{z - \omega_t - \omega_2 - E_3} + \frac{q \rightarrow \tilde{q}}{z - \omega_t - E_1 - \omega_4} \right] \frac{1}{2\omega_t} u_1. \quad (\text{A.11})$$

Type 9: $BV \rightarrow BV$ via pseudoscalar meson exchange,

$$\mathcal{V} = -i \frac{g_a g_b}{m_P m_V} \bar{u}_3 \gamma^5 \left(\frac{\not{q}}{z - \omega_t - \omega_2 - E_3} + \frac{\not{\tilde{q}}}{z - \omega_t - E_1 - \omega_4} \right) \frac{1}{2\omega_t} u_1 \epsilon_{\mu\nu\lambda\tau} p_2^\mu \epsilon_2^\nu p_4^\lambda \epsilon_4^{*\tau}. \quad (\text{A.12})$$

Type 10: $BV \rightarrow BV$ via baryon exchange,

$$\mathcal{V} = g_a g_b \bar{u}_3 (\not{\epsilon}_2 + i \frac{\kappa}{2m} \sigma^{\mu\nu} p_{2\nu} \epsilon_{2\mu}) \left(\frac{\not{q} + m}{z - E_u - \omega_2 - \omega_4} + \frac{\not{\tilde{q}} + m}{z - E_u - E_1 - E_3} \right) \frac{1}{2E_u} (\epsilon_4^* - i \frac{\kappa}{2m} \sigma^{\alpha\beta} p_{4\beta} \epsilon_{4\alpha}^*) u_1. \quad (\text{A.13})$$

Type 11: $BV \rightarrow DP$ via pseudoscalar meson exchange,

$$\mathcal{V} = i \frac{g_a g_b}{m_P} \bar{u}_3^\mu \left(\frac{q_\mu (p_4 + q)_\nu}{z - \omega_t - \omega_2 - E_3} + \frac{\tilde{q}_\mu (p_4 + \tilde{q})_\nu}{z - \omega_t - E_1 - \omega_4} + 2\omega_t \delta_\mu^0 \delta_\nu^0 \right) \times \frac{1}{2\omega_t} u_1 \epsilon_2^\nu. \quad (\text{A.14})$$

Type 12: $BV \rightarrow DP$ via vector meson exchange,

$$\mathcal{V} = \frac{g_a g_b}{m_V^2} \bar{u}_{3,\mu} \gamma^5 \left(\frac{(\gamma^\beta q^\mu - \not{q} g^{\mu\beta}) q^\alpha}{z - \omega_t - \omega_2 - E_3} + \frac{(\gamma^\beta \tilde{q}^\mu - \not{\tilde{q}} g^{\mu\beta}) \tilde{q}^\alpha}{z - \omega_t - E_1 - \omega_4} \right) \times \frac{1}{2\omega_t} u_1 \epsilon_{\alpha\beta\lambda\tau} p_2^\lambda \epsilon_2^\tau. \quad (\text{A.15})$$

Type 13: $BV \rightarrow DP$ via baryon exchange,

$$\mathcal{V} = i \frac{g_a g_b}{m_P m_V} \bar{u}_3^\mu \gamma^5 (p_{2\mu} \not{\epsilon}_2 - \epsilon_{2\mu} \not{p}_2) \left(\frac{\not{q} + m}{z - E_u - \omega_2 - \omega_4} + \frac{\not{\tilde{q}} + m}{z - E_u - E_1 - E_3} \right) \frac{1}{2E_u} \gamma^5 p_4 u_1. \quad (\text{A.16})$$

Type 14: $DP \rightarrow DP$ via vector meson exchange,

$$\mathcal{V} = g_a g_b \bar{u}_3^\tau \left(\frac{\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} q_\nu}{z - \omega_t - \omega_2 - E_3} + \frac{\gamma^\mu - i \frac{\kappa}{2m} \sigma^{\mu\nu} \tilde{q}_\nu}{z - \omega_t - E_1 - \omega_4} \right) \times \frac{(p_2 + p_4)_\mu}{2\omega_t} u_{1,\tau}. \quad (\text{A.17})$$

Type 15: $DP \rightarrow DP$ via baryon exchange,

$$\mathcal{V} = \frac{g_a g_b}{m_p^2} \bar{u}_3^\mu \left(\frac{\not{q} + m}{z - E_u - \omega_2 - \omega_4} + \frac{\not{\tilde{q}} + m}{z - E_u - E_1 - E_3} \right) \frac{p_{2\mu} p_{4\nu}}{2E_u} u_1^\nu. \quad (\text{A.18})$$

Here, q is the exchanged momentum in the first time ordering and $\tilde{q}$ means the second time ordering with $\tilde{q}^0 = -q^0$. The indices 1 and 3 (2 and 4) denote the incoming and outgoing baryon (meson). E_i and ω_i are the on-shell energies for a baryon and a meson, respectively. A “type” number is introduced to identify which potential each process corresponds to.

In Table 5, we list all the processes with the exchanged particle considered in the calculation, and the corresponding “type” of each specific process together with the coupling constants $g_{a,b}$ and cut-off values in the form factor are also given. Note that, for simplicity, we take one cut-off for each process, not for each vertex. For the t -channel pseudoscalar meson exchange and u -channel baryon exchange, we take a uniform value for the same exchanged particle. We adjust the cut-offs for the t -channel vector meson exchanges to obtain positions of the dynamically generated poles near the experimentally observed P_c states. As mentioned above, here we emphasize again that our numerical results depend on the cut-off values, and the results given in this work are for this specific set of cut-offs. In the cases where vector mesons couple to octet baryons and vector mesons couple to decuplet baryons, there is also a tensor coupling $f_b = g_b \kappa$, which is not listed in Table 5, but should be consistent with the vector coupling g_b . The explicit relations of the coupling constants g are given below.

2. Coupling constants

The coupling constants for all the vertices in the calculation are related to each other by SU(3) and SU(4) flavor symmetry [48, 49]. The employed relations of the various coupling constants are given by the following expressions:

- Couplings for pseudoscalar meson, pseudoscalar meson and vector meson:

$$\begin{aligned} g_{DD\rho} &= g_{PPV}, \\ g_{DD\omega} &= g_{PPV}, \\ g_{D\pi D^*} &= -g_{PPV}, \\ g_{D\eta D^*} &= -\frac{1}{\sqrt{3}} g_{PPV}, \\ g_{D\eta' D^*} &= -\sqrt{\frac{2}{3}} g_{PPV}, \end{aligned}$$

$$g_{DDJ/\psi} = -\sqrt{2} g_{PPV}, \quad (\text{A.19})$$

where $g_{PPV} = 3.018$ [51].

- Couplings for vector meson, vector meson and pseudoscalar meson:

$$\begin{aligned} g_{D^* \rho D} &= \frac{1}{\sqrt{2}} g_{VVP}, \\ g_{D^* \omega D} &= \frac{1}{\sqrt{2}} g_{VVP}, \\ g_{D^* D^* \pi} &= \frac{1}{\sqrt{2}} g_{VVP}, \\ g_{D^* D^* \eta} &= \frac{1}{\sqrt{6}} g_{VVP}, \\ g_{D^* D^* \eta'} &= \frac{1}{\sqrt{3}} g_{VVP}, \\ g_{D^* J/\psi D} &= g_{VVP}, \end{aligned} \quad (\text{A.20})$$

where $g_{VVP} = -7.070$ [39].

- Couplings for vector meson, vector meson and vector meson:

$$\begin{aligned} g_{D^* D^* \rho} &= -\frac{1}{\sqrt{2}} g_{VVV}, \\ g_{D^* D^* \omega} &= -\frac{1}{\sqrt{2}} g_{VVV}, \\ g_{D^* D^* J/\psi} &= g_{VVV}, \end{aligned} \quad (\text{A.21})$$

where $g_{VVV} = 2.298$ [39].

- Couplings for octet baryon, octet baryon and pseudoscalar meson:

$$\begin{aligned} g_{\Lambda_c N D} &= -\frac{3\sqrt{3}}{5} g_{BBP}, \\ g_{\Sigma_c N D} &= \frac{1}{5} g_{BBP}, \\ g_{\Sigma_c \Lambda_c \pi} &= \frac{2\sqrt{3}}{5} g_{BBP}, \\ g_{\Sigma_c \Sigma_c \pi} &= \frac{4}{5} g_{BBP}, \\ g_{\Sigma_c \Sigma_c \eta} &= \frac{4}{5\sqrt{3}} g_{BBP}, \\ g_{\Sigma_c \Sigma_c \eta'} &= \frac{4\sqrt{2}}{5\sqrt{3}} g_{BBP}, \\ g_{\Xi_{cc} \Lambda_c D} &= \frac{\sqrt{3}}{5} g_{BBP}, \\ g_{\Xi_{cc} \Sigma_c D} &= -g_{BBP}, \end{aligned} \quad (\text{A.22})$$

where $g_{BBP} = 0.989$ [52].

Table 5 The types of potentials, the couplings and the cut-offs in the form factor for all the considered processes with the exchanged particle (Ex)

Process	Ex	Type	g_a	g_b	Λ [MeV]
$\Lambda_c \bar{D} \rightarrow \Lambda_c \bar{D}$	ω	1	$g_{DD\omega}$	$g_{\Lambda_c \Lambda_c \omega}$	2800
	Ξ_{cc}	2	$g_{\Xi_{cc} \Lambda_c D}$	$g_{\Xi_{cc} \Lambda_c D}$	3000
$\Lambda_c \bar{D} \rightarrow \Sigma_c \bar{D}$	ρ	1	$g_{DD\rho}$	$g_{\Sigma_c \Lambda_c \rho}$	2800
	Ξ_{cc}	2	$g_{\Xi_{cc} \Sigma_c D}$	$g_{\Xi_{cc} \Lambda_c D}$	3000
$\Lambda_c \bar{D} \rightarrow \Lambda_c \bar{D}^*$	ω	4	$g_{D^* \omega D}$	$g_{\Lambda_c \Lambda_c \omega}$	1500
	Ξ_{cc}	5	$g_{\Xi_{cc} \Lambda_c D}$	$g_{\Xi_{cc} \Lambda_c D^*}$	3000
$\Lambda_c \bar{D} \rightarrow \Sigma_c \bar{D}^*$	π	3	$g_{D\pi D^*}$	$g_{\Sigma_c \Lambda_c \pi}$	800
	ρ	4	$g_{D^* \rho D}$	$g_{\Sigma_c \Lambda_c \rho}$	1800
	Ξ_{cc}	5	$g_{\Xi_{cc} \Sigma_c D}$	$g_{\Xi_{cc} \Lambda_c D^*}$	3000
$\Lambda_c \bar{D} \rightarrow \Sigma_c^* \bar{D}$	ρ	6	$g_{DD\rho}$	$g_{\Sigma_c^* \Lambda_c \rho}$	1800
	Ξ_{cc}	7	$g_{\Sigma_c^* \Xi_{cc} D}$	$g_{\Xi_{cc} \Lambda_c D}$	3000
$\Sigma_c \bar{D} \rightarrow \Sigma_c \bar{D}$	ρ	1	$g_{DD\rho}$	$g_{\Sigma_c \Sigma_c \rho}$	2700
	ω	1	$g_{DD\omega}$	$g_{\Sigma_c \Sigma_c \omega}$	2700
	Ξ_{cc}	2	$g_{\Xi_{cc} \Sigma_c D}$	$g_{\Xi_{cc} \Sigma_c D}$	3000
$\Sigma_c \bar{D} \rightarrow \Lambda_c \bar{D}^*$	π	3	$g_{D\pi D^*}$	$g_{\Sigma_c \Lambda_c \pi}$	800
	ρ	4	$g_{D^* \rho D}$	$g_{\Sigma_c \Lambda_c \rho}$	1500
	Ξ_{cc}	5	$g_{\Xi_{cc} \Lambda_c D}$	$g_{\Xi_{cc} \Sigma_c D^*}$	3000
$\Sigma_c \bar{D} \rightarrow \Sigma_c \bar{D}^*$	π	3	$g_{D\pi D^*}$	$g_{\Sigma_c \Sigma_c \pi}$	800
	η	3	$g_{D\eta D^*}$	$g_{\Sigma_c \Sigma_c \eta}$	1000
	η'	3	$g_{D\eta' D^*}$	$g_{\Sigma_c \Sigma_c \eta'}$	1000
	ρ	4	$g_{D^* \rho D}$	$g_{\Sigma_c \Sigma_c \rho}$	1500
	ω	4	$g_{D^* \omega D}$	$g_{\Sigma_c \Sigma_c \omega}$	1500
	Ξ_{cc}	5	$g_{\Xi_{cc} \Sigma_c D}$	$g_{\Xi_{cc} \Sigma_c D^*}$	3000
	ρ	6	$g_{DD\rho}$	$g_{\Sigma_c^* \Sigma_c \rho}$	1500
	ω	6	$g_{DD\omega}$	$g_{\Sigma_c^* \Sigma_c \omega}$	1500
$\Sigma_c \bar{D} \rightarrow \Sigma_c^* \bar{D}$	Ξ_{cc}	7	$g_{\Sigma_c^* \Xi_{cc} D}$	$g_{\Xi_{cc} \Sigma_c D}$	3000
	ω	8	$g_{D^* D^* \omega}$	$g_{\Lambda_c \Lambda_c \omega}$	2000
$\Lambda_c \bar{D}^* \rightarrow \Lambda_c \bar{D}^*$	Ξ_{cc}	10	$g_{\Xi_{cc} \Lambda_c D^*}$	$g_{\Xi_{cc} \Lambda_c D^*}$	3000
	π	9	$g_{D^* D^* \pi}$	$g_{\Sigma_c \Lambda_c \pi}$	800
$\Lambda_c \bar{D}^* \rightarrow \Sigma_c \bar{D}^*$	ρ	8	$g_{D^* D^* \rho}$	$g_{\Sigma_c \Lambda_c \rho}$	3000
	Ξ_{cc}	10	$g_{\Xi_{cc} \Sigma_c D^*}$	$g_{\Xi_{cc} \Lambda_c D^*}$	3000
$\Lambda_c \bar{D}^* \rightarrow \Sigma_c^* \bar{D}$	π	11	$g_{D\pi D^*}$	$g_{\Sigma_c^* \Lambda_c \pi}$	800
	ρ	12	$g_{D^* \rho D}$	$g_{\Sigma_c^* \Lambda_c \rho}$	1300
	Ξ_{cc}	13	$g_{\Sigma_c^* \Xi_{cc} D^*}$	$g_{\Xi_{cc} \Lambda_c D}$	3000
$\Sigma_c \bar{D}^* \rightarrow \Sigma_c \bar{D}^*$	π	9	$g_{D^* D^* \pi}$	$g_{\Sigma_c \Sigma_c \pi}$	800
	η	9	$g_{D^* D^* \eta}$	$g_{\Sigma_c \Sigma_c \eta}$	1000
	η'	9	$g_{D^* D^* \eta'}$	$g_{\Sigma_c \Sigma_c \eta'}$	1000
	ρ	8	$g_{D^* D^* \rho}$	$g_{\Sigma_c \Sigma_c \rho}$	2600
	ω	8	$g_{D^* D^* \omega}$	$g_{\Sigma_c \Sigma_c \omega}$	1000
	Ξ_{cc}	10	$g_{\Xi_{cc} \Sigma_c D^*}$	$g_{\Xi_{cc} \Sigma_c D^*}$	3000
	π	11	$g_{D\pi D^*}$	$g_{\Sigma_c^* \Sigma_c \pi}$	800

Table 5 continued

Process	Ex	Type	g_a	g_b	Λ [MeV]
$\Sigma_c^* \bar{D} \rightarrow \Sigma_c^* \bar{D}$	η	11	$g_{D\eta D^*}$	$g_{\Sigma_c^* \Sigma_c \eta}$	1000
	η'	11	$g_{D\eta' D^*}$	$g_{\Sigma_c^* \Sigma_c \eta'}$	1000
	ρ	12	$g_{D^* \rho D}$	$g_{\Sigma_c^* \Sigma_c \rho}$	1500
	ω	12	$g_{D^* \omega D}$	$g_{\Sigma_c^* \Sigma_c \omega}$	1500
	Ξ_{cc}	13	$g_{\Sigma_c^* \Xi_{cc} D^*}$	$g_{\Xi_{cc} \Sigma_c D}$	3000
	ρ	14	$g_{DD\rho}$	$g_{\Sigma_c^* \Sigma_c^* \rho}$	1700
	ω	14	$g_{DD\omega}$	$g_{\Sigma_c^* \Sigma_c^* \omega}$	1700
	Ξ_{cc}	15	$g_{\Sigma_c^* \Xi_{cc} D}$	$g_{\Sigma_c^* \Xi_{cc} D}$	3000

- Vector couplings for octet baryon, octet baryon and vector meson:

$$\begin{aligned}
 g_{\Lambda_c N D^*} &= -\frac{1}{\sqrt{3}} g_{BBV} (1 + 2\alpha_{BBV}), \\
 g_{\Sigma_c N D^*} &= g_{BBV} (1 - 2\alpha_{BBV}), \\
 g_{\Lambda_c \Lambda_c \omega} &= \frac{2}{3} g_{BBV} (5\alpha_{BBV} - 2), \\
 g_{\Lambda_c \Lambda_c J/\psi} &= \sqrt{2} g_{BBV}, \\
 g_{\Sigma_c \Lambda_c \rho} &= \frac{2}{\sqrt{3}} g_{BBV} (1 - \alpha_{BBV}), \\
 g_{\Sigma_c \Sigma_c \rho} &= 2g_{BBV} \alpha_{BBV}, \\
 g_{\Sigma_c \Sigma_c \omega} &= 2g_{BBV} \alpha_{BBV}, \\
 g_{\Sigma_c \Sigma_c J/\psi} &= \sqrt{2} g_{BBV}, \\
 g_{\Xi_{cc} \Lambda_c D^*} &= \frac{1}{\sqrt{3}} g_{BBV} (4\alpha_{BBV} - 1), \\
 g_{\Xi_{cc} \Sigma_c D^*} &= -g_{BBV},
 \end{aligned} \tag{A.23}$$

where $g_{BBV} = 3.249$ [52], $\alpha_{BBV} = 1.15$ [39].

- Tensor couplings for octet baryon, octet baryon and vector meson:

$$\begin{aligned}
 f_{\Lambda_c N D^*} &= -\frac{1}{2\sqrt{3}} f_{NN\omega} - \frac{\sqrt{3}}{2} f_{NN\rho}, \\
 f_{\Sigma_c N D^*} &= -\frac{1}{2} f_{NN\omega} + \frac{1}{2} f_{NN\rho}, \\
 f_{\Lambda_c \Lambda_c \omega} &= \frac{5}{6} f_{NN\omega} - \frac{1}{2} f_{NN\rho}, \\
 f_{\Lambda_c \Lambda_c J/\psi} &= 0, \\
 f_{\Sigma_c \Lambda_c \rho} &= -\frac{1}{2\sqrt{3}} f_{NN\omega} + \frac{\sqrt{3}}{2} f_{NN\rho}, \\
 f_{\Sigma_c \Sigma_c \rho} &= \frac{1}{2} f_{NN\omega} + \frac{1}{2} f_{NN\rho}, \\
 f_{\Sigma_c \Sigma_c \omega} &= \frac{1}{2} f_{NN\omega} + \frac{1}{2} f_{NN\rho}, \\
 f_{\Sigma_c \Sigma_c J/\psi} &= 0,
 \end{aligned}$$

$$f_{\Xi_{cc} \Lambda_c D^*} = 0,$$

$$f_{\Xi_{cc} \Sigma_c D^*} = 0, \tag{A.24}$$

where $f_{NN\rho} = 19.819$, $f_{NN\omega} = 0$ [52].

- Couplings for decuplet baryon, octet baryon and pseudoscalar meson:

$$\begin{aligned}
 g_{\Sigma_c^* N D} &= -\frac{1}{\sqrt{6}} g_{BDP}, \\
 g_{\Sigma_c^* \Lambda_c \pi} &= \frac{1}{\sqrt{2}} g_{BDP}, \\
 g_{\Sigma_c^* \Sigma_c \pi} &= \frac{1}{\sqrt{6}} g_{BDP}, \\
 g_{\Sigma_c^* \Sigma_c \eta} &= \frac{1}{2\sqrt{2}} g_{BDP}, \\
 g_{\Sigma_c^* \Sigma_c \eta'} &= \frac{1}{3} g_{BDP}, \\
 g_{\Sigma_c^* \Xi_{cc} D} &= -\frac{1}{\sqrt{6}} g_{BDP},
 \end{aligned} \tag{A.25}$$

where $g_{BDP} = 2.127$ [52].

- Couplings for decuplet baryon, octet baryon and vector meson:

$$\begin{aligned}
 g_{\Sigma_c^* N D^*} &= -\frac{1}{\sqrt{6}} g_{BDV}, \\
 g_{\Sigma_c^* \Lambda_c \rho} &= \frac{1}{\sqrt{2}} g_{BDV}, \\
 g_{\Sigma_c^* \Sigma_c \rho} &= \frac{1}{\sqrt{6}} g_{BDV}, \\
 g_{\Sigma_c^* \Sigma_c \omega} &= -\frac{1}{\sqrt{6}} g_{BDV}, \\
 g_{\Sigma_c^* \Sigma_c J/\psi} &= -\frac{1}{\sqrt{3}} g_{BDV}, \\
 g_{\Sigma_c^* \Xi_{cc} D^*} &= -\frac{1}{\sqrt{6}} g_{BDV},
 \end{aligned} \tag{A.26}$$

where $g_{BDV} = 16.031$ [52].

- Vector and tensor couplings for decuplet baryon, decuplet baryon and vector meson:

$$\begin{aligned}
 g_{\Sigma_c^* \Sigma_c^* \rho} &= 2g_{DDV}, \\
 g_{\Sigma_c^* \Sigma_c^* \omega} &= 2g_{DDV}, \\
 f_{\Sigma_c^* \Sigma_c^* \rho} &= \frac{1}{2} \kappa_\rho g_{DDV}, \\
 f_{\Sigma_c^* \Sigma_c^* \omega} &= \frac{1}{2} \kappa_\rho g_{DDV},
 \end{aligned} \tag{A.27}$$

where $g_{DDV} = 7.674$ and $\kappa_\rho = 6.1$ [53,54].

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